

Coupling between two quantum dots through a superconducting island

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Jožef Stefan Institute
Ljubljana, Slovenia



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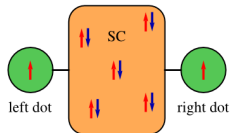
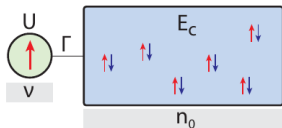
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Motivation

Double quantum dot + superconducting island

- Potential qubit realization by using two dots coupled to a superconducting island SC (subgap states)

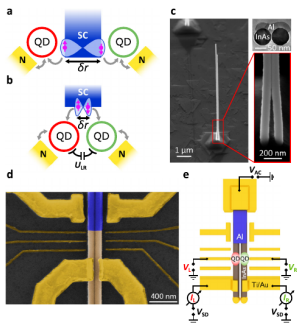


Source: PRB 104, L241409 (2021)

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Source: npj Quantum Materials, 7, 88 (2022)

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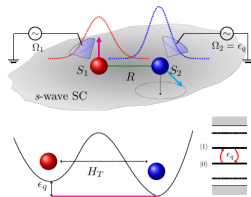
Double quantum dot + superconducting island

- Potential qubit realization by using two dots coupled to a superconducting island SC (subgap states)
- Cooper pair splitter
- Interaction between magnetic impurities embedded in superconducting materials

PRL 113, 087202 (2014)

PRX Quantum 2, 040347 (2021)

⇒



Interaction between magnetic impurities

- RKKY interaction:

$$H_{ij} = J(\mathbf{R}_1 - \mathbf{R}_2)\mathbf{S}_1\mathbf{S}_2$$

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- Between magnetic impurities located closely to each other, the coupling is described by superexchange

Two quantum dots coupled to SC

- Model:

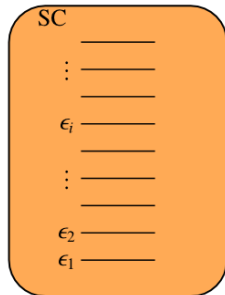
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- Model:

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$$H_{SC} = \sum_{\sigma, i=1}^N \epsilon_i c_{i\sigma}^{\dagger} c_{i\sigma} - \frac{g}{N} \sum_{i,j} c_{i\uparrow}^{\dagger} c_{i\downarrow}^{\dagger} c_{j\downarrow} c_{j\uparrow} \quad (\text{Richardson model})$$



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$$H_{QDs} = \varepsilon_L \sum_{\sigma} d_{L\sigma}^{\dagger} d_{L\sigma} + \varepsilon_R \sum_{\sigma} d_{R\sigma}^{\dagger} d_{R\sigma} + U_L n_{L\uparrow} n_{L\downarrow} + U_R n_{R\uparrow} n_{R\downarrow}$$

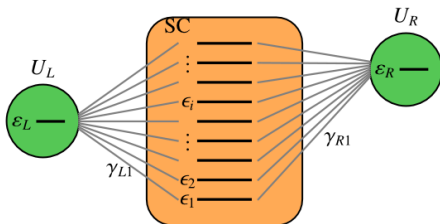


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- Model:

$$H = H_{SC} + H_{QDs} + H_{hyb}$$

$$H_{hyb} = v_L \sum_{\sigma i} (\gamma_{Li}^{\sigma} c_{i\sigma}^{\dagger} d_{L\sigma} + h.c.) + v_R \sum_{\sigma i} (\gamma_{Ri}^{\sigma} c_{i\sigma}^{\dagger} d_{R\sigma} + h.c.)$$



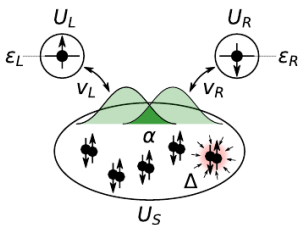
- Strength of couplings: v_L and v_R
- Distribution of couplings: γ_{Li} and γ_{Ri} which fulfill $\sum_i |\gamma_{Li}|^2 = \sum_i |\gamma_{Ri}|^2 = 1$

Two quantum dots coupled to SC

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- Hybridizations may also overlap: $\alpha = \sum_i \gamma_{Li}^* \gamma_{Ri}$ $\alpha \in [0; 1]$



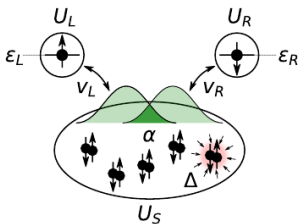
Source: arXiv:2303.14410

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- We will focus on *close* QDs, $\alpha = 1$.
- Main question: What is the ground state depending on $\epsilon_{L/R}$, $v_{L/R}$ and U ?
 What spin configuration?

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- In the single QD - SC system, flat-band limit provided a good description
- Ground state of the superconductor with fixed number of Cooper-pairs M

$$|\Psi_M^N\rangle \sim \left(\sum_i c_{i\uparrow}^{\dagger} c_{i\downarrow}^{\dagger} \right)^M |0\rangle$$

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$$H_{SC} \approx \frac{N|\Delta|^2}{g} - \sum_i \left(\Delta c_{i\uparrow}^+ c_{i\downarrow}^+ + h.c. \right)$$

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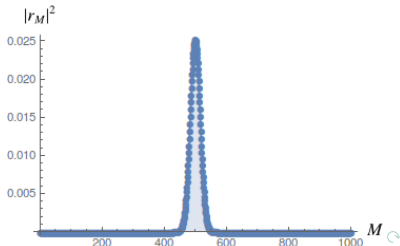
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- Diagonalization with Bogoliubov (finite chemical potential)
- Ground state is a mixture of $|\Psi_M^N\rangle$ with different values of M

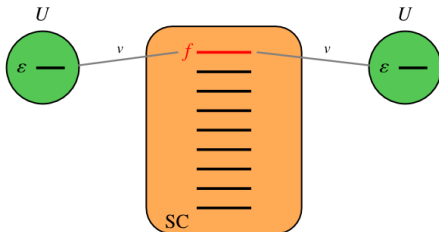
$$|\Psi_{BCS}\rangle = \sum_M r_M |\Psi_M^N\rangle$$

r_M describes a binomial distribution



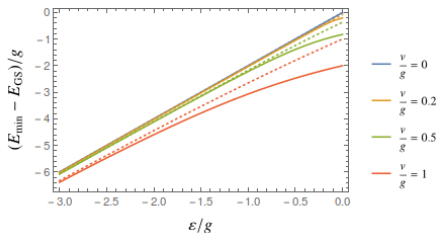
Complete overlap, $\alpha = 1$, analytical results

- Fix the number of Cooper pairs, M , and assume the limit $U \rightarrow \infty$ and the thermodynamic limit
- Symmetric coupling v and $\varepsilon < 0$
- Quantum numbers:
 - total number of particles (even/odd)
 - total spin ($S = 0$ singlet, $S = 1/2$ doublet, $S = 1$ triplet)
 - total S_z
 - parity (symmetric/antisymmetric)



Complete overlap, $\alpha = 1$, analytical results

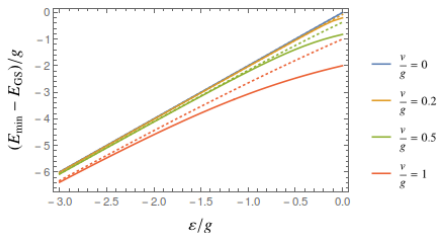
- Even number of particles (singlet/symmetric or triplet/antisymmetric)



—	singlet	$\frac{1}{\sqrt{2}} \left(\uparrow_L, \downarrow_R\rangle - \downarrow_L, \uparrow_R\rangle \right) \otimes \Psi_M^N\rangle_{SC}$
---	triplet	$ \uparrow_L, \uparrow_R\rangle \otimes \Psi_M^N\rangle_{SC}$

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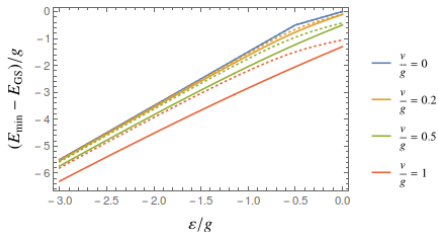


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- The difference in minimal energies: $E_T - E_S = \frac{2v^4}{|\epsilon|^3}$

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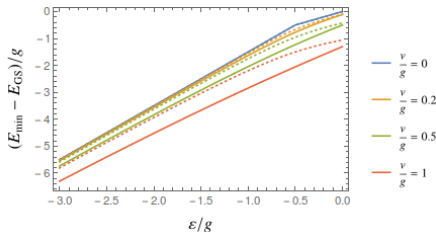
- Odd number of particles, doublet (symmetric or antisymmetric)



———	antisymmetric	$\frac{1}{\sqrt{6}} (2 \uparrow, \uparrow\rangle \otimes \downarrow, \Psi_{M-1}^{N-1}\rangle - (\uparrow, \downarrow\rangle + \downarrow, \uparrow\rangle) \otimes \uparrow, \Psi_{M-1}^{N-1}\rangle)$
----	symmetric	$\frac{1}{\sqrt{2}} (\uparrow, \downarrow\rangle - \downarrow, \uparrow\rangle) \otimes \uparrow, \Psi_{M-1}^{N-1}\rangle$

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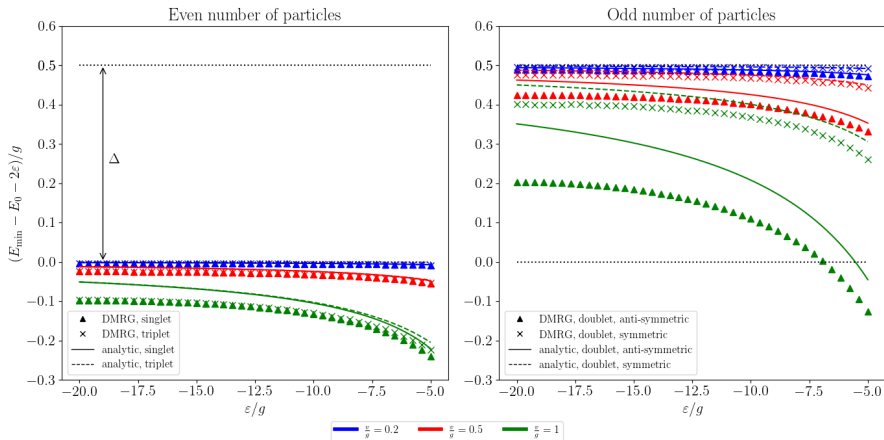


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----	symmetric	$\frac{1}{\sqrt{2}} (\uparrow, \downarrow\rangle - \downarrow, \uparrow\rangle) \otimes \uparrow, \Psi_{M-1}^{N-1}\rangle$

- The difference in minimal energies: $E_{D,sym} - E_{D,asym} = \frac{2v^2}{|\epsilon|}$

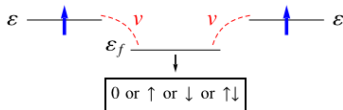
Comparison with DMRG

- Density Matrix Renormalization Group simulation with $U/g = 40$ and $N = 80$



Superexchange through a single level

- Two dots coupled to a single level

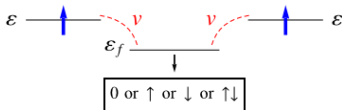


- Hamiltonian

$$H = \varepsilon \sum_{\sigma \delta=L,R} d_{\delta\sigma}^{\dagger} d_{\delta\sigma} + U \sum_{\delta=L,R} n_{\delta\uparrow} n_{\delta\downarrow} + v \sum_{\sigma \delta=L,R} (d_{\delta\sigma}^{\dagger} f_{\sigma} + h.c.) + \varepsilon_f \sum_{\sigma} f_{\sigma}^{\dagger} f_{\sigma}$$

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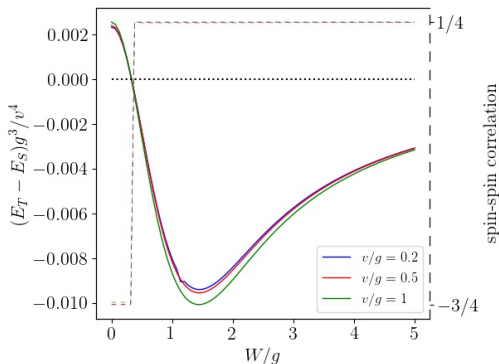
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- The occupation of the mediating level determines whether ferromagnetic or antiferromagnetic alignment is favoured
- Doubly occupied or empty intermediate level favours antiferromagnetic alignment
- Singly occupied intermediate level favours antiferromagnetic alignment

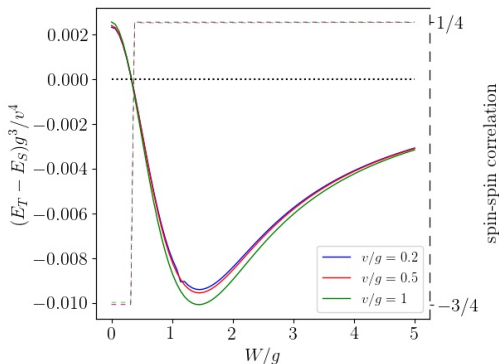
Effects of finite bandwidth

- DMRG results:
 - even number of particles: transition from singlet to triplet
 - odd number of particles: no transition, always ferromagnetic alignment



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- Kinetic energy enhances the probability of single occupancy on the distinguished level

Summary

- In the flatband limit, QD-SC-QD system features
 - singlet (antiferromagnetic alignment) for even number of particles
 - parity-antisymmetric doublet (ferromagnetic alignment) for odd number of particles
- The features can be understood through the occupancy of the distinguished level of the superconductor
- Finite bandwidth plays an important role as opposed to single QD - SC

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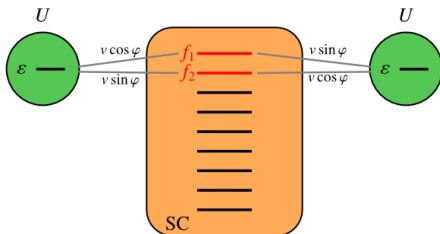
Thank you for your attention!

Löwdin orthogonalization

- Two quantum dots coupled to the SC in a symmetric way: $\varepsilon_L = \varepsilon_R = \varepsilon$, $v_L = v_R = v$ and $U_L = U_R = U$
- In the flatband limit, the SC levels are identical

⇒ freedom to choose basis of SC levels $f_{i\uparrow} = \sum_{j=1}^N U_{ij} c_{i\uparrow}$ $f_{i\downarrow} = \sum_{j=1}^N U_{ij}^* c_{i\downarrow}$

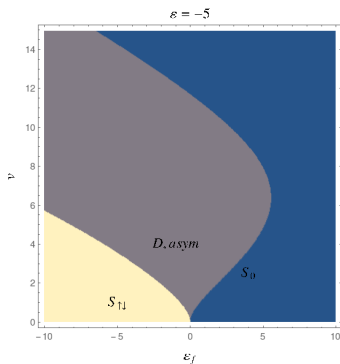
- We choose a basis in which the dots are coupled to two SC level only



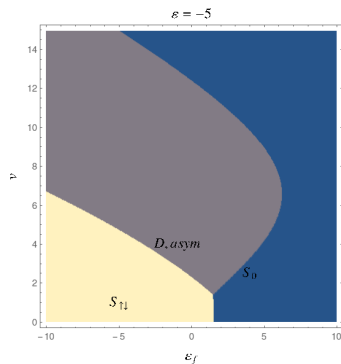
$$\alpha = \sin(2\varphi)$$

Superexchange through a single level

- Phase diagram for the ground state for $U = 10$ and $\varepsilon = -5$



$g = 0$



$g = 1$