

Quantum correlations of dissipative oscillators: Driving and synchronization effects

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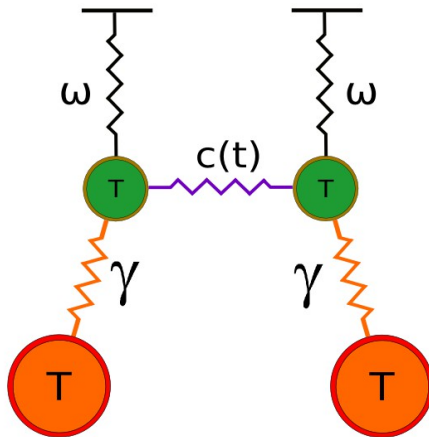


2 dissipative driven harmonic oscillators

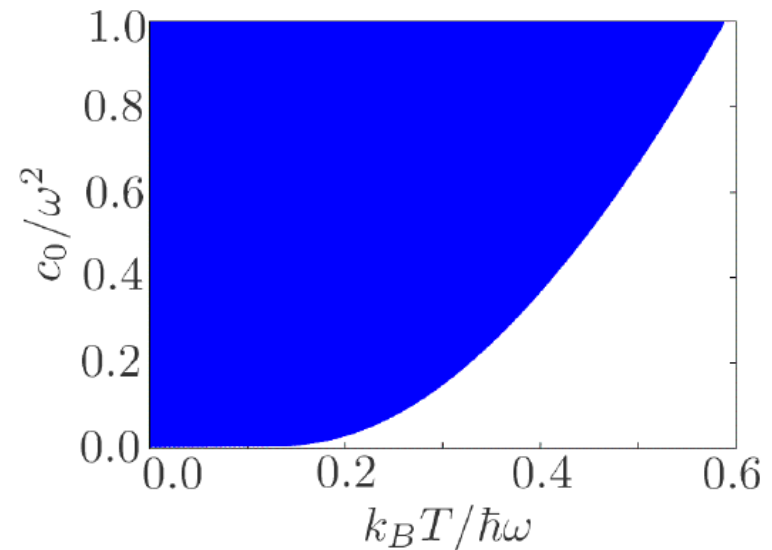
$$\mathcal{H}_s = \sum_{\alpha=1}^2 \left(\frac{P_{\alpha}^2}{2m} + \frac{1}{2} m \omega^2 Q_{\alpha}^2 \right) + c(t) Q_1 Q_2 \quad c(t) = mc_0 + mc_1 \cos(\omega_d t)$$

$$\mathcal{H} = \mathcal{H}_s + \sum_{\alpha,k=1}^{2,\infty} \frac{p_{\alpha,k}^2}{2m_k} + \frac{m_k \omega_k^2}{2} \left(x_{\alpha,k} - \frac{c_k Q_{\alpha}}{m_k \omega_k^2} \right)^2$$

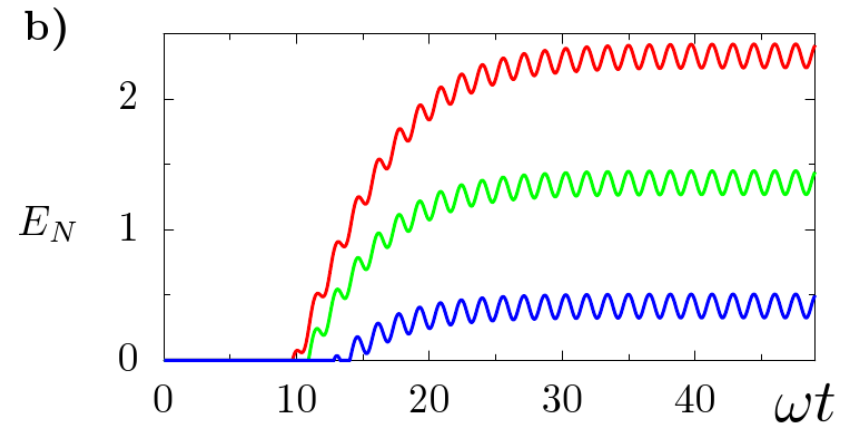
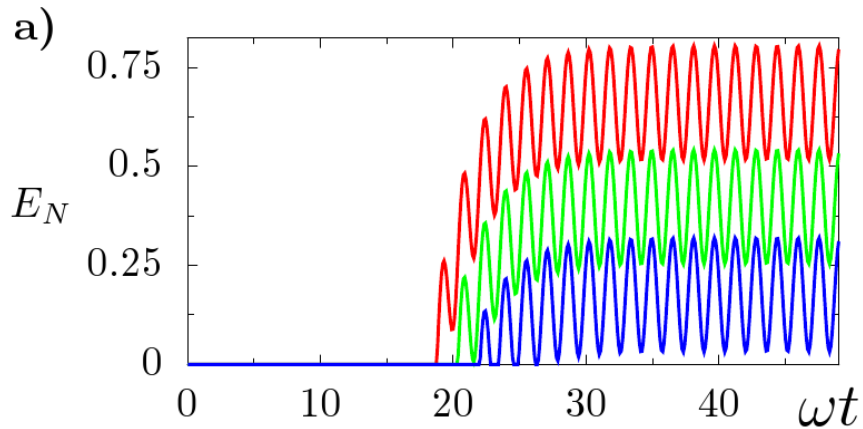
– We choose separate baths –



– Thermal entanglement –



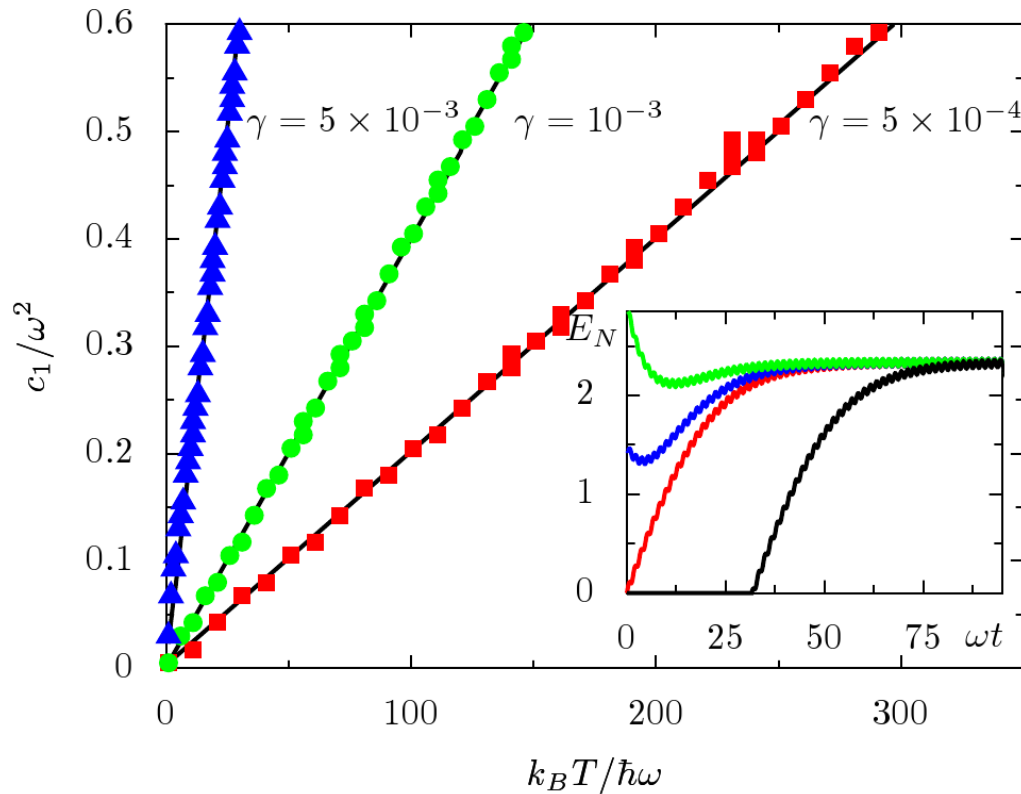
Driving produces nonequilibrium stationary entanglement



a) $\gamma = 0.005\omega$, $c_1 = 0.5m\omega^2$,
 $k_B T / \hbar\omega = 250$ (red), 300 (green), 350 (blue)

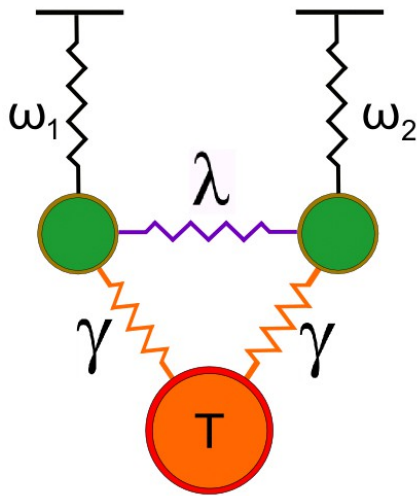
b) $c_1 = 0.5m\omega^2$, $k_B T / \hbar\omega = 5$,
 $\gamma = 0.005\omega$ (red), 0.01ω (green), 0.02ω (blue)

Which translates into a high-T entanglement diagram



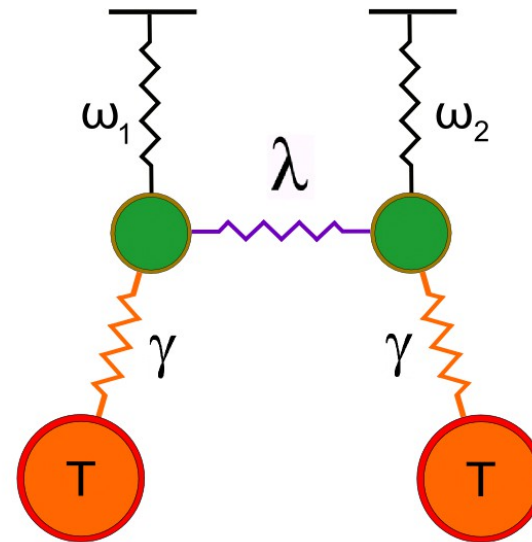
$$\frac{k_B T}{\hbar\omega} \leq \frac{|\text{Im}(\mu_M)|}{\gamma}$$

2 nonidentical dissipative harmonic oscillators (synchronization)



Common bath

(oscs. variances synchronize)

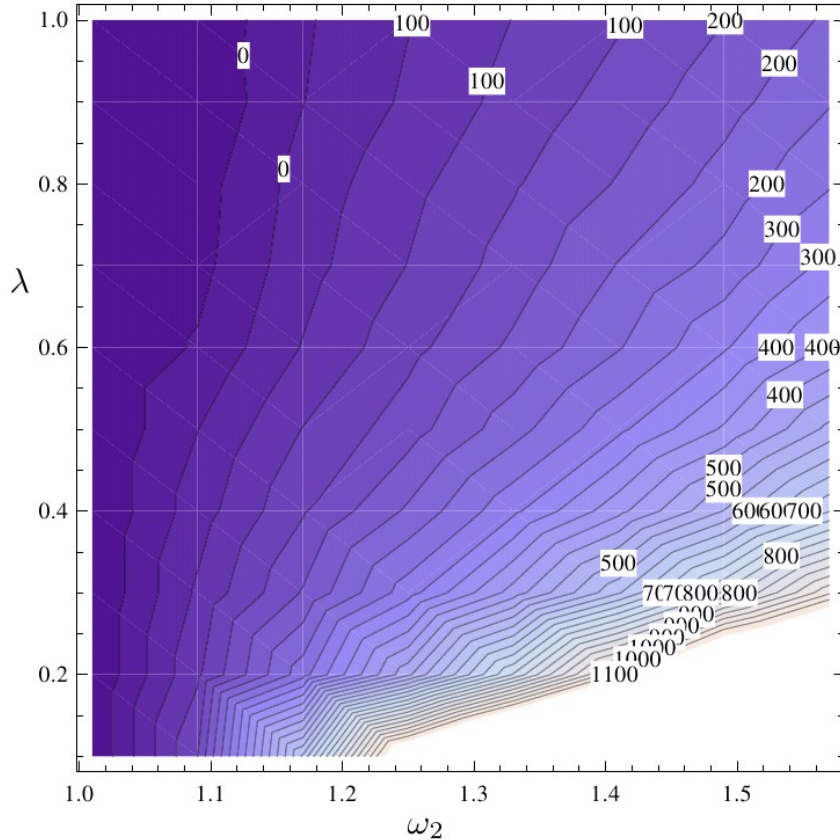


Separate baths

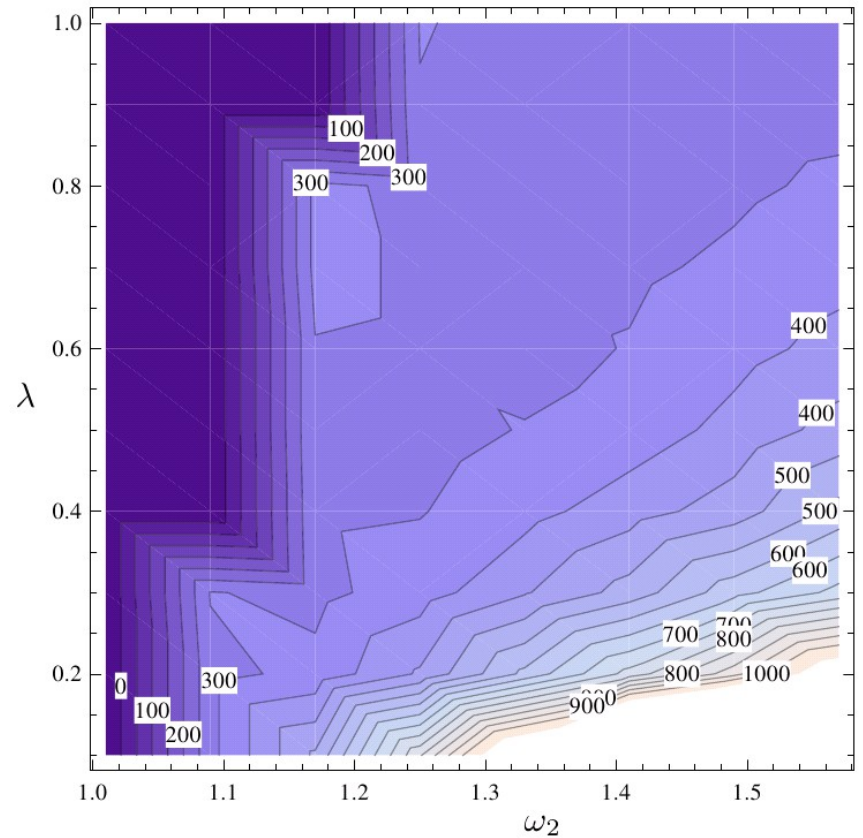
(oscs. var. do not synchronize)

Synchronization of first and second moments shows similar behaviour (common bath):

2nd moments sync. time



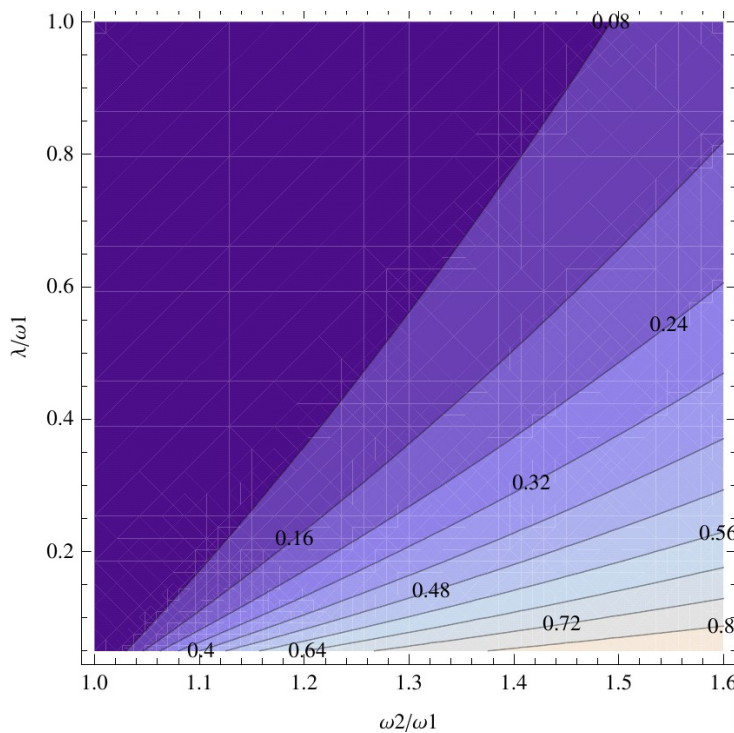
1st moments sync. Time



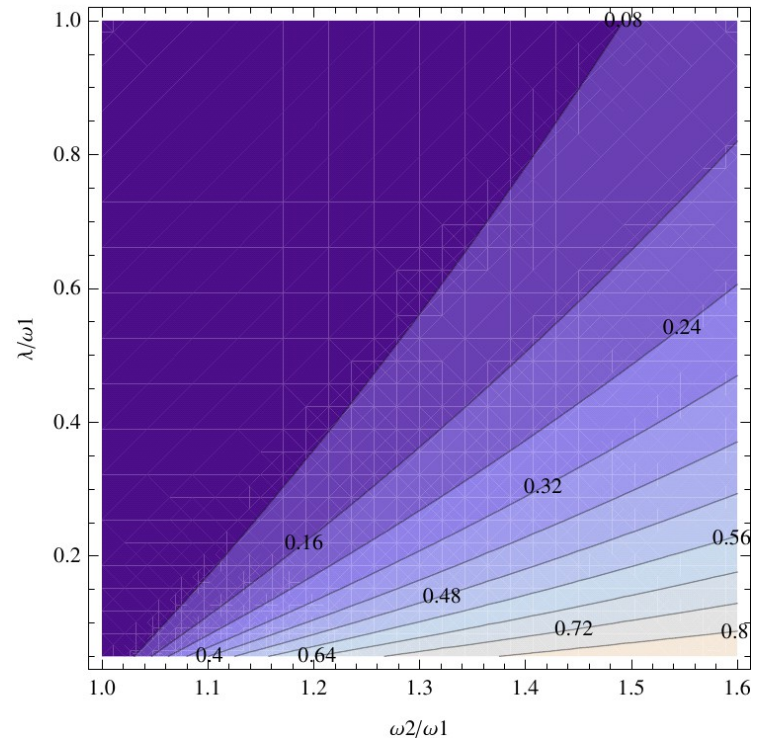
Synchronization times can be understood from the real parts of the eigenvalues of the dynamical equations (equations for all variances).

-- Smallest real part over largest one --

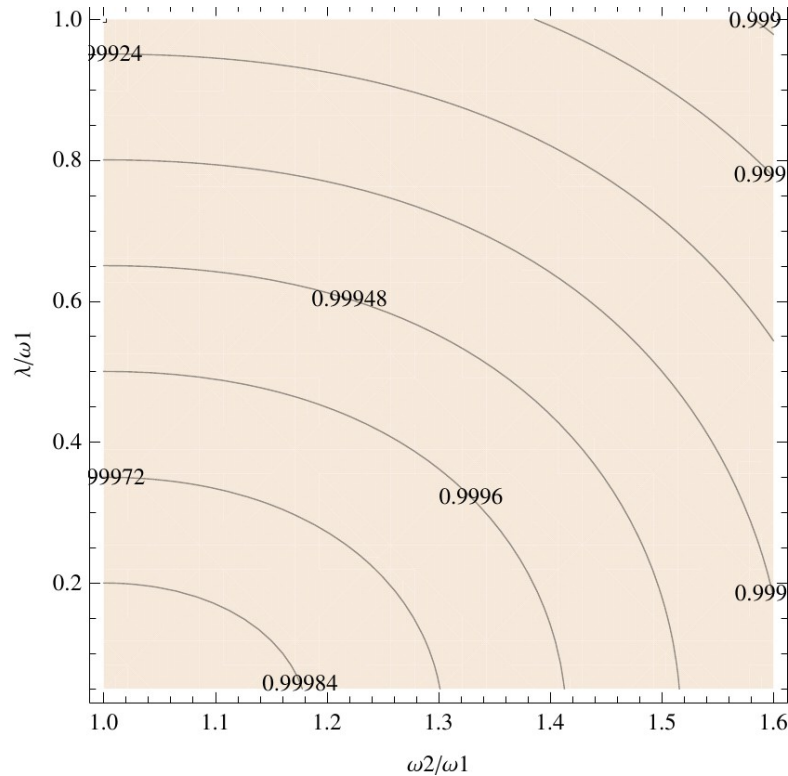
(2nd moments, common bath)



(1st moments, common bath)



-- Smallest real part over largest one --
 (2nd moments, Separate baths)



No “mode” relaxes faster than others, so NO synchronization

Synchronization of 2nd momenta leads to reduced information loss, both quantum correlations and total correlations (initial state: 2-mode squeezed):

Common bath

Separate baths

