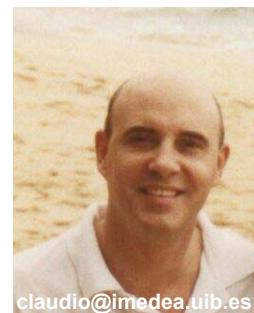
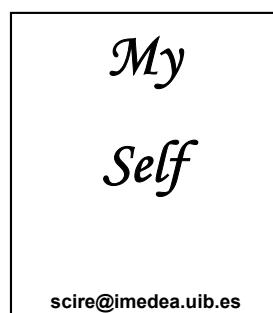


Two-mode dynamics in different semiconductor laser structures



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GLASGOW



- **Vertical Cavity Surface Emitting Lasers**

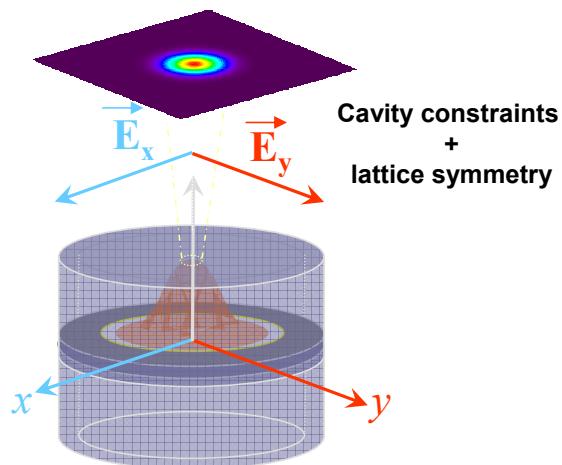


- **Twin Stripes Semiconductor lasers**

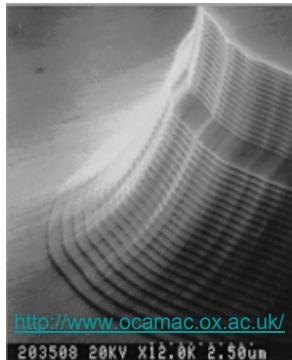
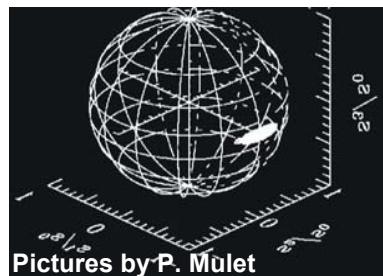
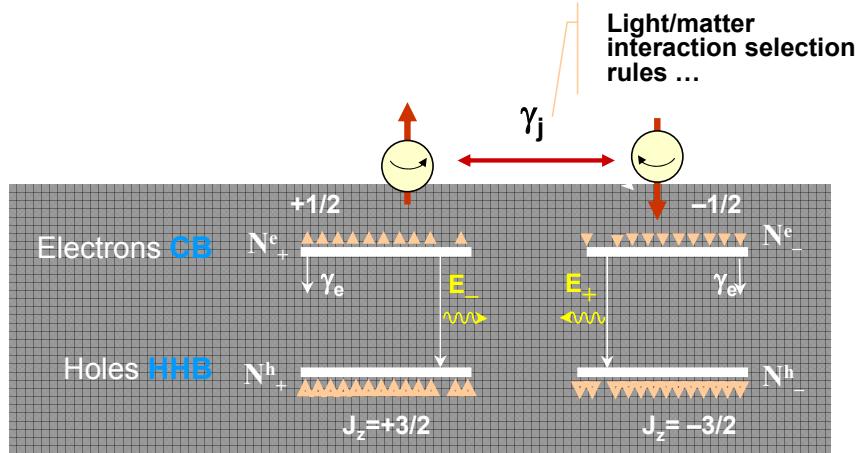


- **Semiconductor Ring-lasers**

VCSELs – two-polarization modes



Cavity constraints
+
lattice symmetry



Polarization switching



Intensity
instab.

Phase
instab.

Thermal (lattice)

Material gain

- gain shift
- gain/loss shifts
- Strain (thermo-elastic)

Non-thermal

- Carrier heating
- type I PS in SFM

Modal gain

- Thermal lensing

- Spatial hole burning

No gain mech.

Spin Flip + α -factor
(Pol./Phase/Ampl-coupling)
Type II PS

G.Verschaffelt, et al.

types of PS in # types of VCSELs. VISTA RTN



VCSELs: Proposed theoretical descriptions *Spin-Flip Model*

1. Processes and features

- SVA.Polarization d. of f.
 - Intensity-Phase-polarization interaction
(cavity-anisotropies & alpha-factor)
 - Polarization resolved Light/matter int.selection rules

2. Achievements

- Good description of some PS processes
(non-thermal PS, phase instabilities)
- Good description of intensity/phase/polarization non-linear dynamics

Extensions:

- multi-transverse mode dynamics
(mesoscopic susceptibility beyond the two-level approach)
- External cavity, phase conjugated feedback, ext. Injection, sat.absorber.
- *Two mode rate equations:*

Explain jitter (variances)

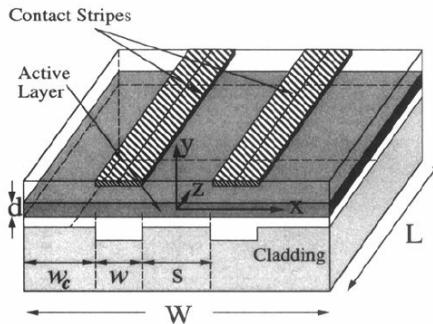
Investigate the switching time statistics

frequency-limit current-induced PS

Origin of PS in VCSELs

Relationship SFM-TMRE
Van der Sande et al, PRA 2003.

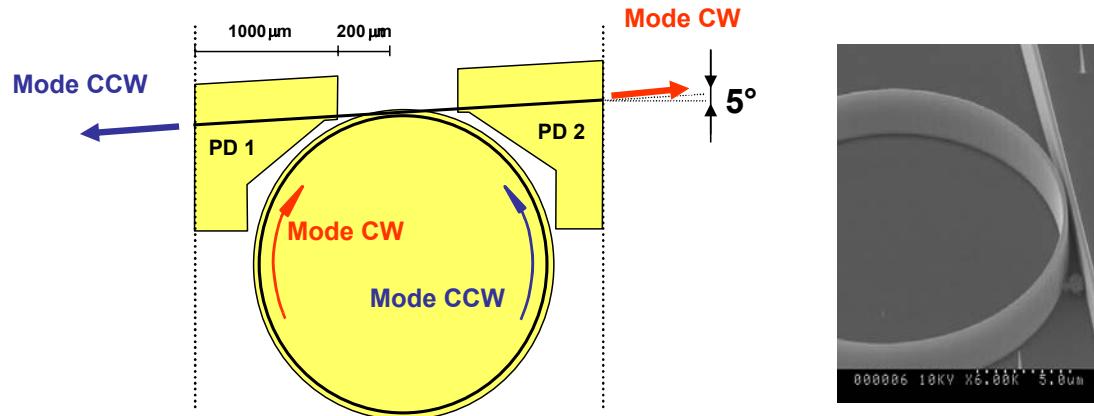
- Twin Stripe lasers: two-spatial modes
- Fast two-mode dynamics (e.g. anti-phase current driving)
- Bistability
- FALCON-TMR



[O.Hess, E.Scholl, PRA 50, 787 1994]

[I.Paiss, A.Hardy, IEEE JQE 25, 1609 1989]

- Semiconductor Ring-lasers: two-counterpropagating-modes
- Bistability
- Inertial Rotation Sensors



[M.Sorel et al. Opt.Lett. 27 1992 (2002)]
[...Nature...]

VCSEL

Spin-Flip model
Four levels + spinflip
 $\text{div}(E)=0$
Linear anisotropy

Rate Eqs.
 $\gamma_s = \infty$

$\gamma < 1$
[M. San Miguel
PRL 75, 425 (1995)]

Maxwell – Bloch Eqs.

Two-levels $\chi + \alpha$ -factor
+ phen. Dissipation
 $\delta_t P = 0$

Single mode operation
SVEA
Mean field

Twin Stripes

Twin-Stripe.TwoModeSVA
Coupled-Mode theory

VC Ginzburg-Landau Eq

[I.R.Aranson Rev.Mod.Phys. 74 99 (2002)]

Time scales < $f_{\text{rel.ox}}$, $\delta_t N = 0$

$\gamma \ll 1$

Sem. Ring Laser

Ring.TwoModeSVA
Perturbative symmetry br.

$\gamma = 2$ (!)...

VCGLE $\dot{E}_{1,2} = (\mu + i\omega)E_{1,2} + (b + i\theta)\left(\left|E_{1,2}\right|^2 + \gamma\left|E_{2,1}\right|^2\right)E_{1,2} - (k_d + ik_c)E_{2,1}$

Cross-saturation
Linear-field-coupling

Gain Cross-saturation in SRL

Given form of $\delta\hat{N}(z)$ can pick out any synchronous terms besides the unsaturated gain:

$$\text{#12.5: } \delta\hat{N}(z) = \frac{I_f + I_b}{I_{sat}} + \frac{\sqrt{I_f I_b}}{I_{sat}} (e^{-i2kz} + e^{i2kz})$$

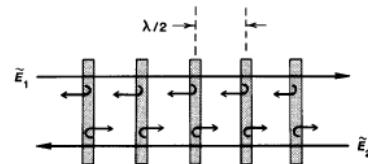
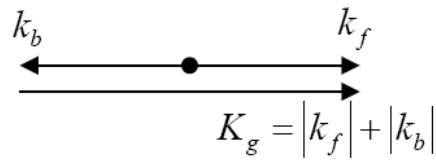
Besides usual uniform self- and cross-saturation terms, circled terms also synchronous

$$\frac{dI_f}{dz} = g_0 \left[1 - \frac{I_f}{I_{sat}} - \frac{2I_b}{I_{sat}} \right] I_f$$

Leads to cross-saturation twice as strong as self-saturation: major impact on stability analysis

Result of coherent scattering off the Bragg-matched grating

Similar phenomena appear in many contexts: photorefractives, FWM, ...



www.stanford.edu

Gain Cross-saturation in VCSEL

- Cross-sat is dynamically included through spin-deviation var. $d(t)$ in SFM model
Spin deviations affects polarization, coupled to phase (γ_p) and to amplitude (α).
- In two mode rate equations cross sat. Coef. are euristically included or derived from SFM ...

$$\begin{aligned}\mathcal{E}_{xy} &= \frac{2\kappa\gamma_s}{\gamma_s^2 + 4\gamma_p^2} \left(1 - 2\alpha \frac{\gamma_p}{\gamma_s} \right) \\ \mathcal{E}_{yx} &= \frac{2\kappa\gamma_s}{\gamma_s^2 + 4\gamma_p^2} \left(1 + 2\alpha \frac{\gamma_p}{\gamma_s} \right)\end{aligned}$$

Van Exter et al, JOB 1999.

Van der Sande et al, PRA 2003.

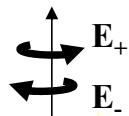
Gain Cross-saturation in TSs

- Often neglected ($\gamma=0$)
- Incoherent coupling due to Cross-carrier diffusion

Complex linear fields coupling

VCSEL

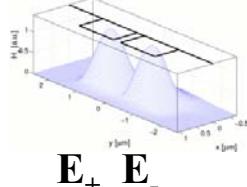
Spin-Flip. RE model
Circular pol. basis



- Cavity anisotropy

Twin Stripes

Twin-Stripe.TwoModeSVA
Local fields



E_+ E_-

- Diffraction - Evanescent wave

Sem. Ring Laser

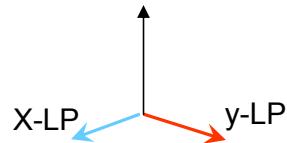
Ring.TwoModeSVA
Travelling waves



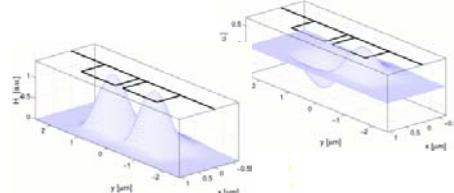
- localized reflection

VCGLE: O(2) group symmetry breaking

Linear pol. basis
0 or π phase diff



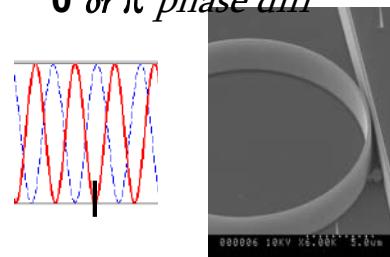
Supermodes
0 or π phase diff



O. Hess et al. Phys. Rev. A 50, 787 (1994)

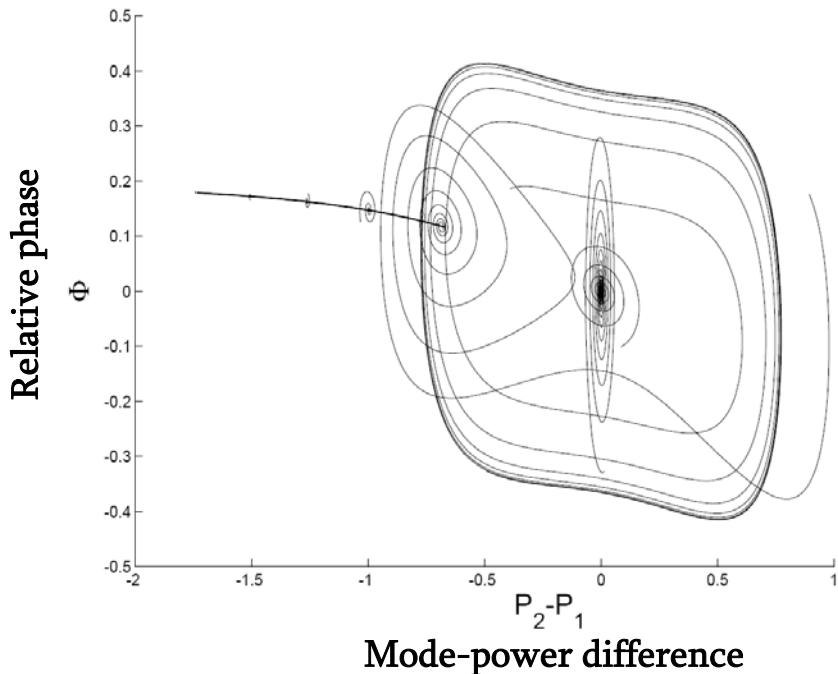
Different effective ref. index: Different Frequency and losses

Standing waves
0 or π phase diff



R.J.C. Spreew et al. PRA 45 1213 (1992)

VCGLE. Bifurcations. $\gamma = 2$. $k_d, k_c \sim 10^3, 10^4$

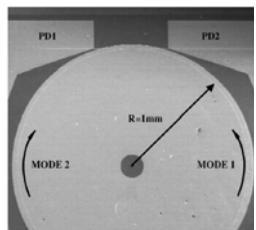


VCGLS \equiv experiments on AlGaAlAs 1mm radius Ring Laser

Why?

- $f_{Hopf} \sim 10-100 \text{ MHz} \ll f_{Rel.Ox}$

Experimental Results. SRL modeled by TwoMode. SVA Eqs.



2mm \varnothing , DQW GaAs-AlGaAs

Three regimes

A: bi-CW

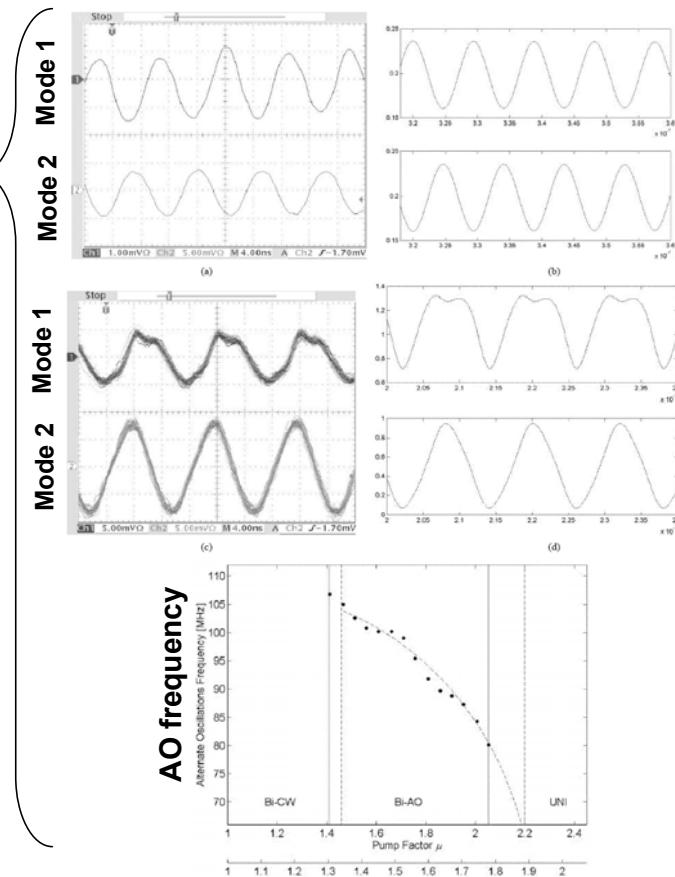
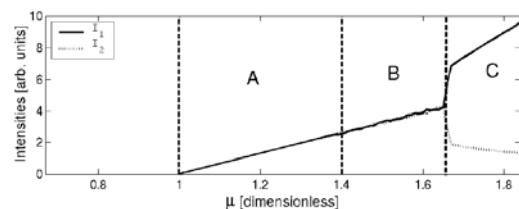
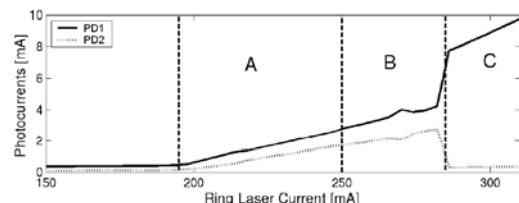
Symmetric fixed point

B: bi-AO

Hopf induced limit cycle

C: UNI

Pitchfork induced bistability



Sorel et al, IEEE JQE, 10 1187 (2003), Opt.Lett. 27 1992 (2002)

Switching in bistable SRL

Two-mode SVA model results

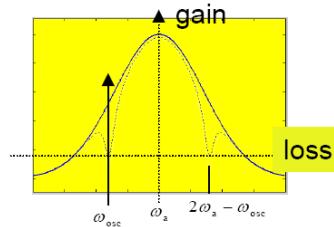
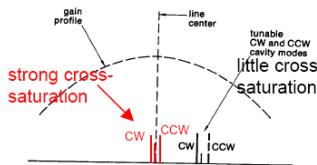
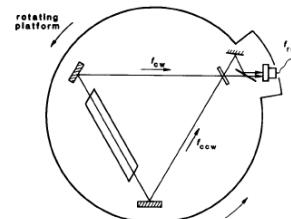
- Single mode SRL with perturbative backscattering is intrinsically bistable
- Structural constraints to achieve bistability
- Some insight on the switching process: energy redistribution between CW and CCW
- Switching time $10\text{-}30\text{ ps}$. Switching energy *few Joules*

..... *But*

- (⌚ Model not suitable @ ps timescale
- (⌚ Model not suitable for microRings
- ➔ TW approach (local SHB due to pulse propagation)

He-Ne Ring-Laser-Gyro

- Ring lasers used as rotation sensors
 - CCW and CW mode frequencies split
 - $\omega_{CCW} - \omega_{CW} \propto$ rotation rate
 - beat note measures rotation
- Obviously need both modes oscillating simultaneously
 - usually OK as gain medium (He-Ne) is inhomogeneously broadened
 - spectral holes usually don't overlap
 - cross-saturation weak: $C < 1$



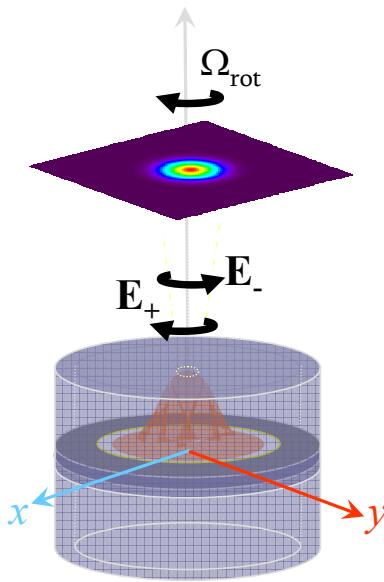
Phase-locking (Adler eq.) $d\Delta\phi/dt = \Delta\omega + k_d \sin(\Delta\phi)$

General coupled oscillators behavior $\Delta\phi = \arcsin(\Delta\omega/k_d)$

Fokker-Plank description of phase noise (NER)

[W.W.Chow et al. Reviews of Modern Physics 57 61, 1985.]

Spinning-VCSEL



Inertial rotation induces circular anisotropy In VCSELs

Theoretical approach

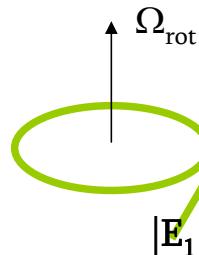
- Maxwell Eqs in rotating framework (*Lorentz tr.*)
+ Bloch Eqs. (medium at rest) see discussion
[Phys. Rev. 1966-1969]

Modified Spin-Flip model + circular birefringence $\gamma_c \sim \Omega_{\text{rot}}$
• Indirect effect on polarization through γ_p :

phase-ellipticity coupling

$$\tan(2\chi) = -\frac{\gamma_c}{\gamma_p} \left(1 + \frac{\alpha k(\mu-1)\gamma}{\gamma_s \gamma_p} \right)$$

Rotating SRL



$$\Delta = \frac{2\pi R \tau_p}{\lambda} \Omega_{rot}, \quad \begin{cases} \text{Inertial rotation term} \\ \text{In two-mode SVA eq} \end{cases}$$

$$\dot{E}_{1,2} = \pm i\Delta E_{1,2} + \dots - (k_d + ik_c) E_{2,1}$$

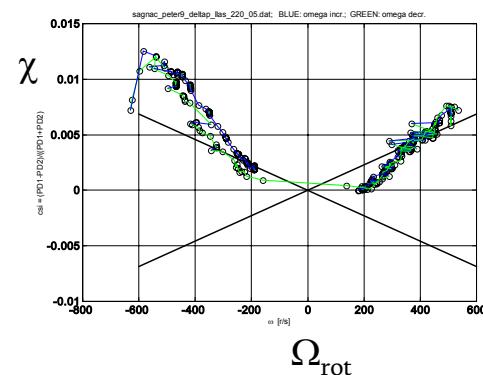
$$E_{1,2} = Q(1 \pm \delta) \exp(i\omega + i\phi_{1,2}) \quad \chi = \frac{|E_1|^2 - |E_2|^2}{|E_1|^2 + |E_2|^2} \approx 2\delta \quad \phi = \phi_2 - \phi_1$$

$$\begin{aligned} \dot{\delta} &= (\eta Q^2 N_0 - k_d)\delta - k_c \phi + \sqrt{\beta_{sp} \tau_p N_0} W_\delta \\ \dot{\phi} &= (\alpha \eta Q^2 N_0 - k_c)\delta + k_d \phi + \Delta + \frac{\sqrt{\beta_{sp} \tau_p N_0}}{Q} W_\phi \end{aligned}$$

↓

$$\delta = \frac{-k_c}{[(k_d + \alpha k_c)2\eta Q^2 N_0 - k_d^2 - k_c^2]} \times \Delta.$$

CW measure of inertial rotation

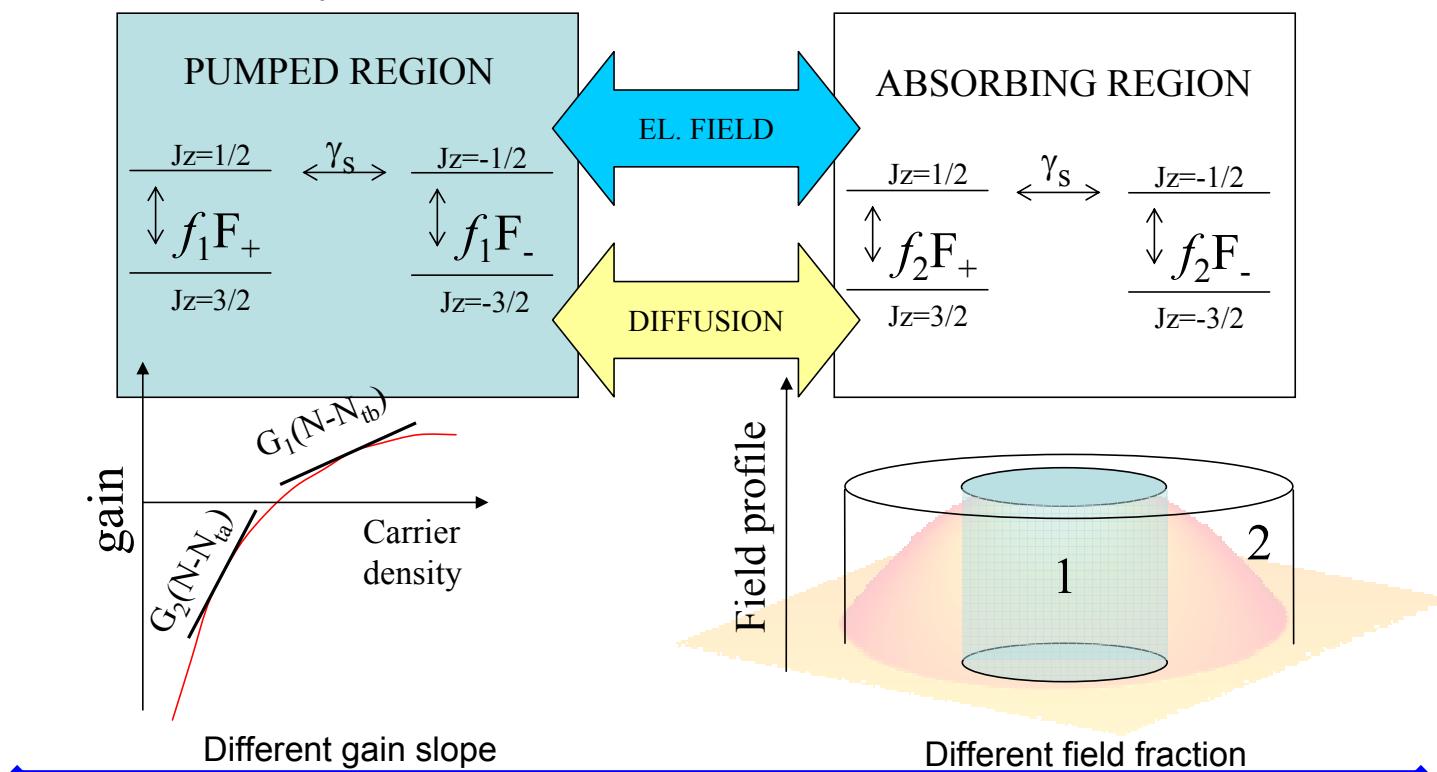


Exp. by G.Giuliani, Università di Pavia
GaAs-AlGaAs SRL

Two-mode models + saturable absorber: small radius VCSEL

Spin-Flip + Yamada Model [M. Yamada IEEE JQE 29 1330 1993]

- Self-pulsations with polarization d. of f. [Scirè et al. Opt. Lett. 27 391, 2002]
- Polarization chaos: application to encoded communications [Scirè et al. PRL 90, 113901, 2003]
- Vectorial excitability (?)



Two-mode models + saturable absorber Twin-Stripes

Two.modeSVA + Yamada Model

- Synchronization of SP [Scirè et al. JQE 41 272, 2005] in TS+sat.abs.
- View of synchronizing a many element SP_EELs array
- View of Synchronized array of mode-locked EELs**

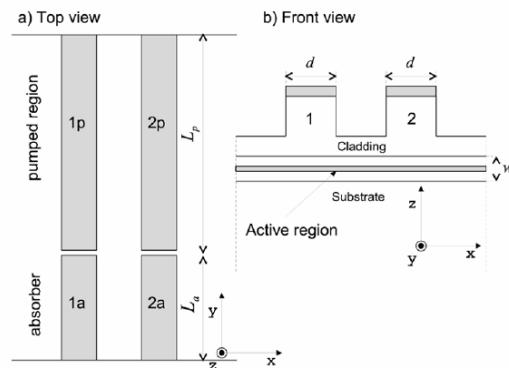
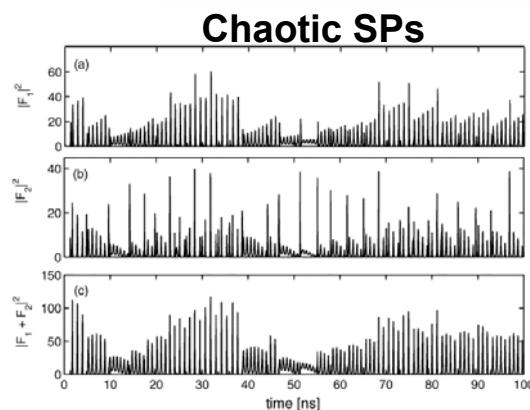
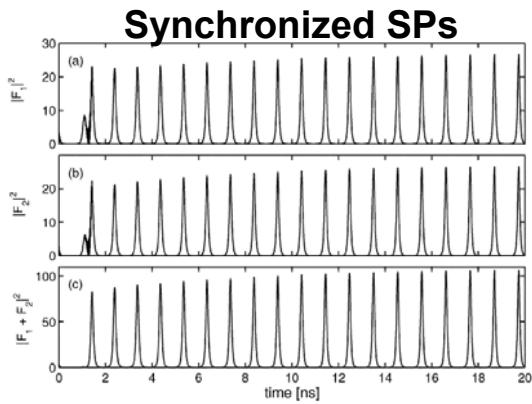


Fig. 1. Schematic picture of the device.



⌚ Linearized gain Vs large carrier excursions ☹ Fast dynamics

- Arrays

Early works on arrays as a phase-oscillators-synchronization problem

[H.G.Winful, ch.5 *Diode Laser Arrays*, Cambridge 1994]

Theory:

- Rate equations with first neighbour coupling

Perfect synchronization – Kuramoto Model

- From network theory

All-to-all coupling, $k_c = 1/\sigma$, av. path length $\ell=1$

Given topology, $k_c = f(\ell)/\sigma$. E.g. linear chain of osc. $\ell = \ln(N)$

Coupled – mode theory is perturbative

e.g. [O.Hess PRA 50, 787 1994] – three regimes dep. on interelement spacing

Array synchronization is a size dependent problem

Small arrays: coupled mode theory.

Big arrays: Bloch functions [D.Botez Diode laser arrays, Cambridge 1994].

Int. Case: device oriented modeling [P.Debernardi]

Peculiarities in synchronizing pulsating lasers

- **Conclusions**



- **Acknowledgments**
- **Salvador Balle, Guido Giuliani, Jan Danckaert, Siyuan Yu, Sandor Fuster, Pep Mulet *and his nice figures*, Francesco Marino, Francesco d'Ovidio, Iacyel Gomez da Silva, Victor M. Eguiluz, and many others....**



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