

Large-scale transport in oceans

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Statistical Physics and Dynamical Systems approaches in Lagrangian Fluid Dynamics



STATISTICAL PHYSICS AND DYNAMICAL SYSTEMS APPROACHES IN LAGRANGIAN FLUID DYNAMICS

OUTLINE

1. Lagrangian fluid dynamics and introduction to chaotic advection. Hamiltonian dynamics, KAM tori, Lyapunov exponents, open flows
2. **Dispersion, diffusion and coherent structures in flows.** Turbulent, pair and chaotic dispersion, gradient production, FTLE, FSLE, Lagrangian Coherent Structures
3. Chemical and biological processes in flows. Fisher and excitable plankton waves, filamental transitions, lamellar approaches, burning manifolds
4. Complex networks of fluid transport. Directed and weighted flow networks. Community detection

Diffusion equation: $\frac{\partial C}{\partial t} = D\nabla^2 C.$

$$C(\mathbf{x}, t) = \int d\mathbf{x}' G(\mathbf{x} - \mathbf{x}', t) C_0(\mathbf{x}') , \quad G(\mathbf{x}, t) = \frac{1}{(4\pi Dt)^{d/2}} e^{-\frac{\mathbf{x}^2}{4Dt}}.$$

$$w^2 \equiv \frac{\int \mathbf{x}^2 G(\mathbf{x}, t) d\mathbf{x}}{\int G(\mathbf{x}, t) d\mathbf{x}}, \quad w \approx (2dDt)^{1/2} ,$$

Turbulent diffusion- diffusion-like behavior by advection:

$$\frac{d}{dt} \langle (\mathbf{r} - \mathbf{r}_0)^2 \rangle = 2 \langle (\mathbf{r} - \mathbf{r}_0) \cdot \mathbf{v}[\mathbf{r}(t)] \rangle.$$

$$\langle (\mathbf{r} - \mathbf{r}_0)^2 \rangle = 2 \int_0^t dt' \int_0^{t'} dt'' \langle \mathbf{v}[\mathbf{r}(t')] \cdot \mathbf{v}[\mathbf{r}(t' - t'')] \rangle$$

$$f_L(t', t'') \equiv \frac{1}{\langle v^2 \rangle} \langle \mathbf{v}[\mathbf{r}(t')] \cdot \mathbf{v}[\mathbf{r}(t' - t'')] \rangle \quad T_L \equiv \int_0^\infty d\tau f(\tau) . \quad \text{Lagrangian correlation time}$$

$$\langle (\mathbf{r} - \mathbf{r}_0)^2 \rangle \simeq \langle v^2 \rangle t^2 \quad t \ll T_L \quad \langle (\mathbf{r} - \mathbf{r}_0)^2 \rangle = \langle v^2 \rangle T_L t \quad t \gg T_L,$$

$$\langle (\mathbf{r} - \mathbf{r}_0)^2 \rangle = \langle v^2 \rangle T_L t$$

$$D_T \simeq \frac{1}{2d} \langle v^2 \rangle T_L.$$

Taylor dispersion

$$\frac{d}{dt} \langle (\mathbf{r}_2(t) - \mathbf{r}_1(t))^2 \rangle = 2 \langle (\mathbf{r}_2(t) - \mathbf{r}_1(t)) (\mathbf{v}(\mathbf{r}_2) - \mathbf{v}(\mathbf{r}_1)) \rangle$$

Small separations – similar to infinitesimal dispersion:

$$\langle (\mathbf{r}_2(t) - \mathbf{r}_1(t))^2 \rangle \approx \exp(\lambda t)$$

Large separations – similar to Taylor turbulent dispersion: $D=2D_T$

In between ... correlated dispersion.

For example, in the inertial range of 3d turbulence:

$$\frac{\langle [\mathbf{v}(\mathbf{r}_2) - \mathbf{v}(\mathbf{r}_1)] (\mathbf{r}_2 - \mathbf{r}_1) \rangle}{|\mathbf{r}_2 - \mathbf{r}_1|} = C \epsilon^{1/3} |\mathbf{r}_2 - \mathbf{r}_1|^{1/3} \quad \text{Kolmogorov scaling}$$

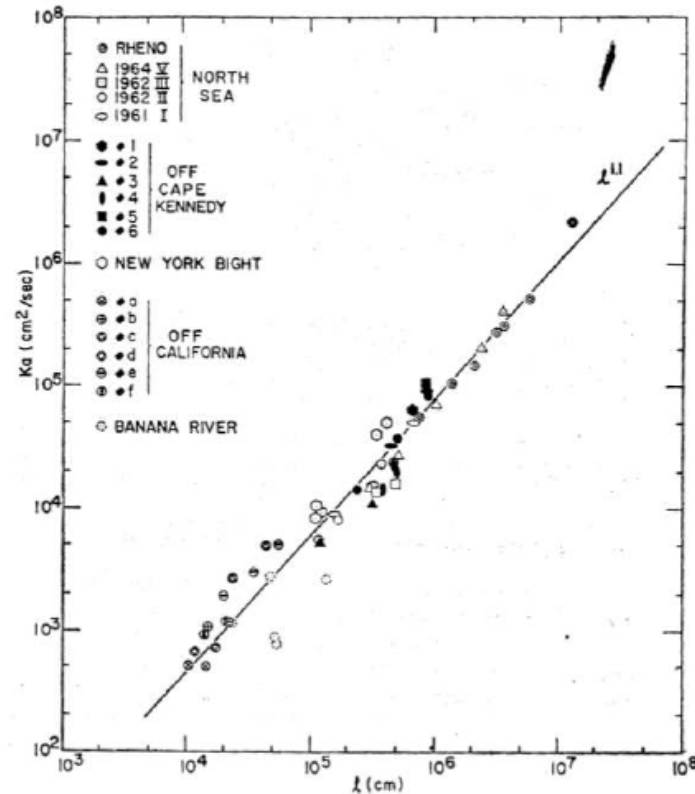
$$\frac{d}{dt} \langle (\mathbf{r}_2(t) - \mathbf{r}_1(t))^2 \rangle = 2C \epsilon^{1/3} |\mathbf{r}_2 - \mathbf{r}_1|^{4/3} \quad \langle |\mathbf{r}_2 - \mathbf{r}_1|^2 \rangle = \left(\frac{2}{3} C \right)^3 \epsilon t^3$$

Richardson law

This is somehow equivalent to a scale-dependent diffusivity: $D = C\epsilon^{1/3}|r_2 - r_1|^{4/3}$

Empirical effective (pair) diffusivity
Okubo, Dee Sea Res. 18, 789 (1971)

$$D_{eff}(l) \sim l^{1.15}$$

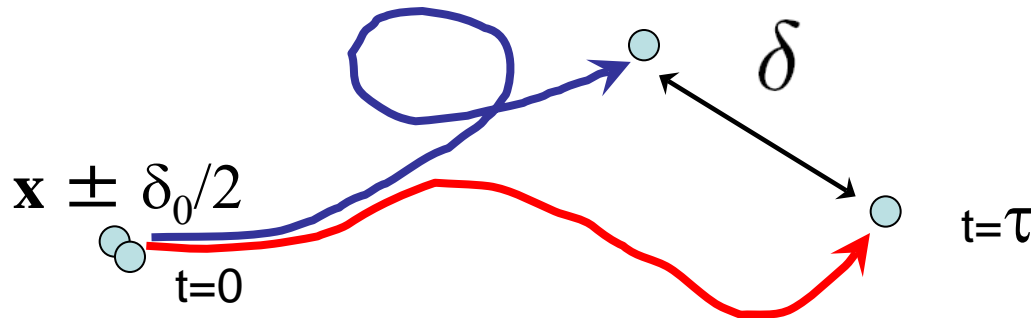


$$\lambda(t) = \lim_{\|\delta(0)\| \rightarrow 0} \frac{1}{t} \ln \frac{\|\delta(t)\|}{\|\delta(0)\|}$$

Finite-time Lyapunov exponent

$$\lambda = \lim_{t \rightarrow \infty} \lambda(t)$$

Lyapunov exponent



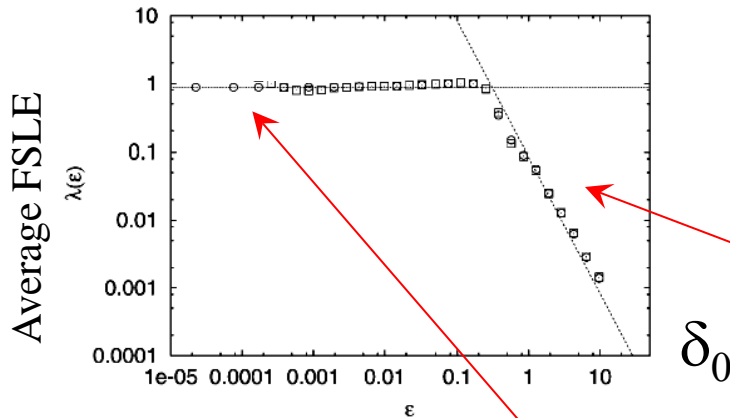
$$\lambda(\delta_0, \delta_f) \equiv \frac{1}{\tau} \log \frac{\delta_f}{\delta_0}$$

Finite-size Lyapunov exponent
FSLE

All the quantities are also functions of the initial position and time:

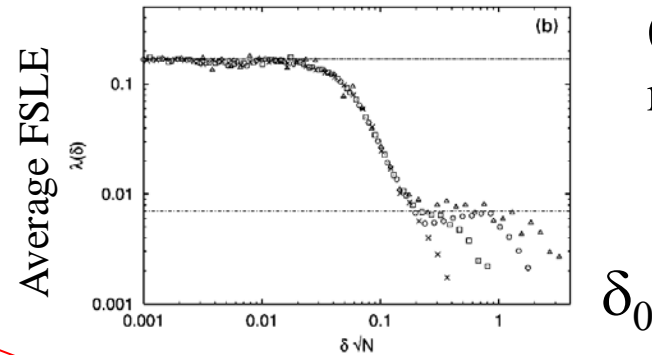
$$\lambda(\mathbf{x}, t, \delta_0, \delta_f)$$

A chaotic map



Exponential growth of separations (chaotic regime)

System with several time scales



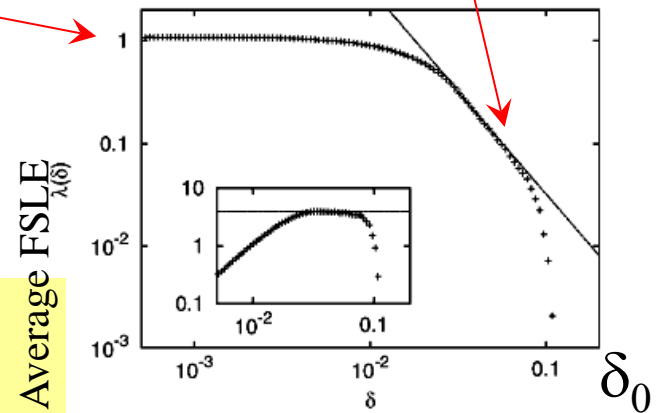
(coupled maps)

Subexponential growth (diffusion regime)

$$\lambda(\delta) \sim \delta^{-2}$$

When $\delta_0 \rightarrow 0$,
 FSLE \rightarrow Lyapunov
 and when $t \rightarrow \infty$,
 FTLE \rightarrow Lyapunov

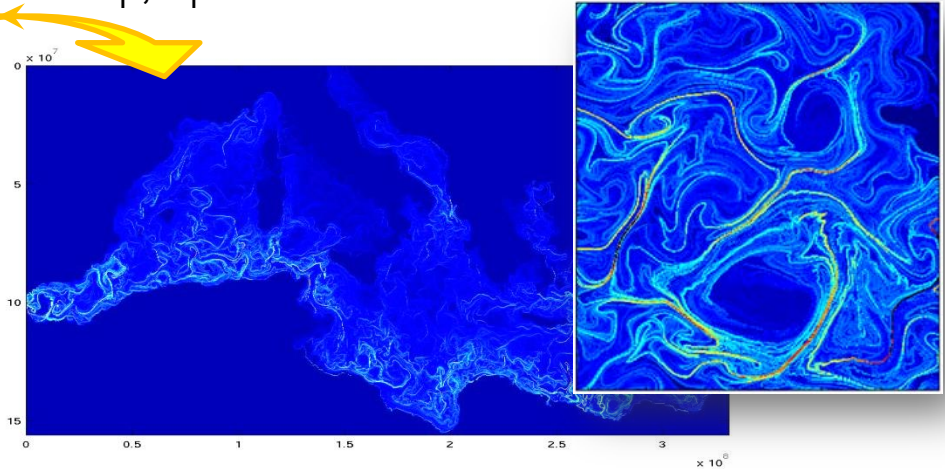
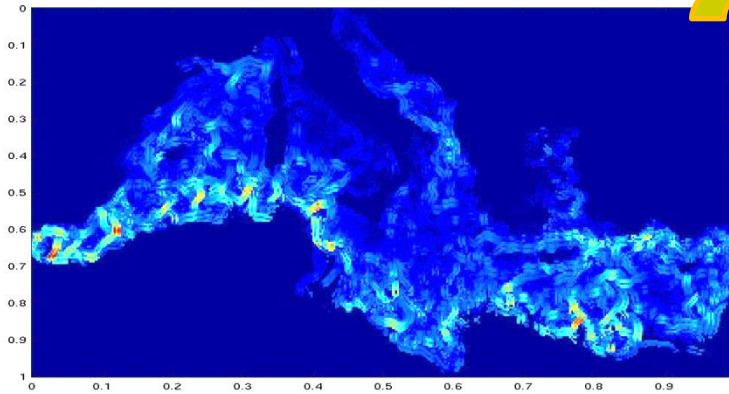
G. Boffetta et al. / Physics Reports 356 (2002) 367-474



2D turbulence

The FSLE was originally introduced to quantify dispersion from non-infinitesimal initial separations (Aurell et al. 1997)

Reducing scales δ_i, δ_f

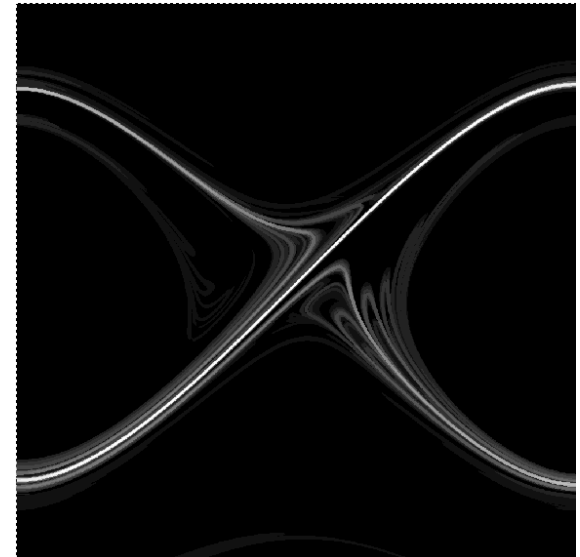
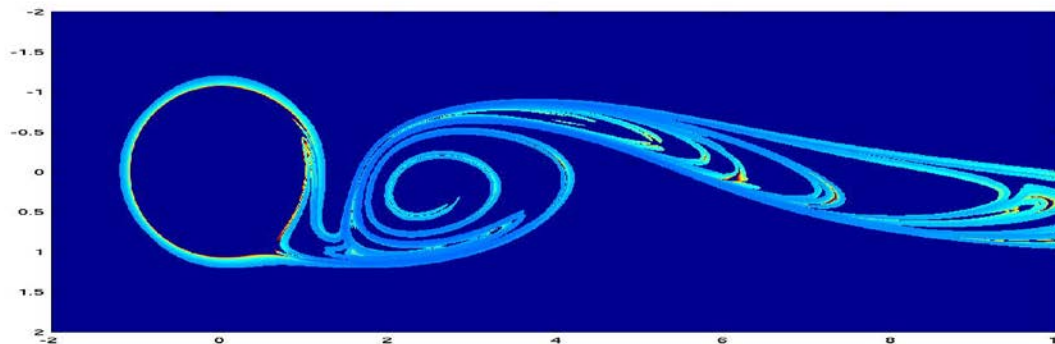


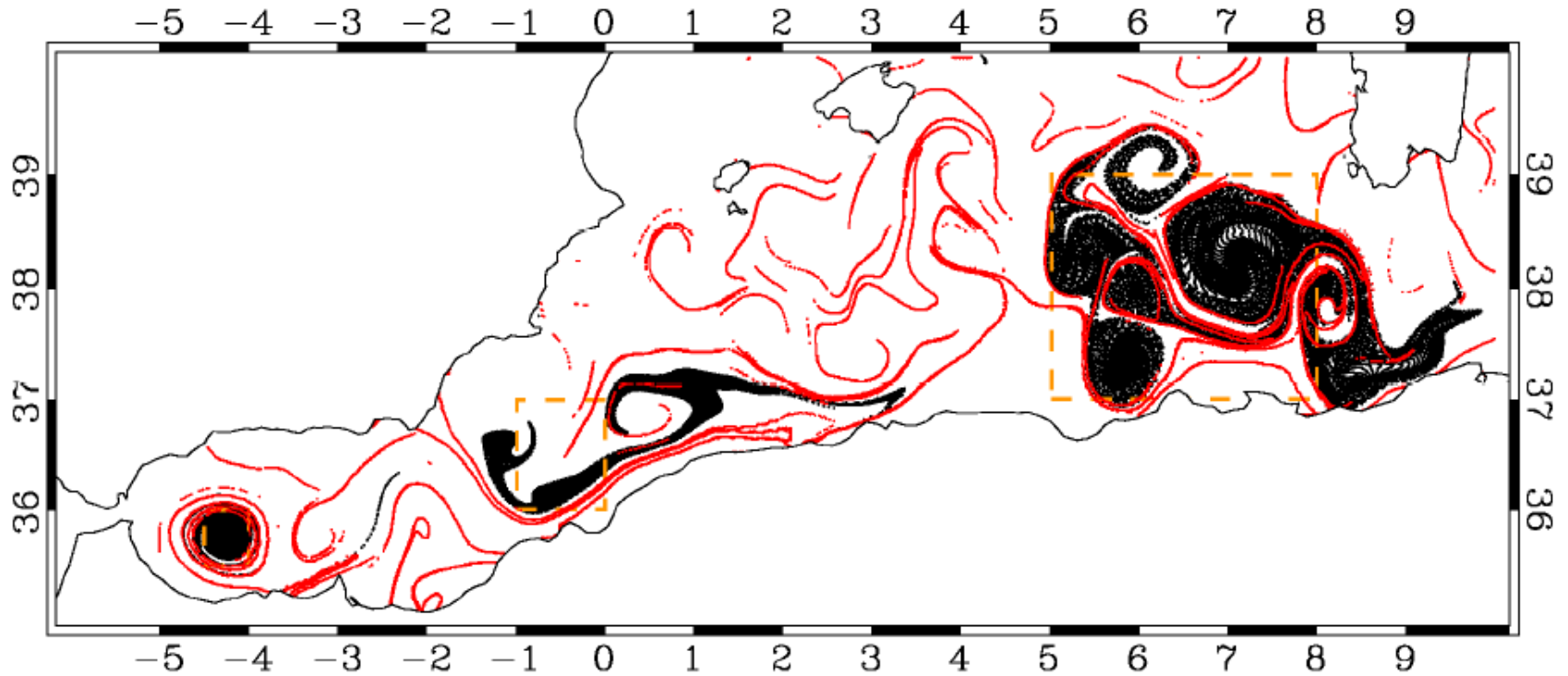
FSLE for small enough scales, \leftrightarrow **FTLE** for large enough times

Forward in time: repelling manifolds

Backward in time: attracting manifolds

LAGRANGIAN COHERENT STRUCTURES





FSLE are Lagrangian, but not direct advection:

- shorter simulations
- no problems with exponentially increasing line lengths
- exhaustive consideration of initial conditions

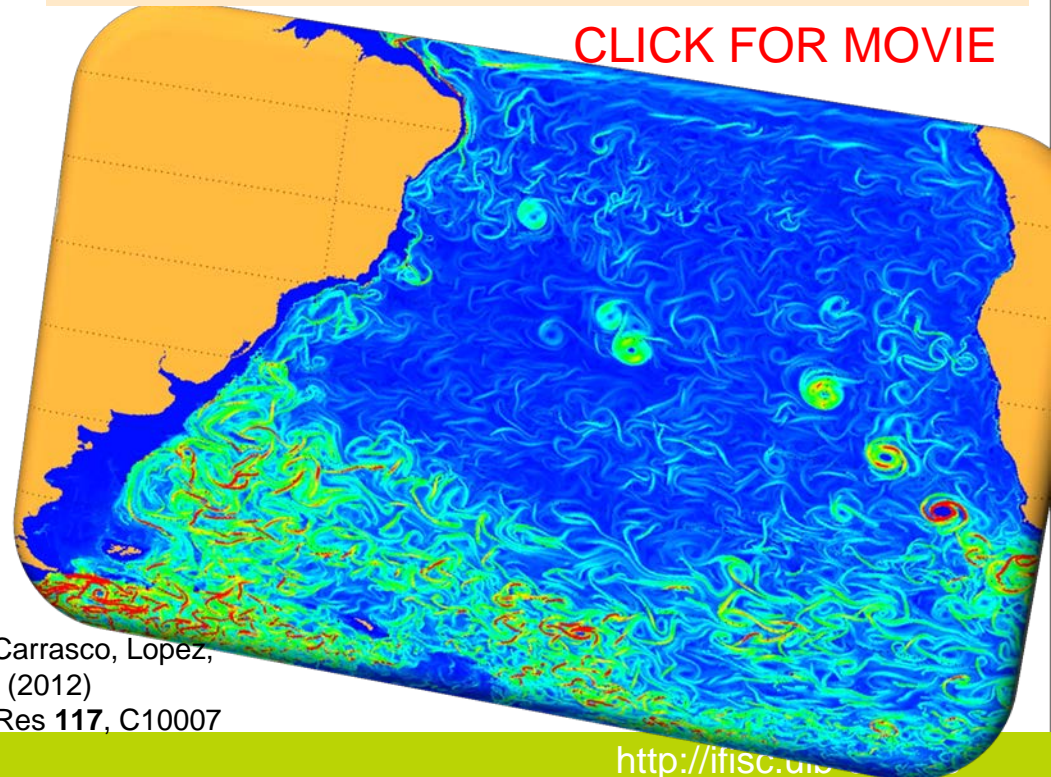
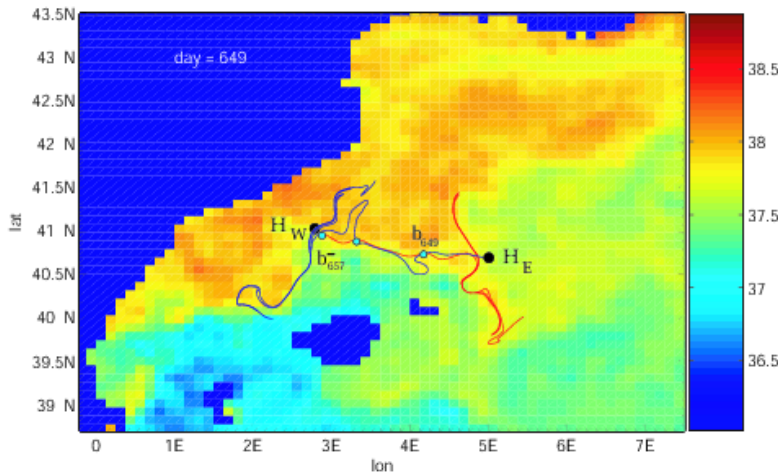
Lagrangian approaches to transport and mixing

- ❑ Geometric, local, ... : FTLE, FSLE, geodesics, variational theory, M function, ...
- ❑ Set-oriented, probabilistic ,...: Transfer operator, coherent sets, eigenvectors and singular vectors, networks, ...

- ❑ Detailed view of single events
- ❑ Statistical (climatological) descriptions

BIBLIOGRAPHY at 'Resources' for the School :
www.gefenol.es/school2014/resources/

CLICK FOR MOVIE



A.M. Mancho, E. Hernandez-Garcia, D. Small, S. Wiggins, V. Fernandez, J. Physical Oceanography **38**, 1222-1237 (2008).

Hernandez-Carrasco, Lopez, EHG, Turiel, (2012) J. Geophys Res **117**, C10007

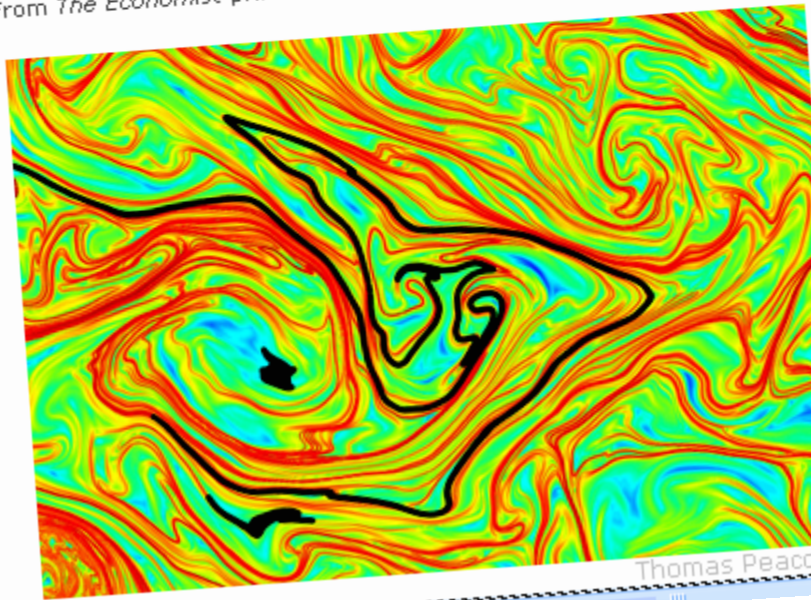


Lagrangian Coherent Structures

Lagrangian coherent structures The skeleton of water

Research is revealing a hidden structure within liquids and gases that guides the movement of everything from pollution to aeroplanes

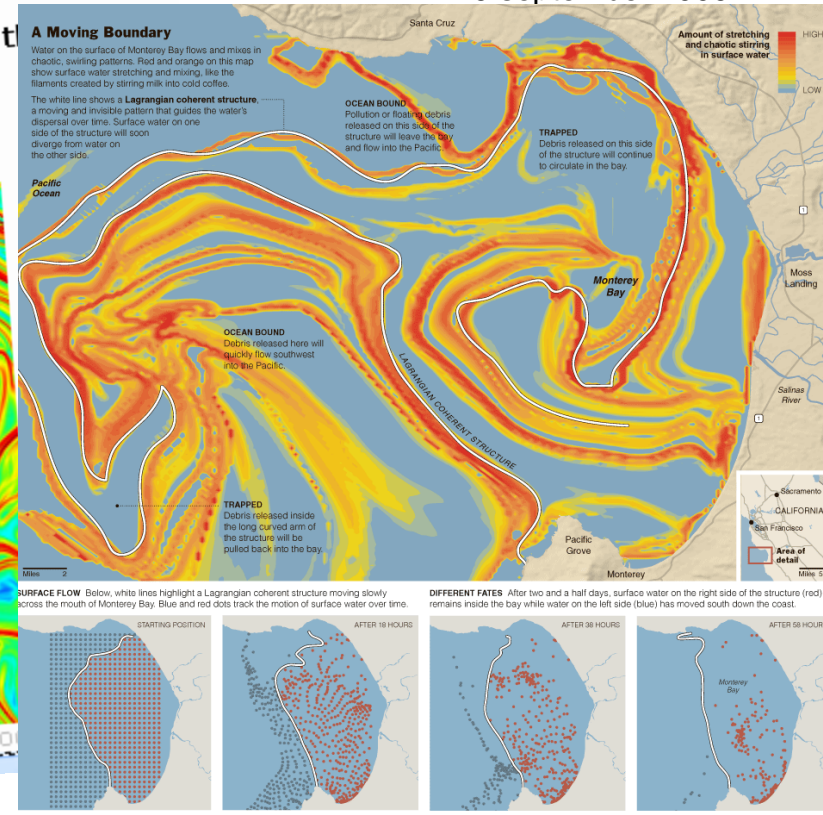
Nov 12th 2009 | From *The Economist* print edition



Thomas Peacock

The New York Times

29 september 2009

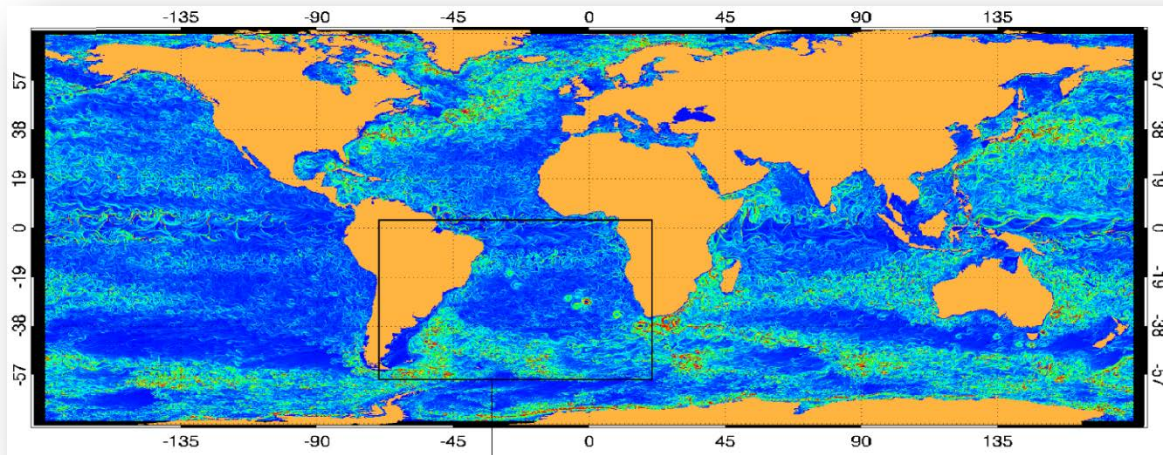


Sources: Francois Lekien, Université Libre de Bruxelles; Chad Coulllette, California Institute of Technology; Shawn C. Shadden, Illinois Institute of Technology

JONATHAN COBLE/THE NEW YORK TIMES

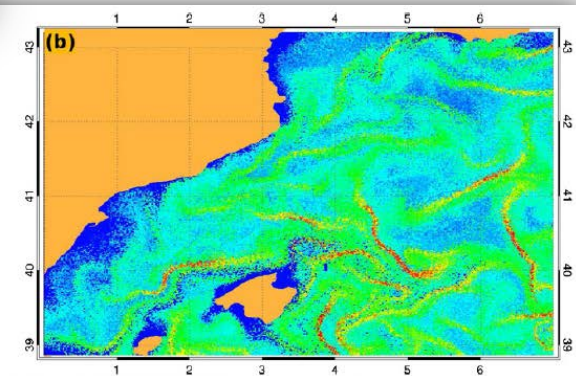
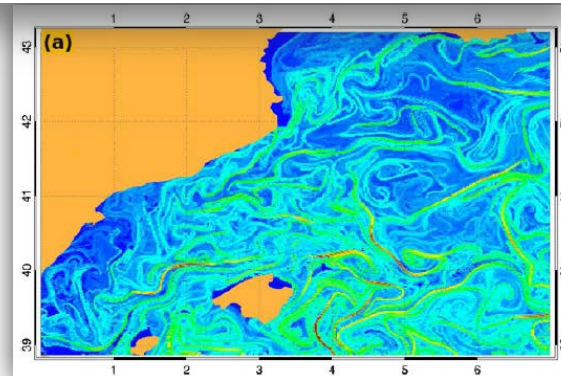
Any advantage in using FSLE to locate LCS?

In oceanographic contexts it is usually straightforward to identify the relevant spatial scales: Rossby radius, coastal features



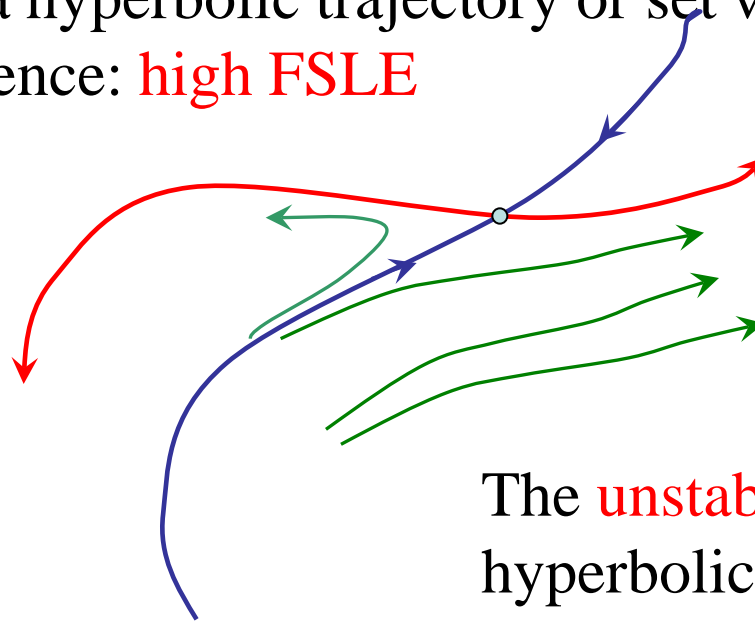
Trajectories can be nonsmooth
(noise ...)

I. Hernández-Carrasco et al.
Ocean Mod. 36, 208 (2011)



Disadvantage:
No theorems ...

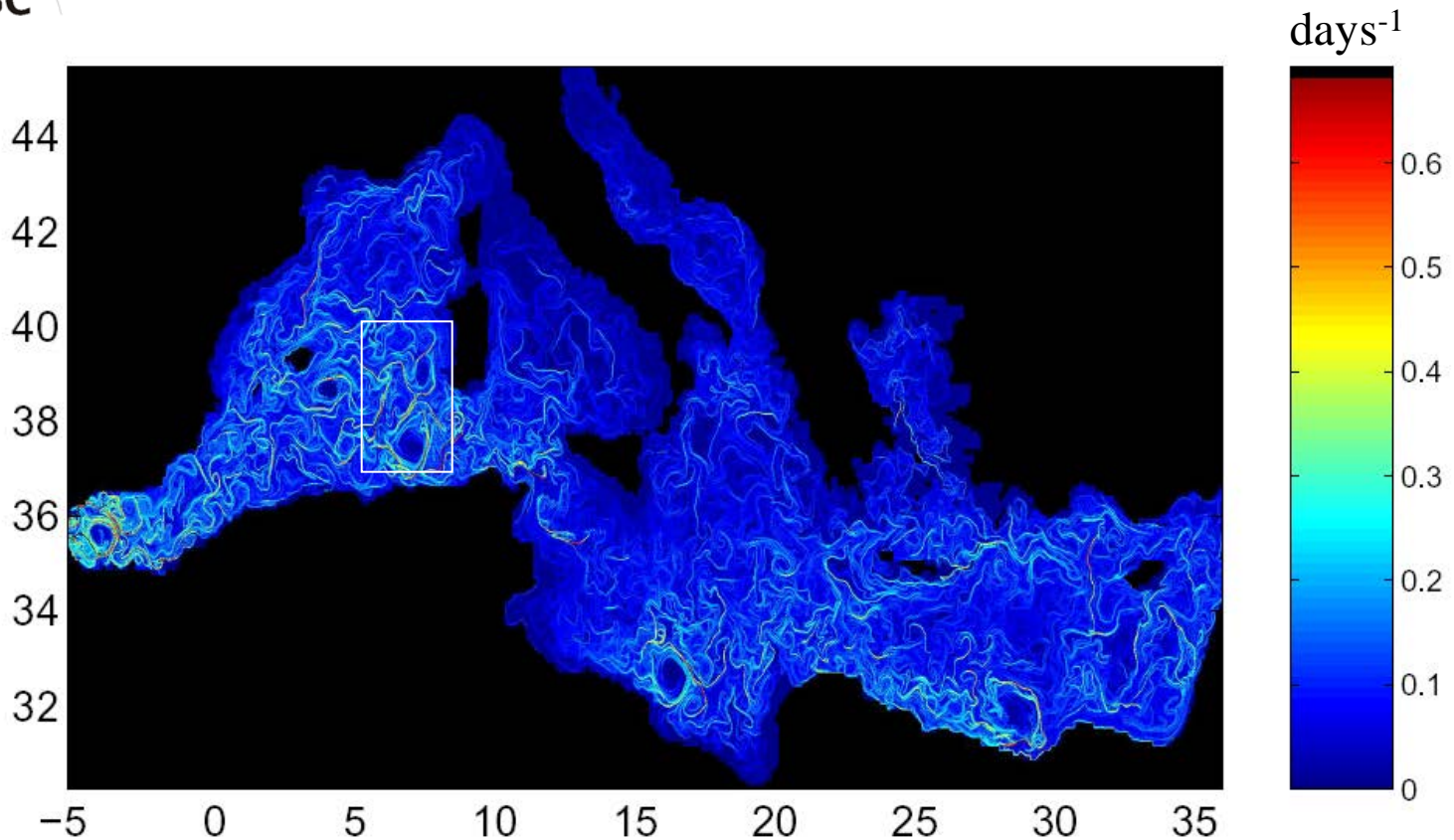
The idea is that initial conditions close to the **stable manifold** of a hyperbolic trajectory or set will show strong divergence: **high FSLE**



The **unstable manifold** of hyperbolic sets would be marked by **high FSLE in the time backwards** direction

Other types of Lyapunov exponents would display similar information, but FSLE is less affected by saturation

REMARK: these are heuristic consideration. Theorems needed (some available for FTLE)



DieCAST model for the full Mediterranean Primitive equations,
 48 vertical levels, $1/8^\circ$ horizontal resolution,
 climatological forcings ... \rightarrow 5 years of daily velocity fields

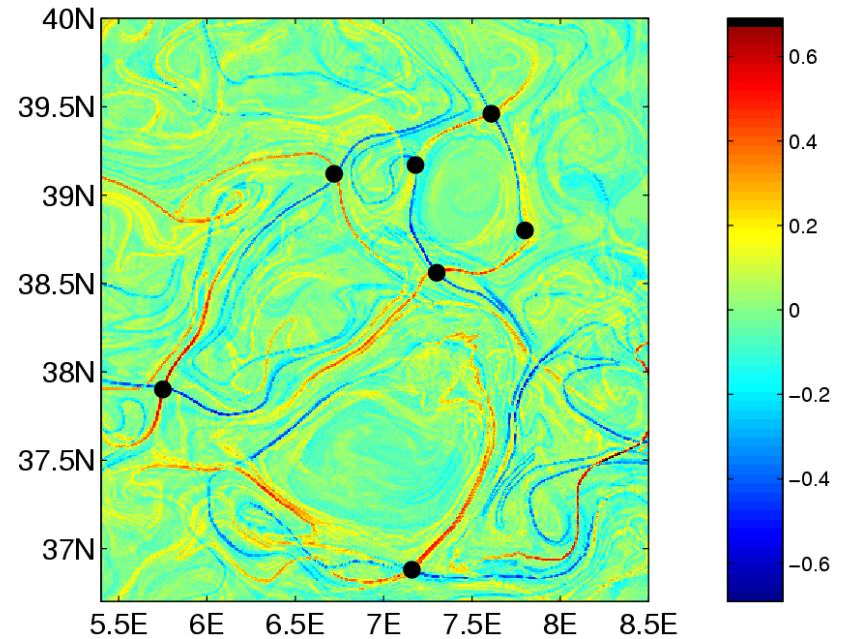
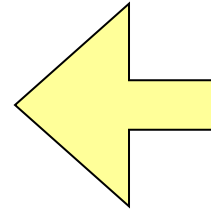
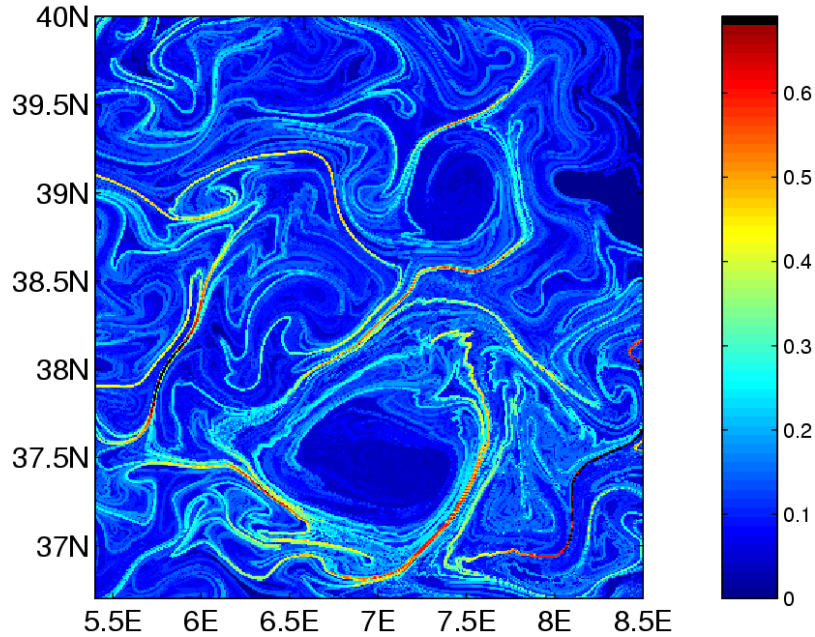
$\delta_0 = 0.02^\circ \rightarrow \delta_f = 1^\circ$ (mesoscale transport)

$\delta_0 \approx 2 \text{ km} \rightarrow \delta_f \approx 110 \text{ km}$ twodimensional

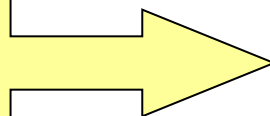
d'Ovidio, Fernández, Hernández-García, López, Geophys. Res. Lett. 31, L17203 (2004)

CLICK THE FIGURES FOR MOVIES

FSLE from time-backwards
Integrations.
Are they really unstable
manifolds of hyperbolic
trajectories?



FSLE from **forward**
and **backwards**
integrations

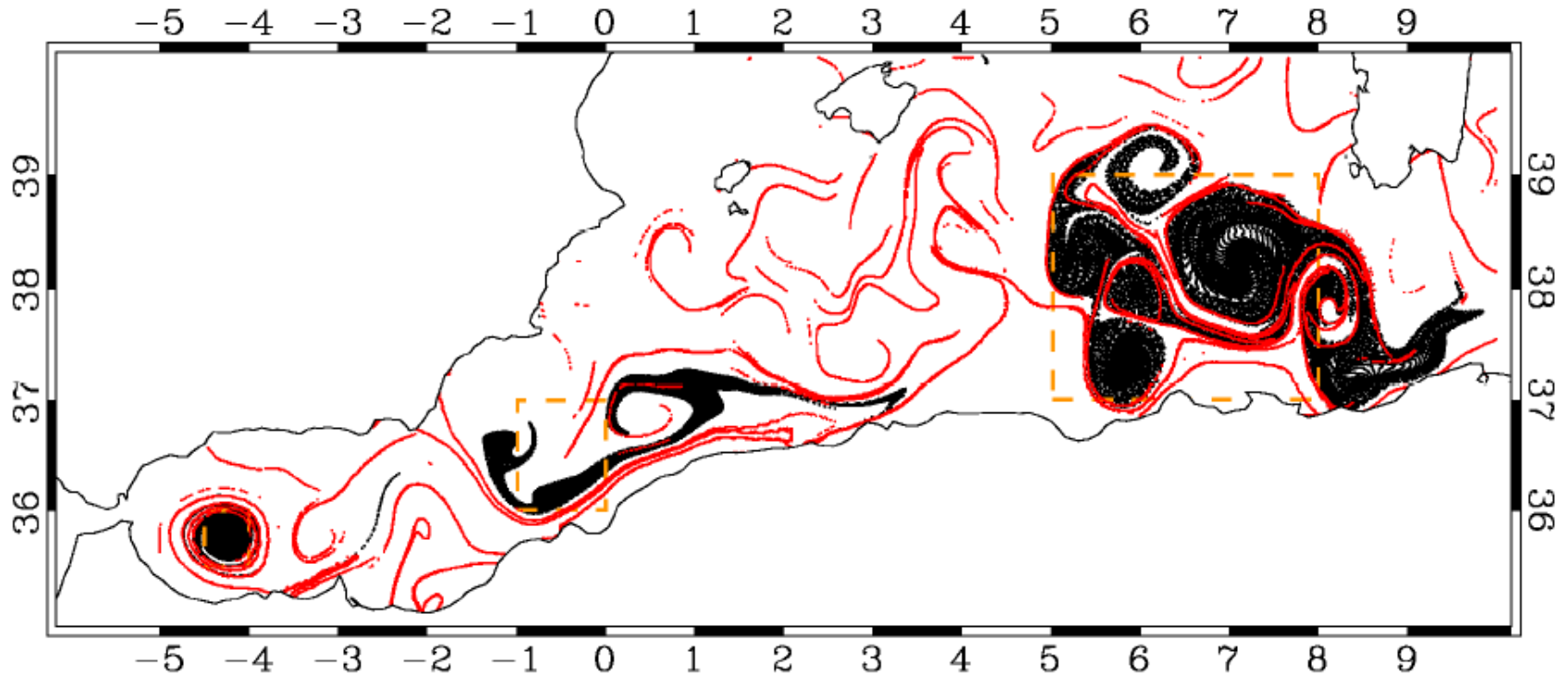


The strongest lines are seen to organize tracer flow

G, ShR

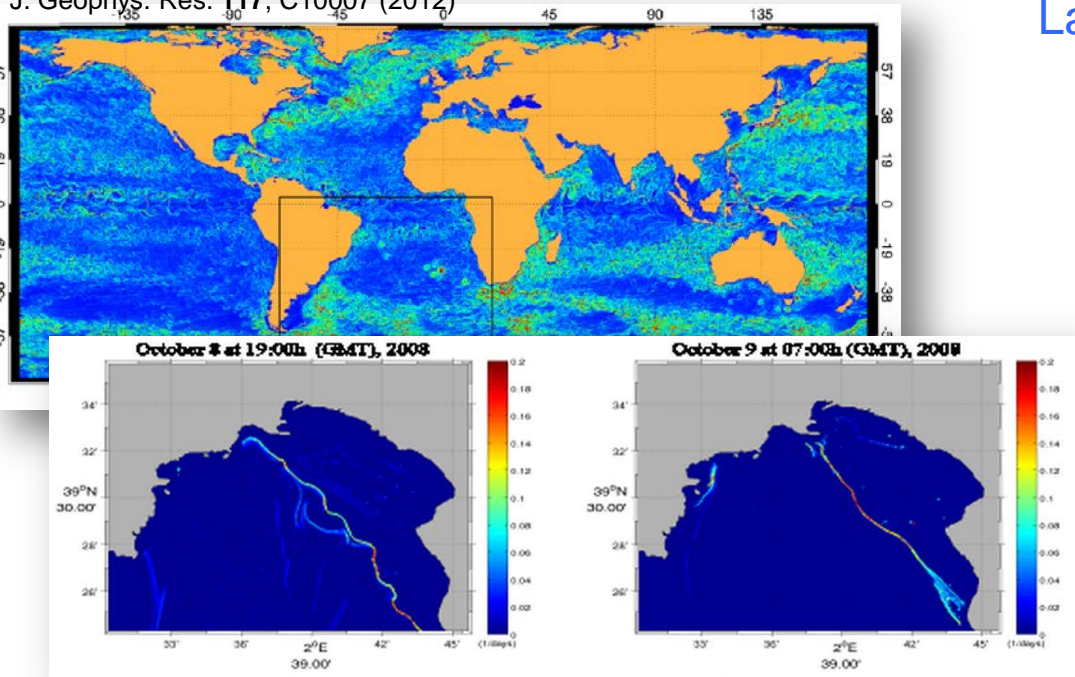
Lines: $FSLE > 0.2 \text{ days}^{-1}$

Tracer advection for 2 or 1 weeks



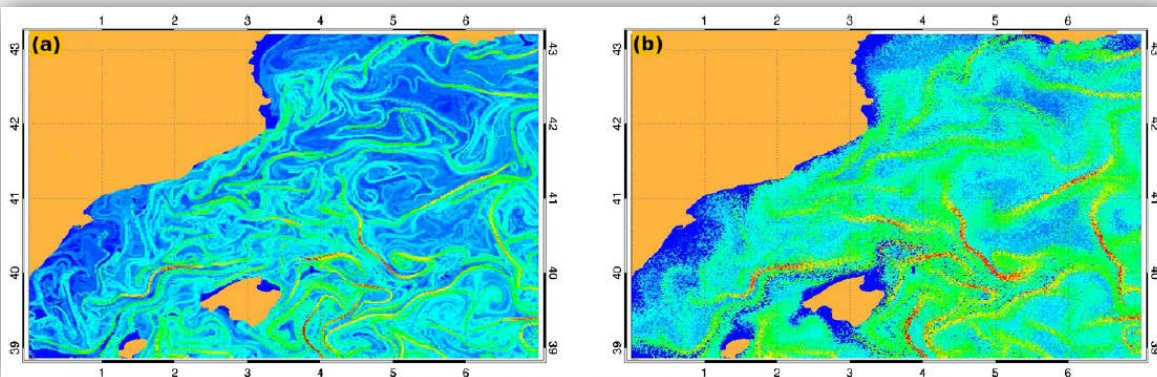
FSLE are Lagrangian, but not direct advection:

- shorter simulations
- no problems with exponentially increasing line lengths
- exhaustive consideration of initial conditions



Hernández-Carrasco, López, Orfila, Hernández-García,
 Nonlinear Processes in Geophysics **20**, 921-933 (2013)

Bahía de Palma



Any advantage in using FSLE in Lagrangian studies?

- Easy switching between local and statistical approaches
- In oceanographic contexts it is usually straightforward to identify the relevant spatial scales: Rossby radius, coastal features
- Trajectories can be nonsmooth (noise ...)

Disadvantages:

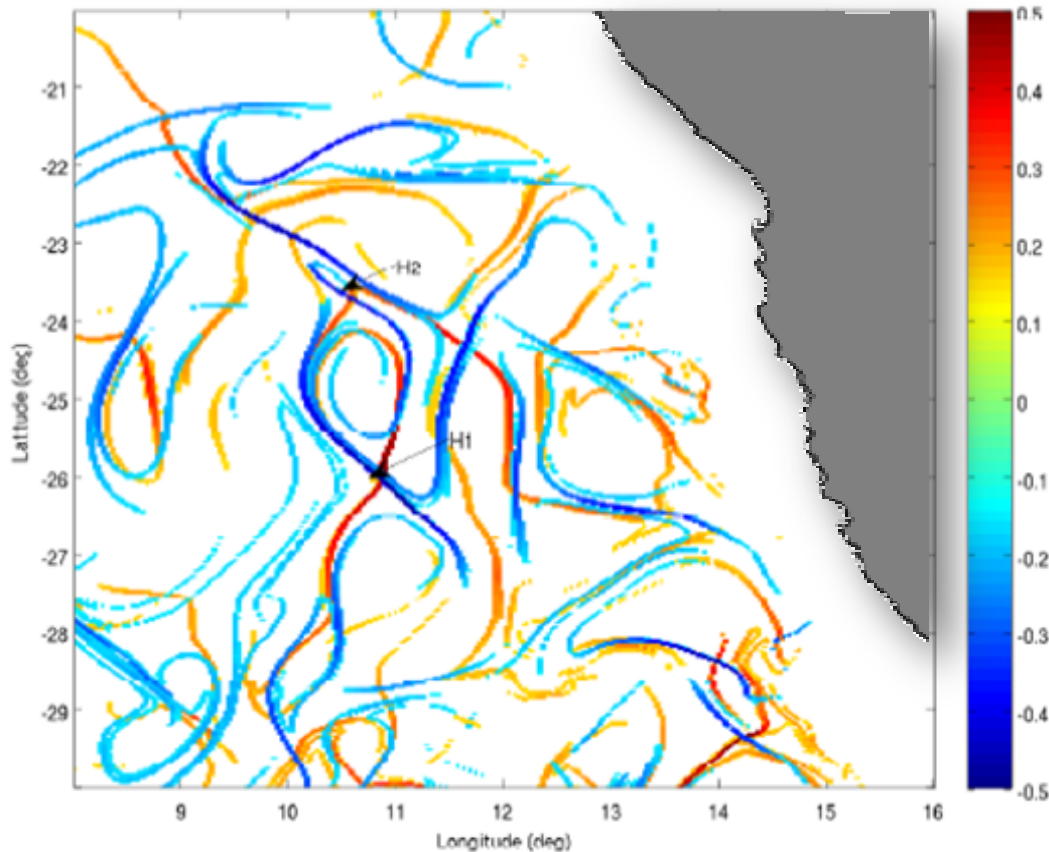
- No distinction between hyperbolic, shear, ... structures
- Lack of analytical approaches (but see Tzella and Haynes, PRE 2010, Karrasch and Haller, Chaos 2013)
- As for FTLE, not all high FSLE structures have a clear impact on flows. Need to check with actual particle trajectories

Hernández-Carrasco et al.
 Ocean Mod. **36**, 208 (2011)

Some examples of recent Lagrangian studies in the ocean using Finite Size Lyapunov Exponents

Three-dimensional characterization flow and eddies in Benguela

J.H. Bettencourt, C. Lopez, E. Hernandez-Garcia, Ocean Modelling 51 (2012) 73–83



ROMS model:

(from Gutknecht et al.(2013).
and Le Vu et al.)

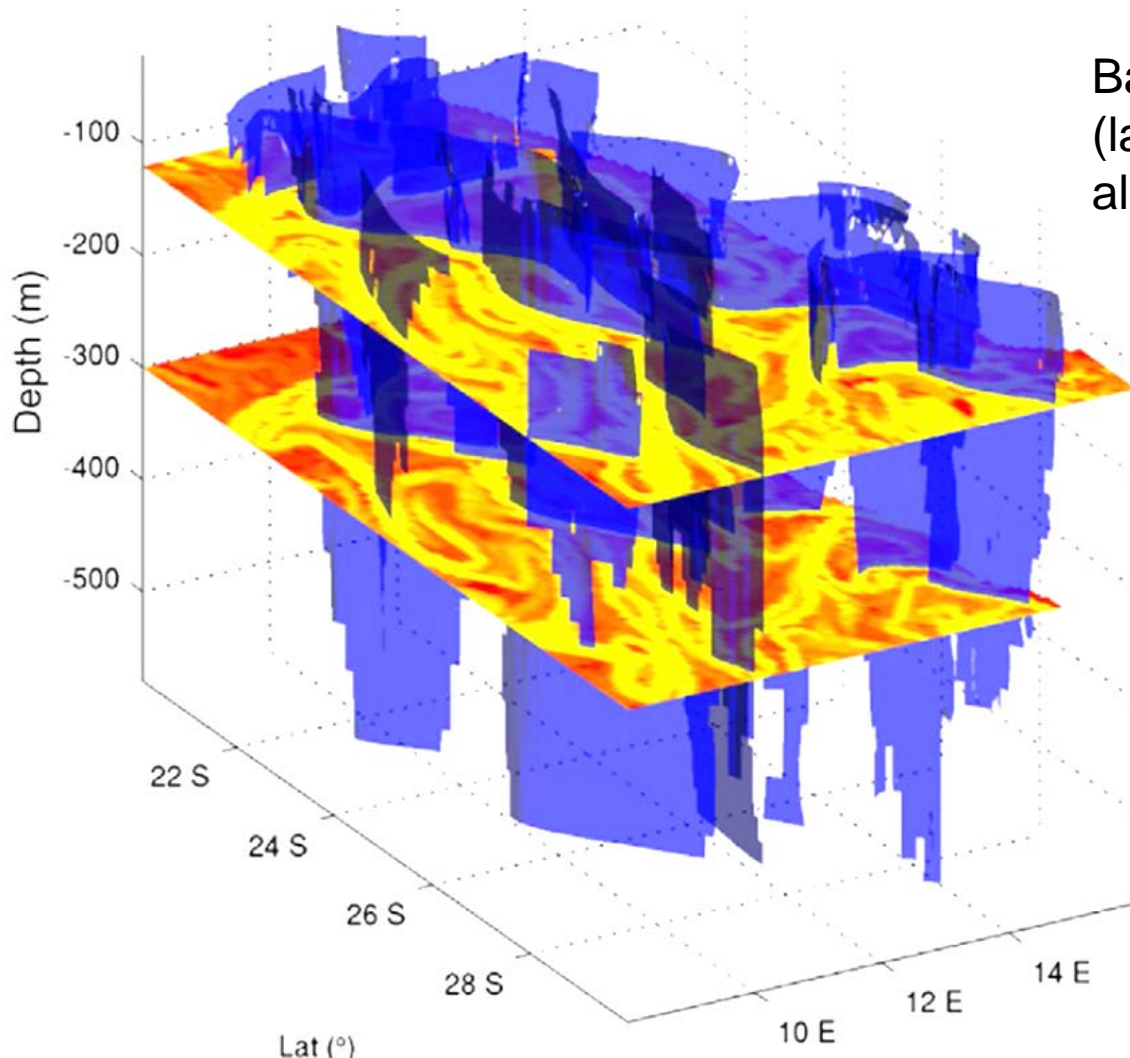
2 years of simulation,
climatologically forced.

Horizontal resolution
1/12 degrees (8 km)
32 vertical terrain-following
levels

Forward and backward
FSLE fields

$\delta_0=2$ km ; $\delta_f=100$ km

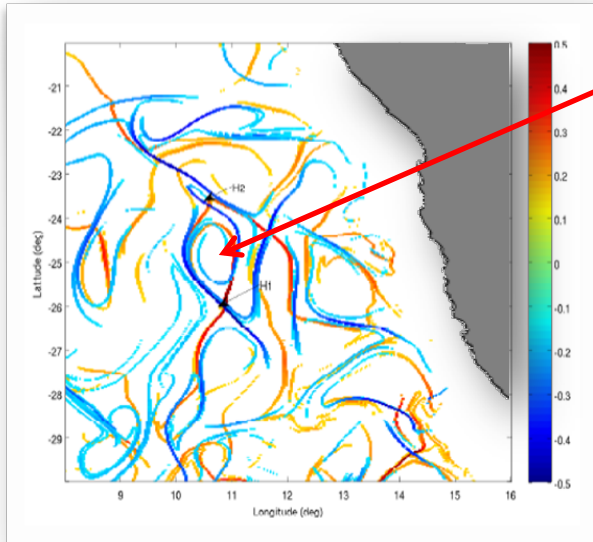
Particles released in horizontal planes every 20 m and integrated in 3D



Backward FSLE from a (largest) ridge extracting algorithm

Curtain-like structure as arising when vertical shear of horizontal velocities much smaller than horizontal velocities

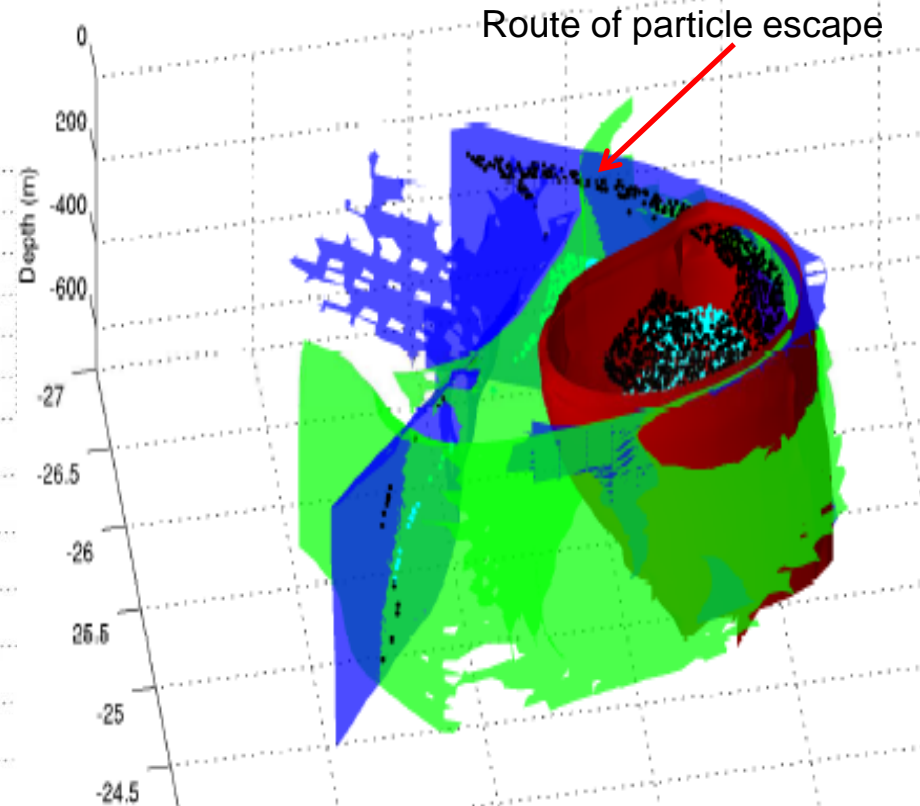
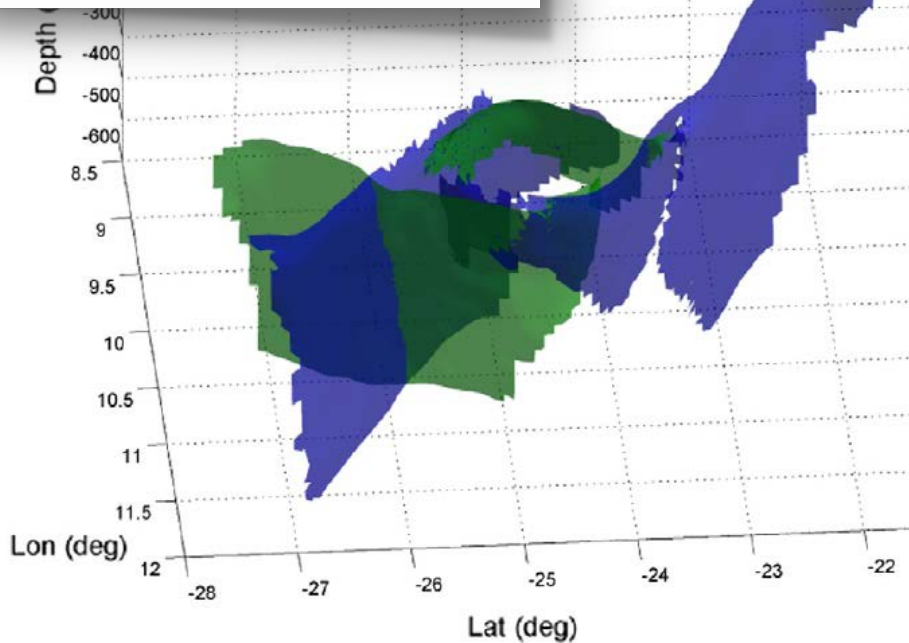
(Branicki, Mancho, Wiggins, Physica D 240 (2011) 282–304)



Particular eddy enclosed by hyperbolic manifolds

BACKWARDS FSLE
FORWARD FSLE
Q-criterion isosurface

FSLE methodology is giving the hyperbolic filamentation region, not the coherent core

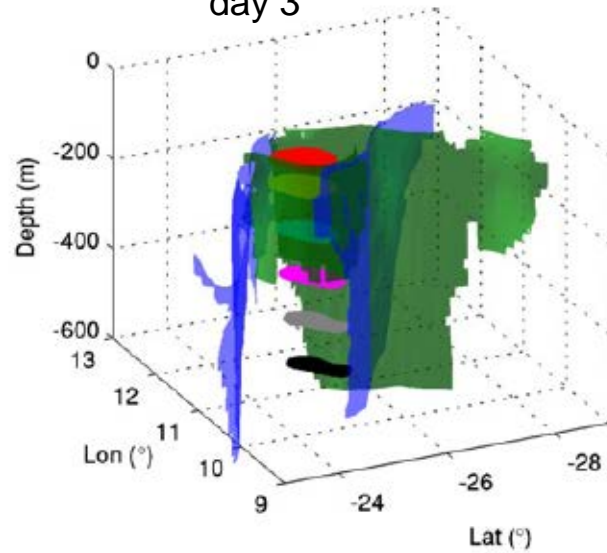


J.H. Bettencourt, C. Lopez, E. Hernandez-Garcia
Ocean Modelling 51 (2012) 73–83
J. Phys. A 46 (2013), 254022

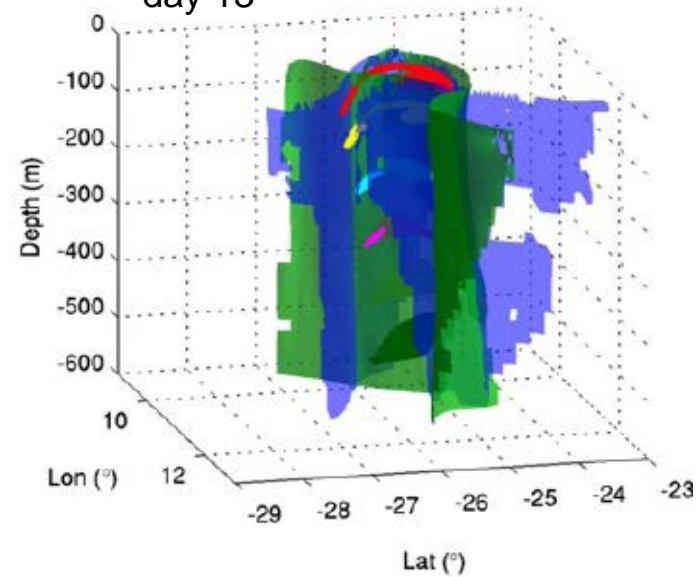
BACKWARDS FSLE
FORWARD FSLE

Red: 40 m
yellow: 100 m
cyan: 200 m
magenta: 300 m
grey: 400 m
black: 500 m

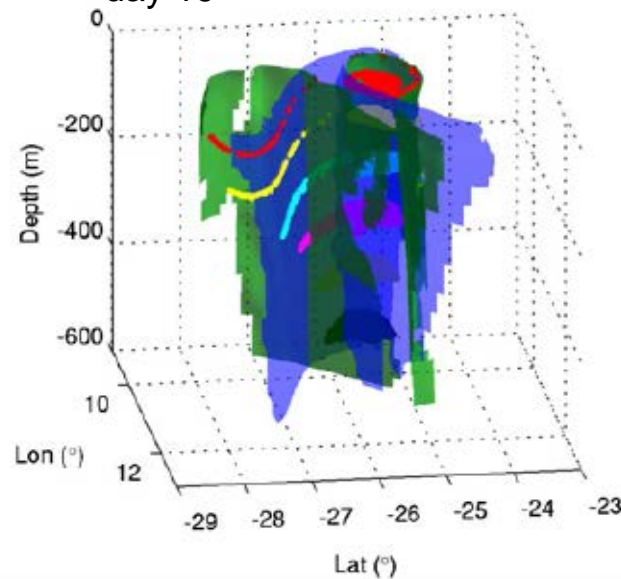
day 3



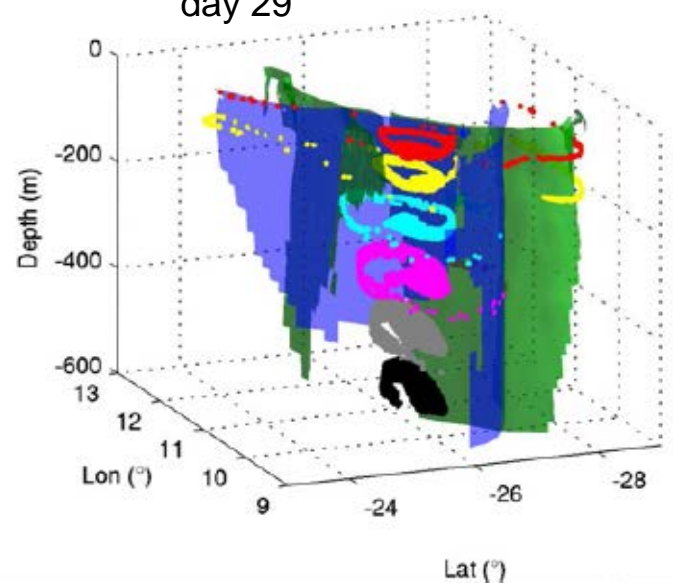
day 13



day 19



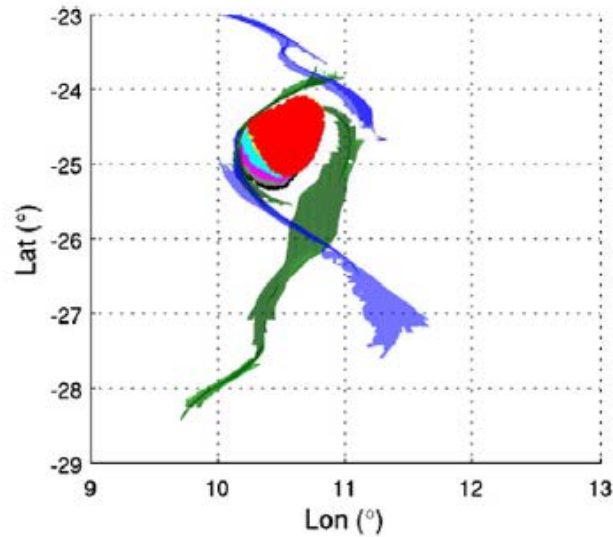
day 29



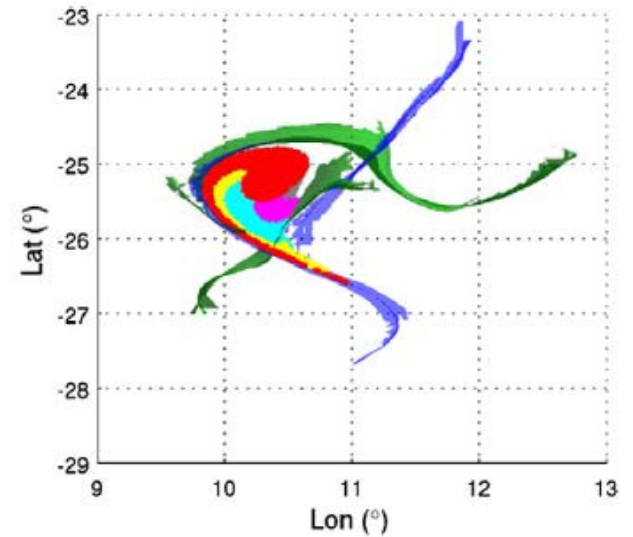
BACKWARDS FSLE
FORWARD FSLE

Red: 40 m
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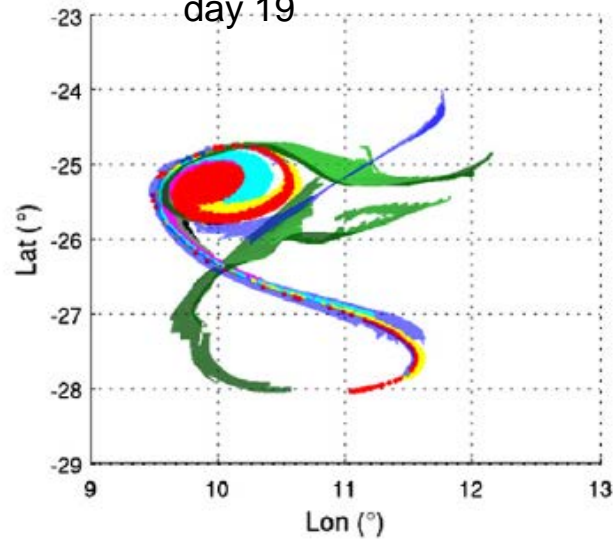
day 3



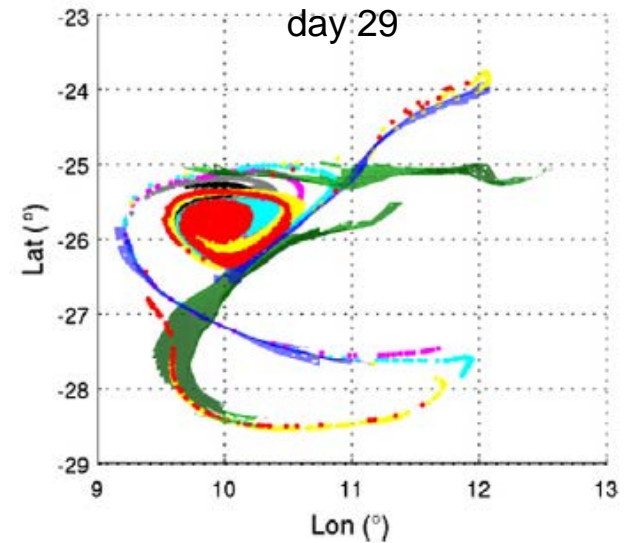
day 13

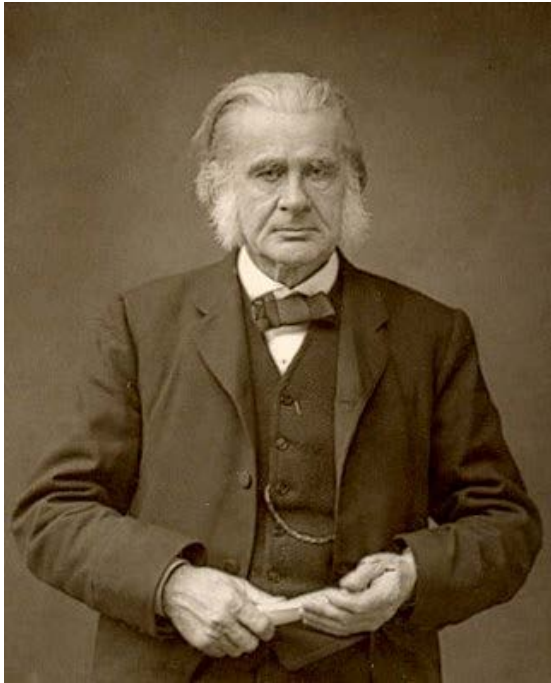


day 19



day 29





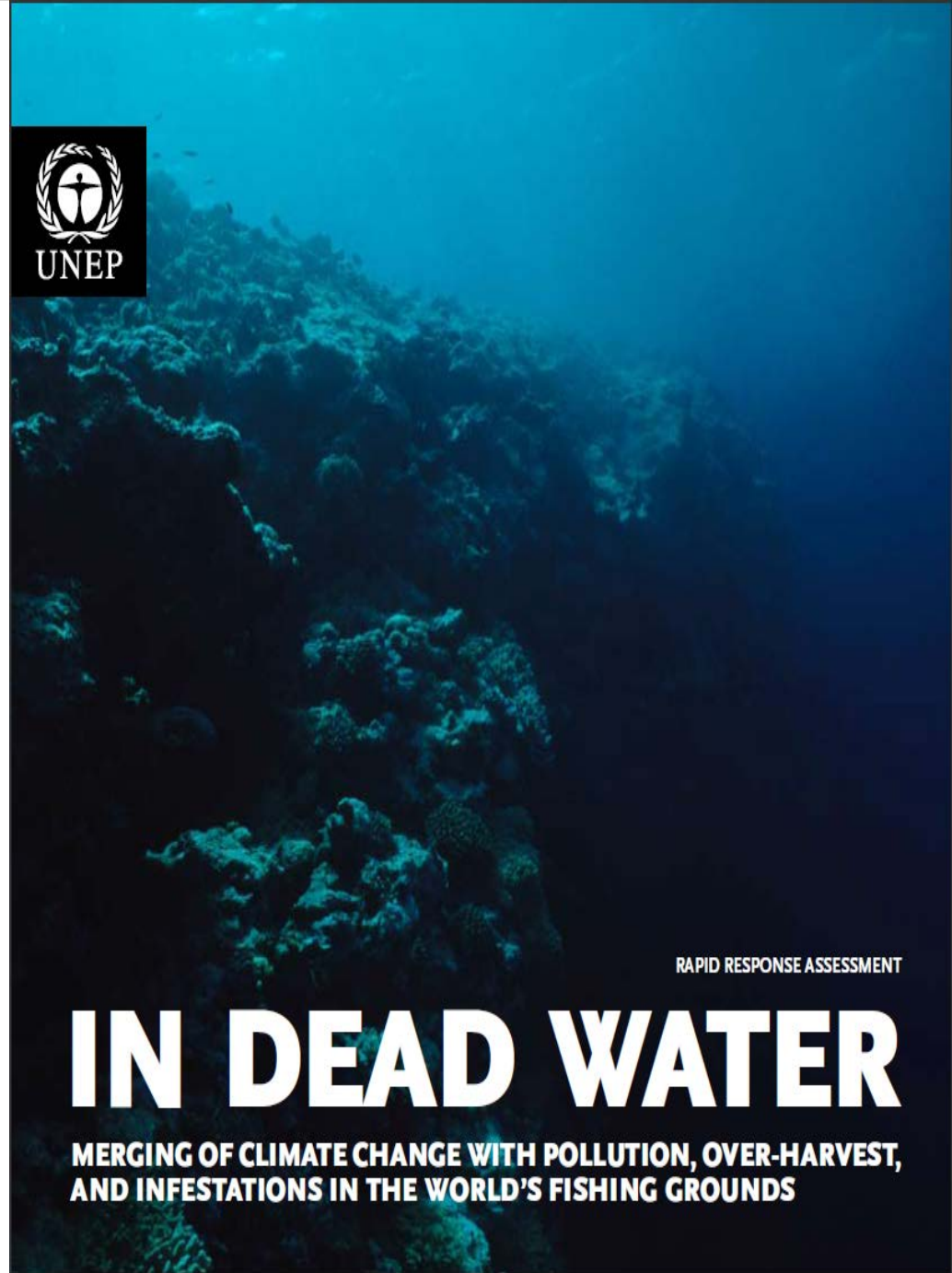
“ I believe, then, that the cod fishery, the herring fishery, the pilchard fishery, the mackerel fishery, and probably all the great sea fisheries, are inexhaustible; that is to say, that nothing we do seriously affects the number of the fish. And any attempt to regulate these fisheries seems, consequently, from the nature of the case, to be useless.”

Thomas H. Huxley, Intern. Fisheries Exhibition, London (1883)



“Increased development, coastal pollution and climate change impacts on ocean currents will accelerate the spreading of marine dead zones, many around or in primary fishing grounds.”

*United Nations Environmental Programme
(2008)*

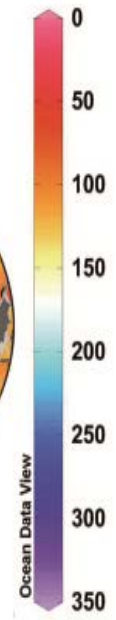
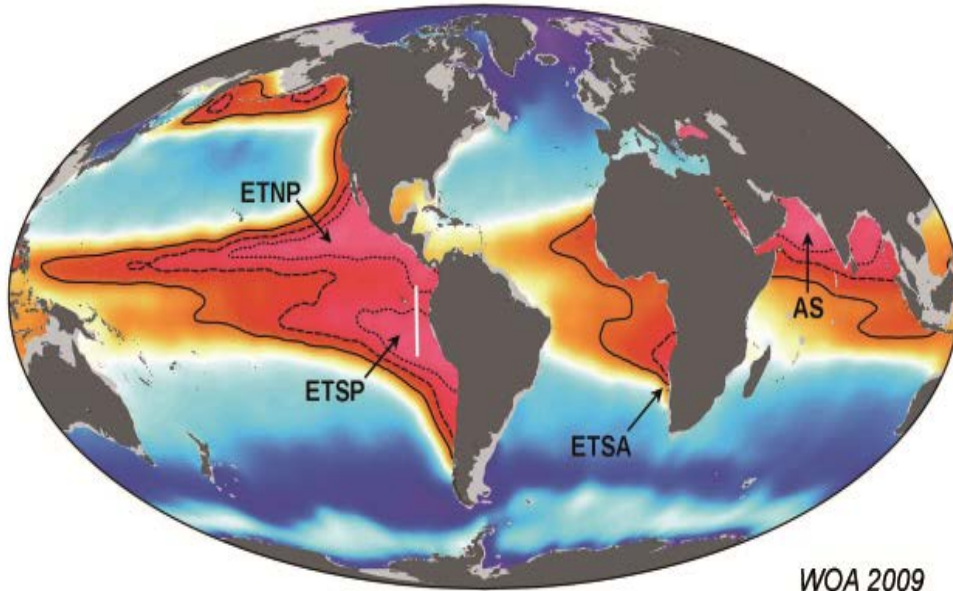


RAPID RESPONSE ASSESSMENT

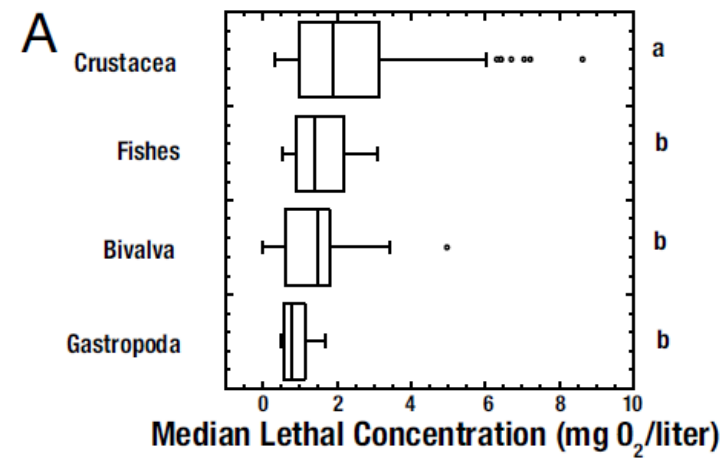
IN DEAD WATER

**MERGING OF CLIMATE CHANGE WITH POLLUTION, OVER-HARVEST,
AND INFESTATIONS IN THE WORLD'S FISHING GROUNDS**

Oxygen (μM) at 300 m



Hypoxic levels - $\text{O}_2 < 88 \mu\text{M}$



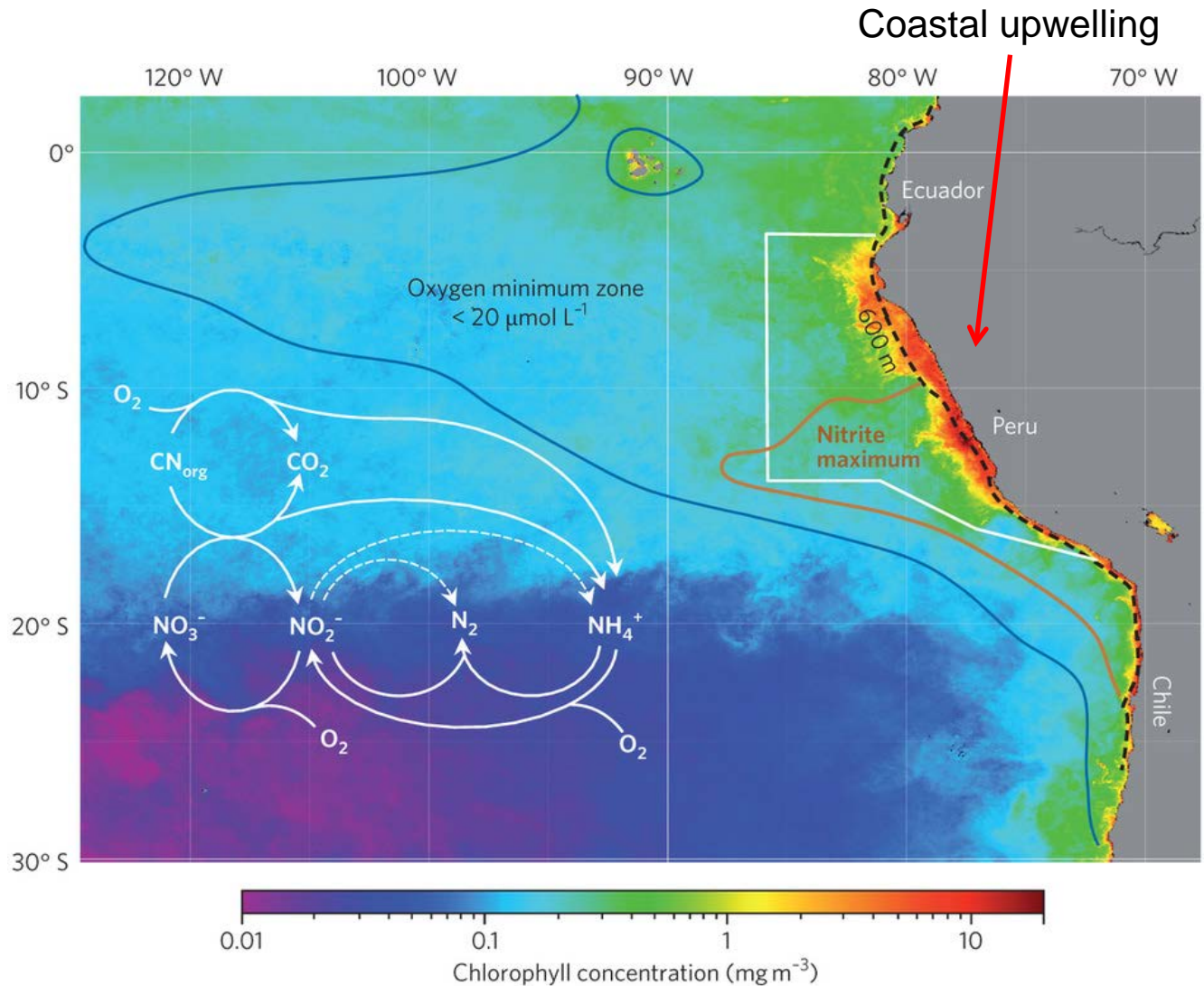
Vaquer-Sunyer & Duarte (2008)

Respiration and nitrification consume oxygen

Increased stratification associated to global warming will make things worse

Role of flow: Large scale patterns induce low ventilation areas.

What about horizontal stirring and mixing?



Thamdrup (2013)

ROMS hydrodynamic model:

3D primitive equations

Hydrostatic

Terrain following

Forced by climatology

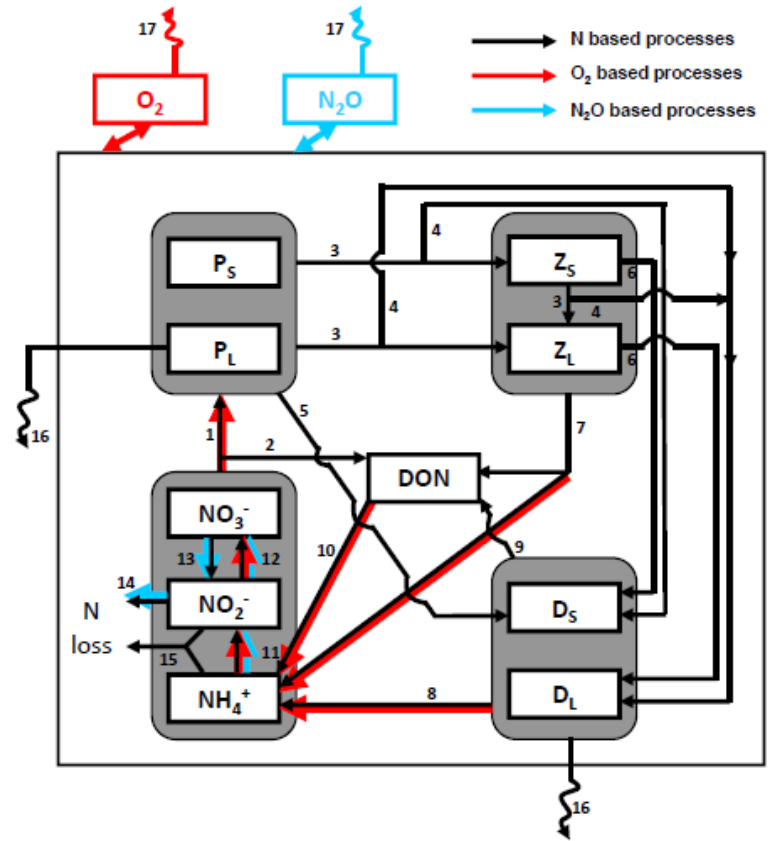
horizontal resolution of 1/9 degrees (~ 12 km)

32 terrain-following vertical levels

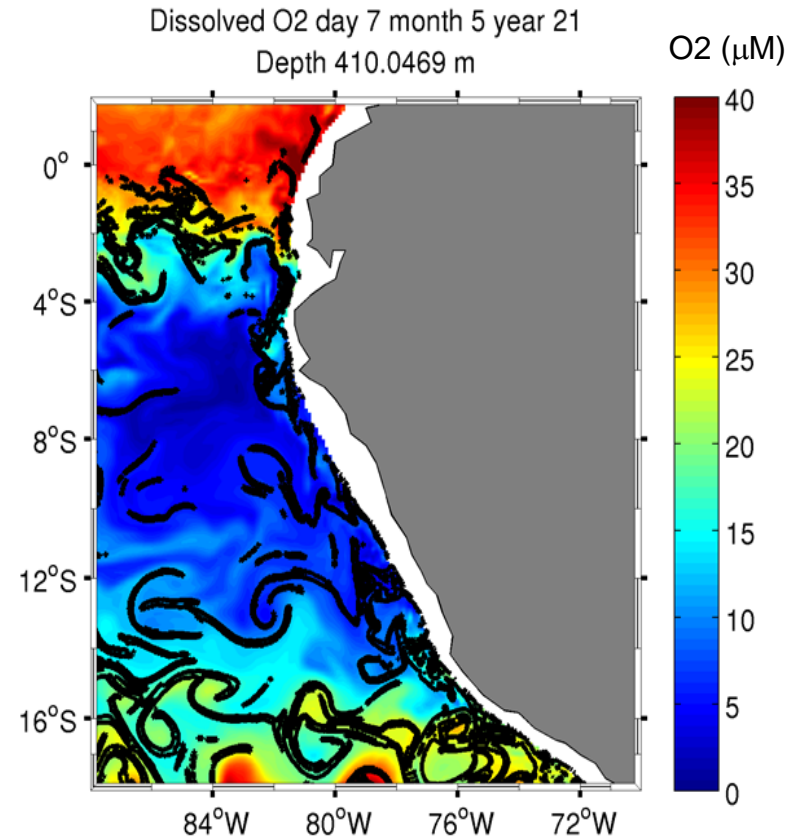
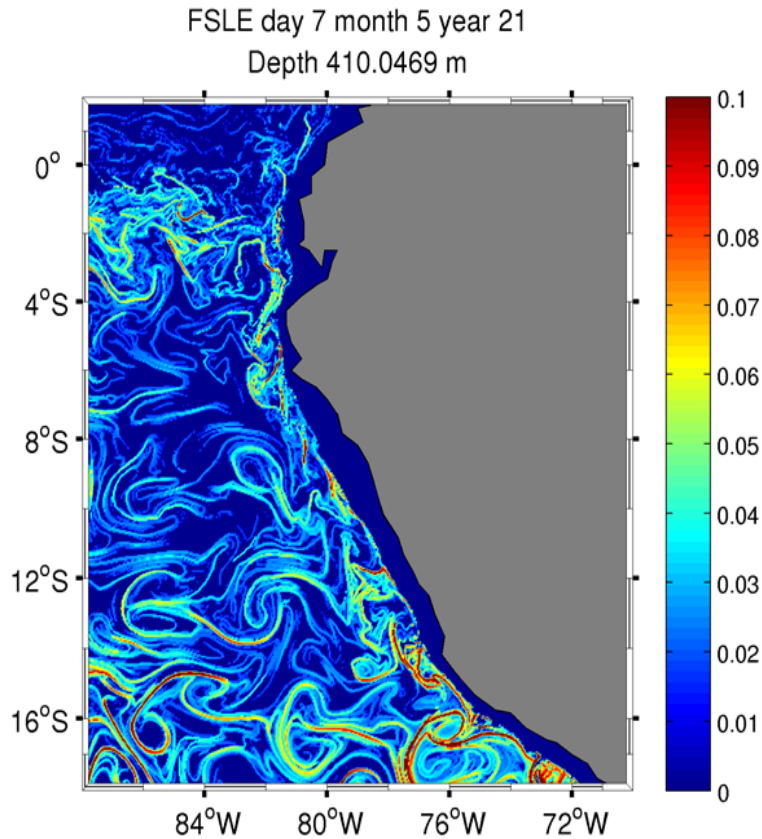
BioEBUS biogeochem. model:

$$\frac{\partial C_i}{\partial t} = -\nabla \cdot (\mathbf{u}C_i) + K_h \nabla^2 C_i + \frac{\partial}{\partial z} \left(K_z \frac{\partial C_i}{\partial z} \right) + SM$$

(Gutknecht et al, 2013)



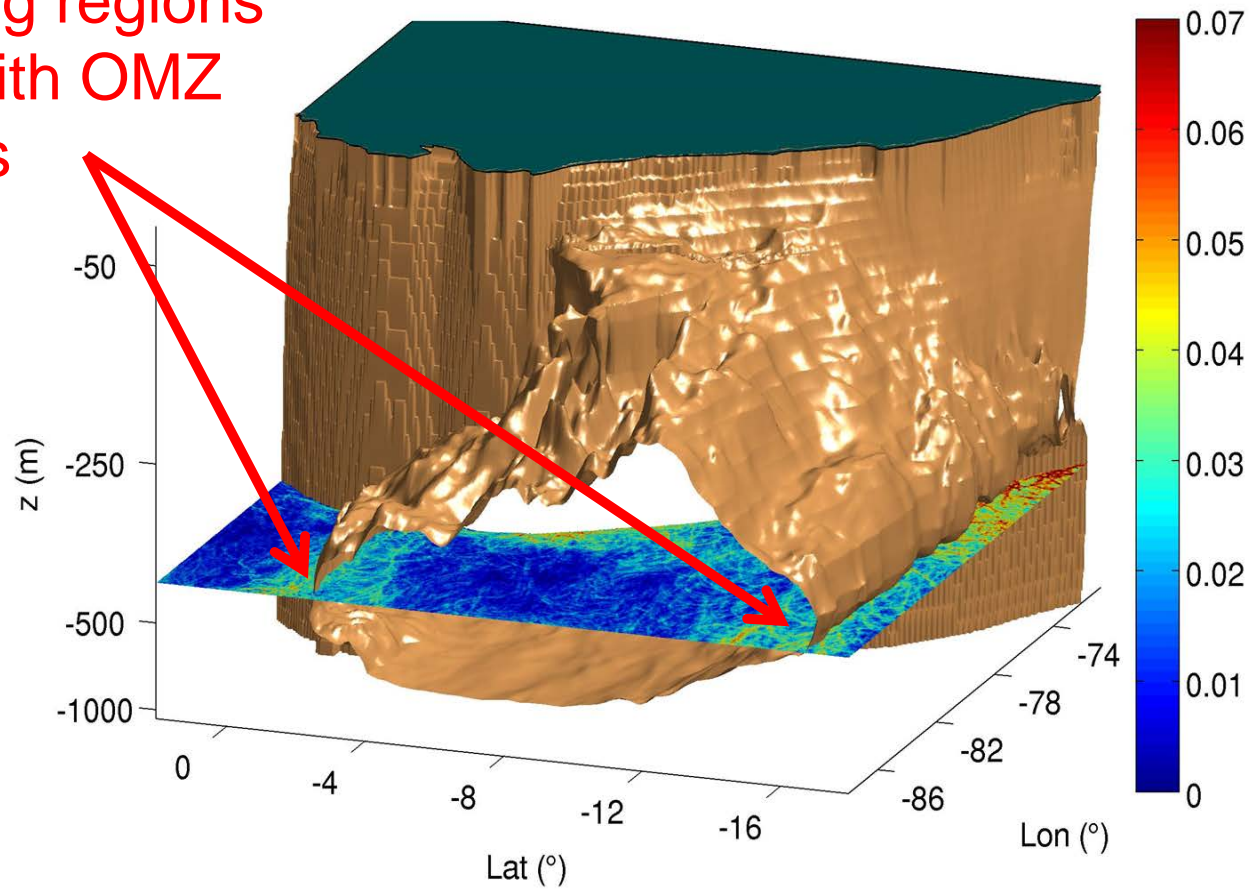
- | | | |
|-------------------------------|------------------------------|------------------------|
| 1. Assimilation of nutrients | 6. Mortality of zooplankton | 11,12. Nitrification |
| 2. Exudation | 7. Excretion | 13,14. Denitrification |
| 3. Grazing | 8. Decomposition of detritus | 15. Anammox |
| 4. Fecal pellets | 9. Hydrolysis | 16. Vertical sinking |
| 5. Mortality of phytoplankton | 10. Decomposition of DON | 17. Sea-air flux |



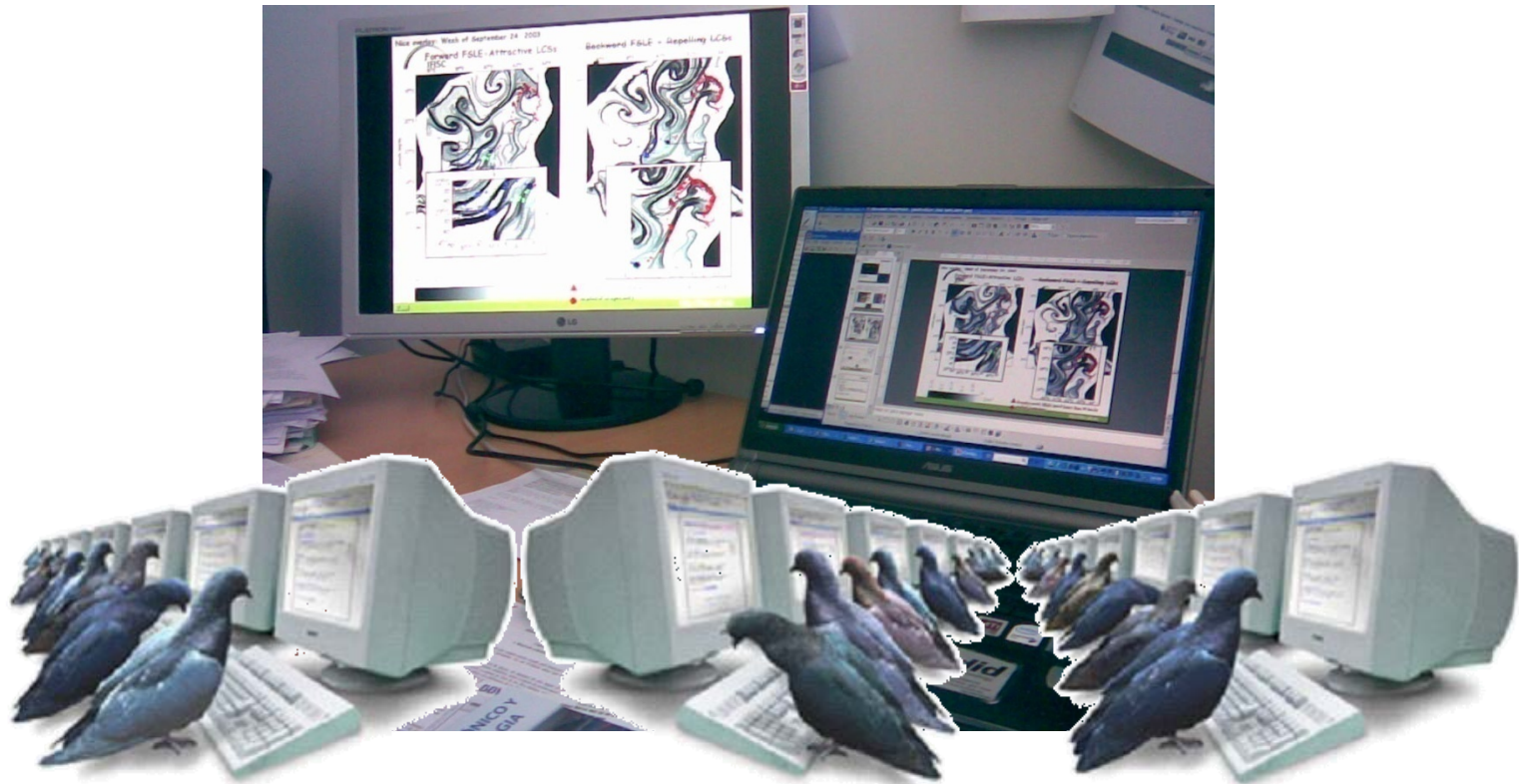
Backward FSLE (day^{-1})
 Particles released in horizontal planes and integrated in 3D
 $\delta_0=4 \text{ km}$; $\delta_f=100 \text{ km}$

Mean Oxygen Minimum Zone boundary at 20 μM

High stirring regions coincide with OMZ boundaries



Do birds know about Lyapunov exponents?



Tew Kai, Rossi, Sudre, Weimerskirch, Lopez, Hernandez-Garcia, Marsac, Garçon,
PNAS 106, 8245 (2009)



Satellite transmitter and altimeter
(total weight : 1 to 3% mass of adults,
max 45g)

8 birds (from Europa Island community) fitted with satellite transmitter and altimeter.

Followed for their foraging trips from August 18 to September 30, 2003.

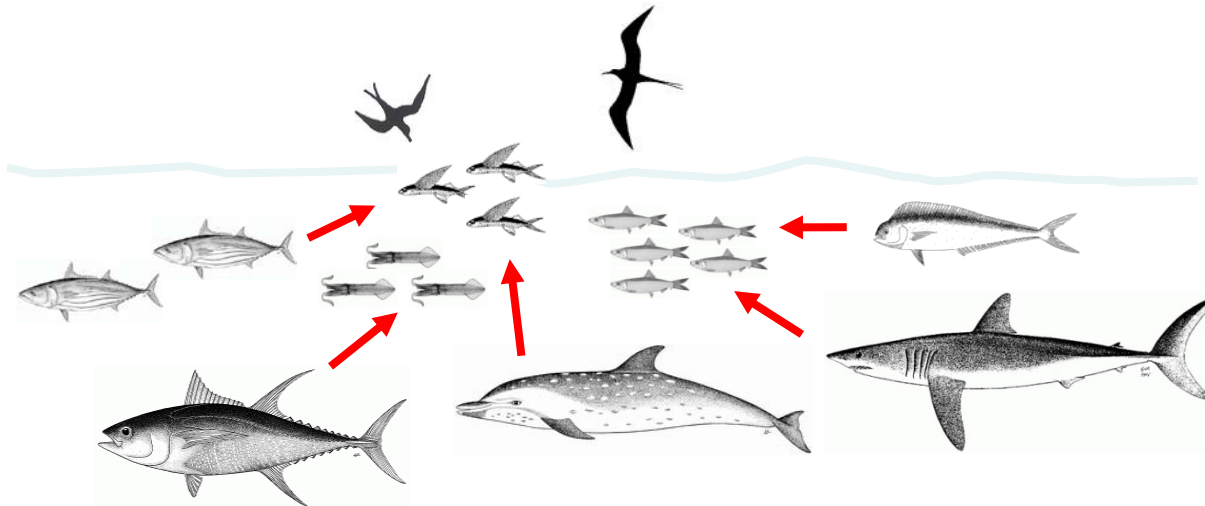
1600 Argos from 50 trips positions, distributed into 17 long trips (> 614 km) and 33 short trips.

(Weimerskirch et al., 2004)

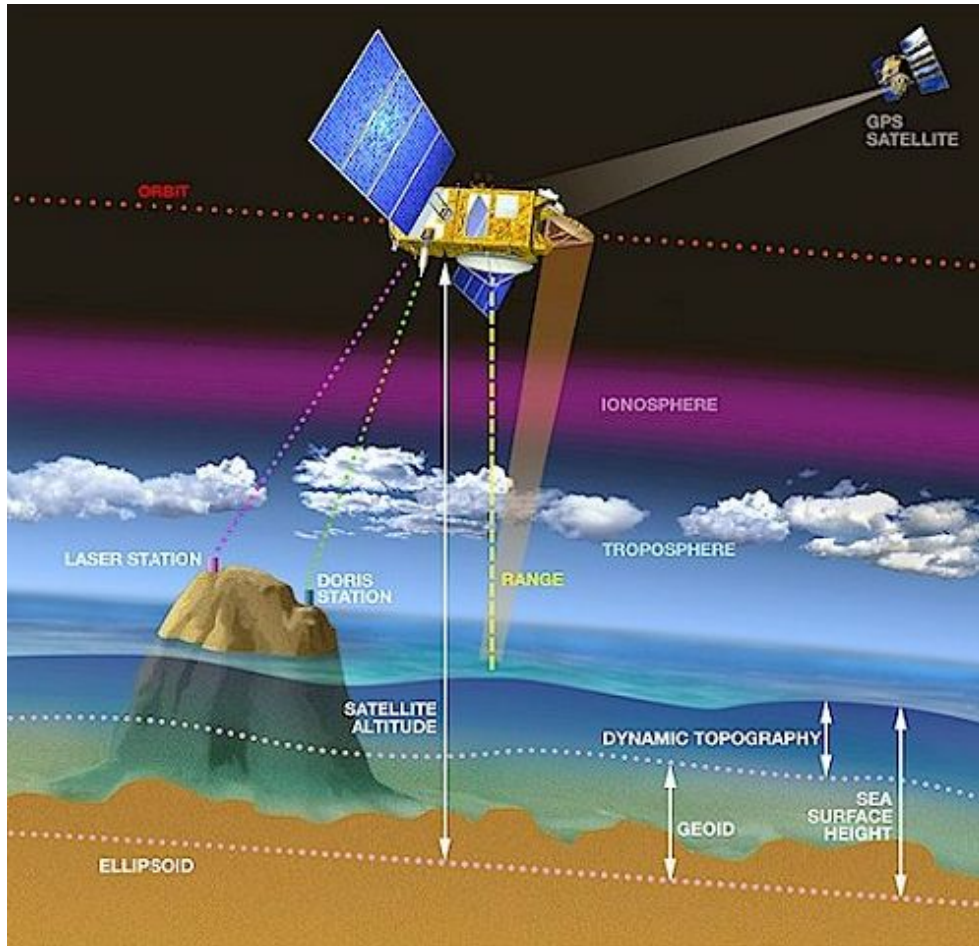


Great frigatebird (*fregata minor*):

- Large seabirds (light weight < 5 kg and large wings > 2m). Use thermals to soar before gliding over long distances and time (days/nights over weeks).
- Traveling at high altitudes to locate patches of prey and come close to surface to feed (reduced flight speed indicates foraging).
- Feeding occurs only during daytime (peaks in the morning and evening).
- Unable to dive or rest on the water surface (permeable plumage) → in association with subsurface predators (tuna, ...): **fisheries indicators**



SATELLITE ALTIMETRY FROM TOPEX/POSEIDON, ERS-2, JASON, ENVISAT, ...



Dynamic Topography (DT)=
Sea Surface Height (SSH) – Geoid (G)

SSH \approx 3 cm

G \approx meters ...

Sea Level Anomalies (SLA) =
SSH - \langle SSH \rangle_t = DT - \langle DT \rangle_t

Dynamic topography determines, via the Coriolis force, the velocity field (at large scales, geostrophic approximation)

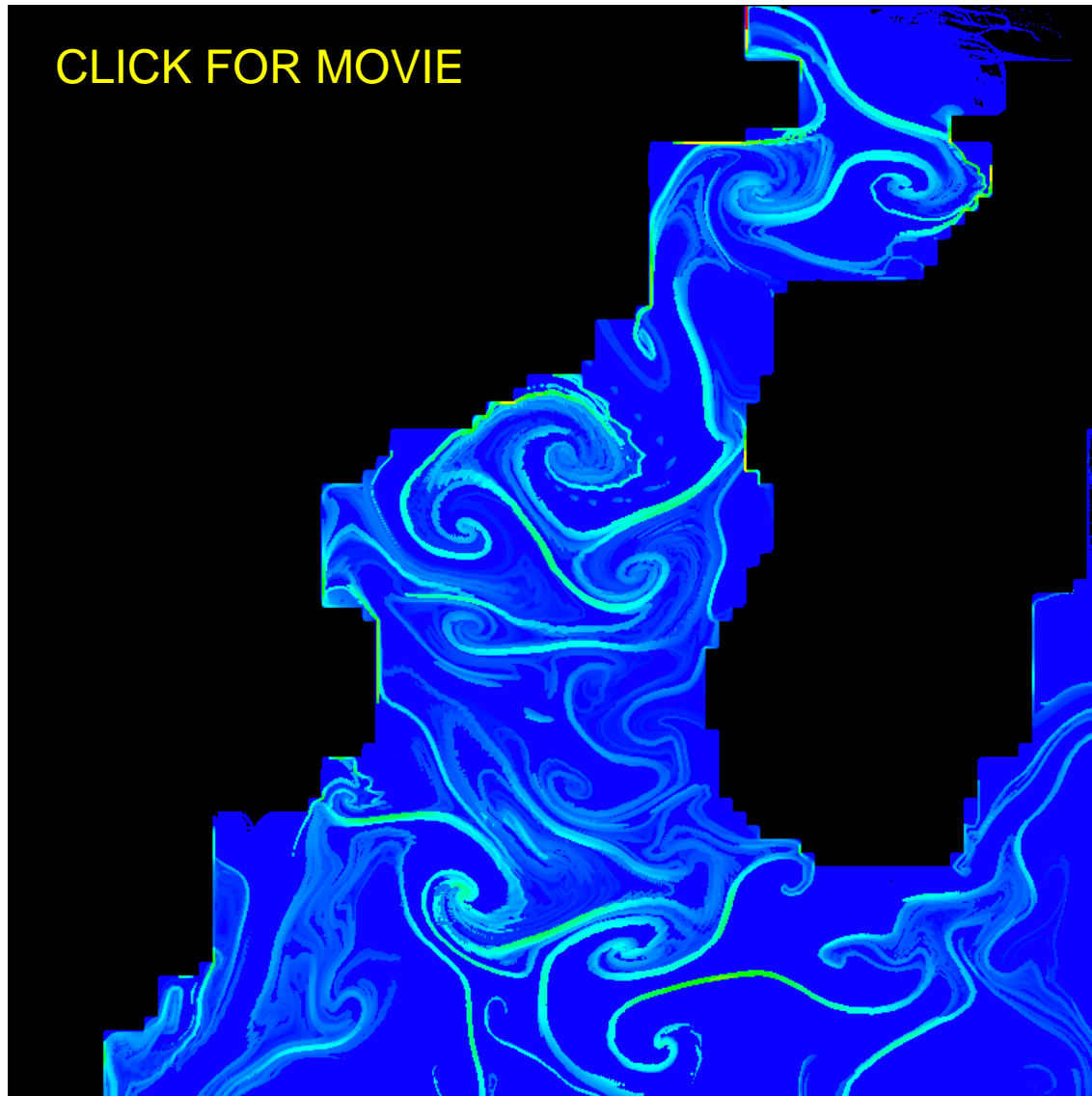
Ageostrophic components
Can be estimated from
scatterometer data

(Surface roughness \rightarrow wind \rightarrow Ekman component)

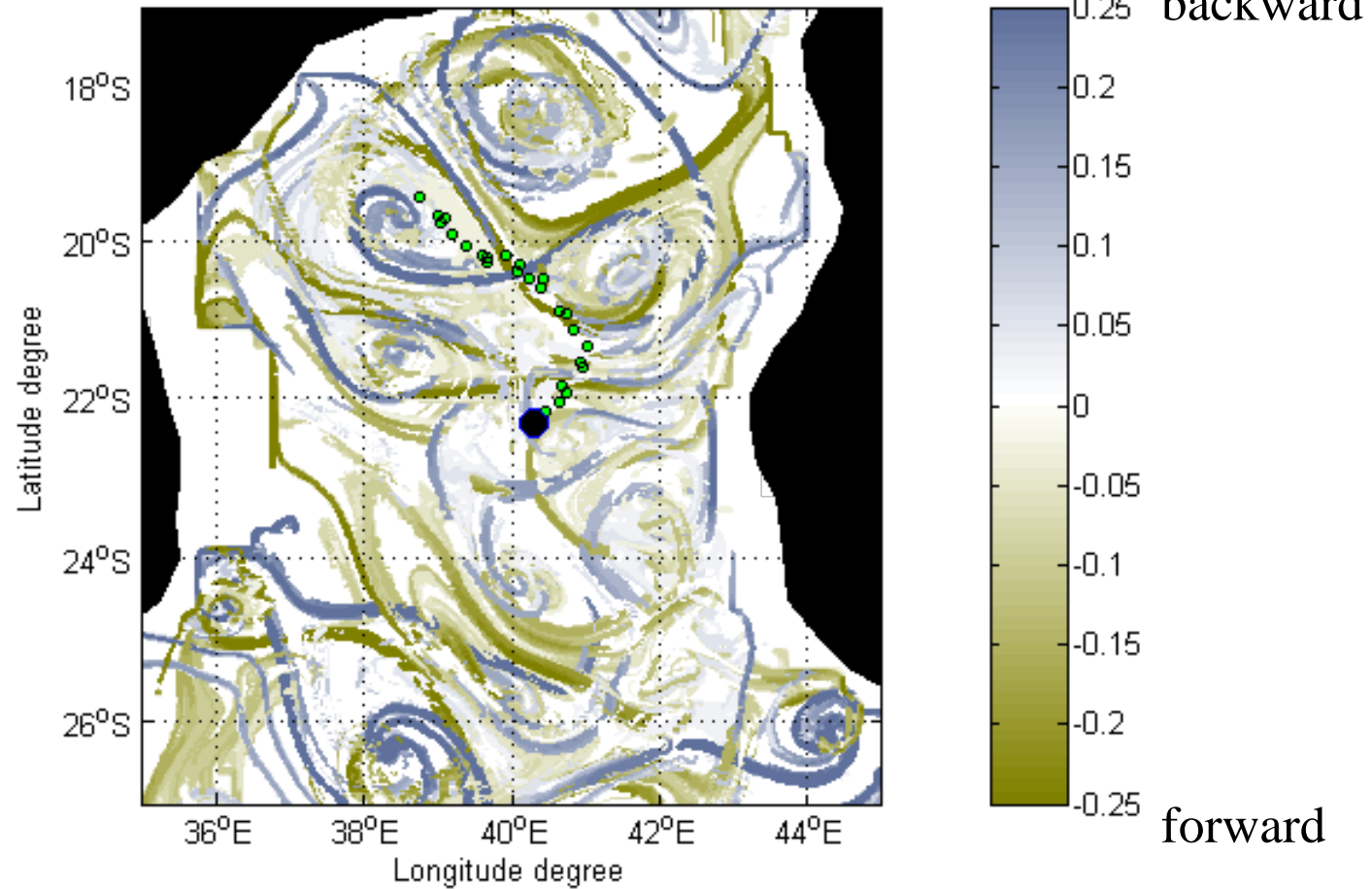
CLICK FOR MOVIE

Backwards
FSLE

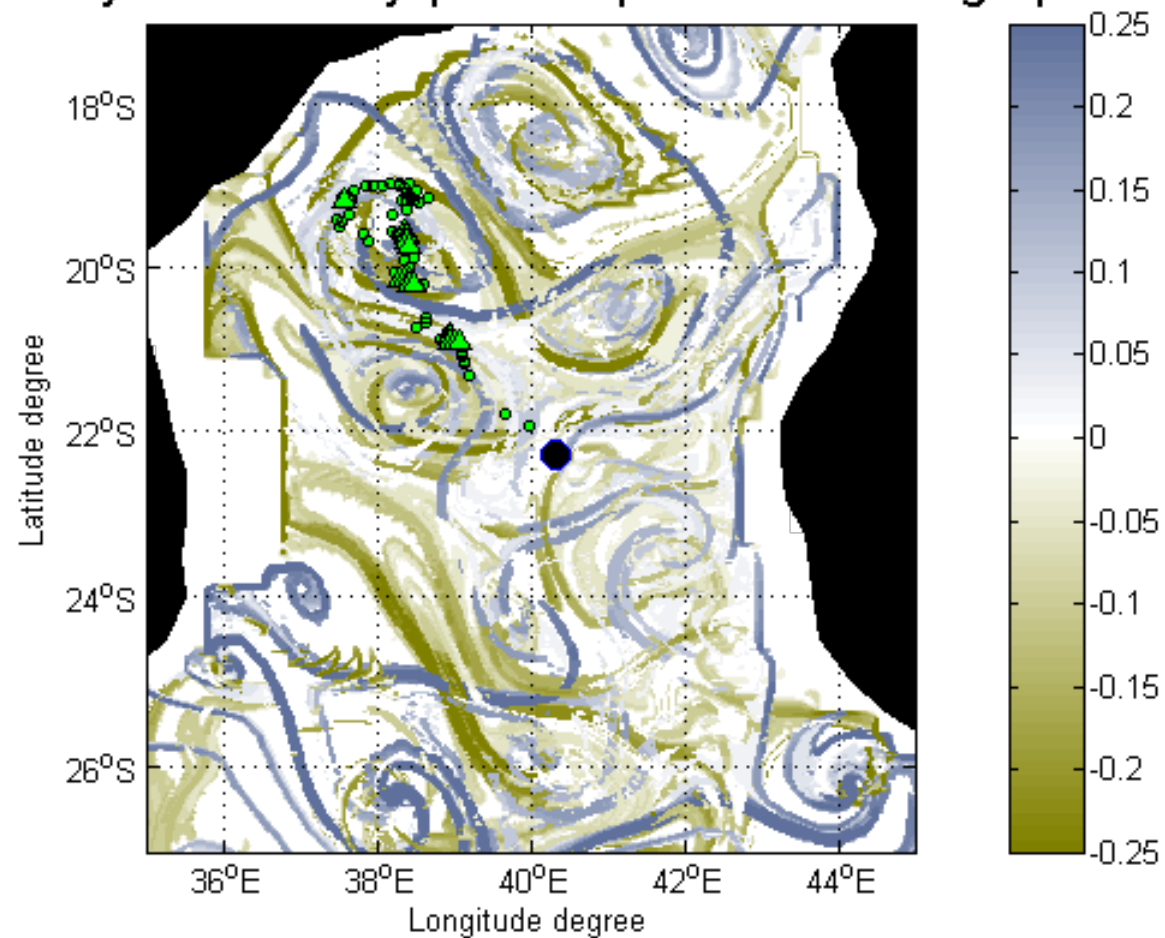
August 18 -
September 30,
2003.



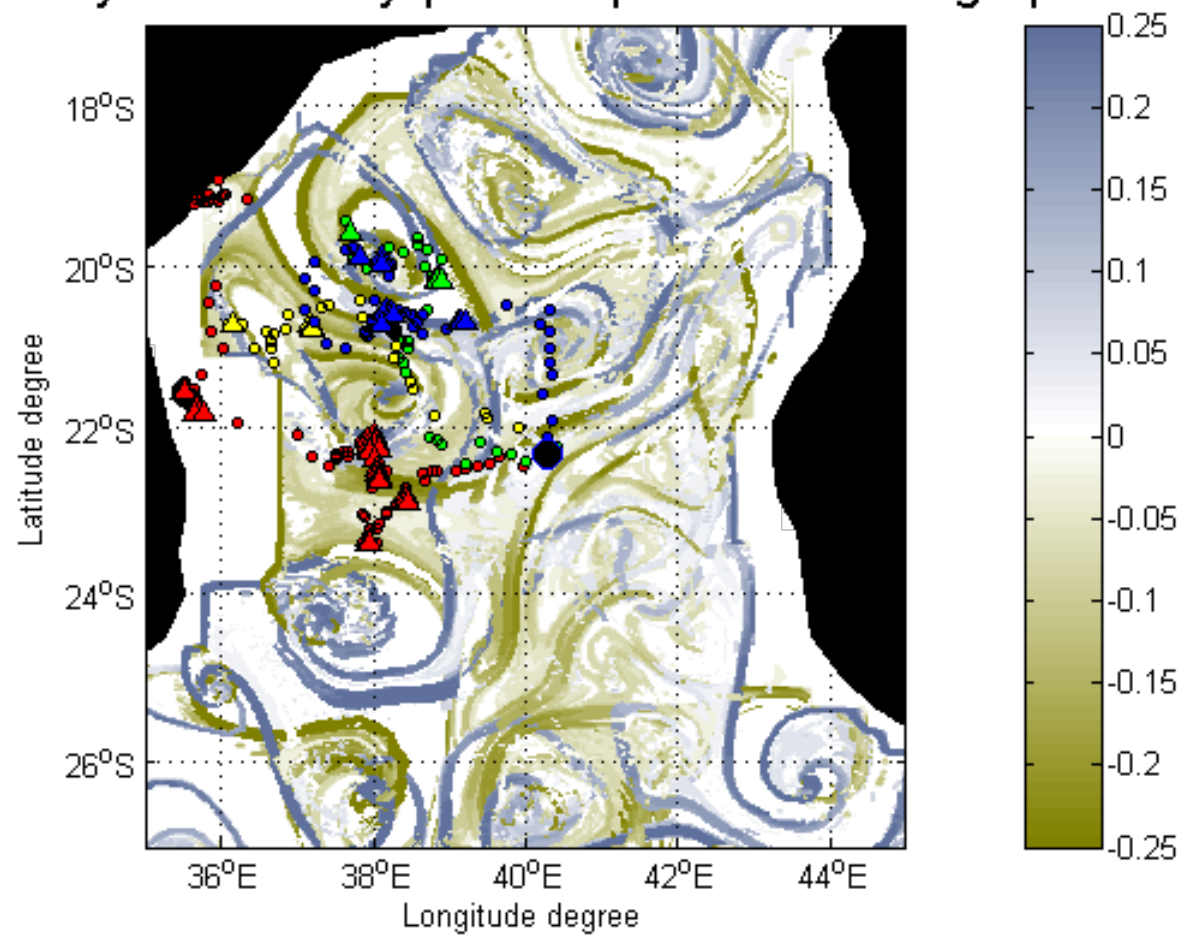
Overlay Finite Size Lyapunov Exponent - 1496 long trips



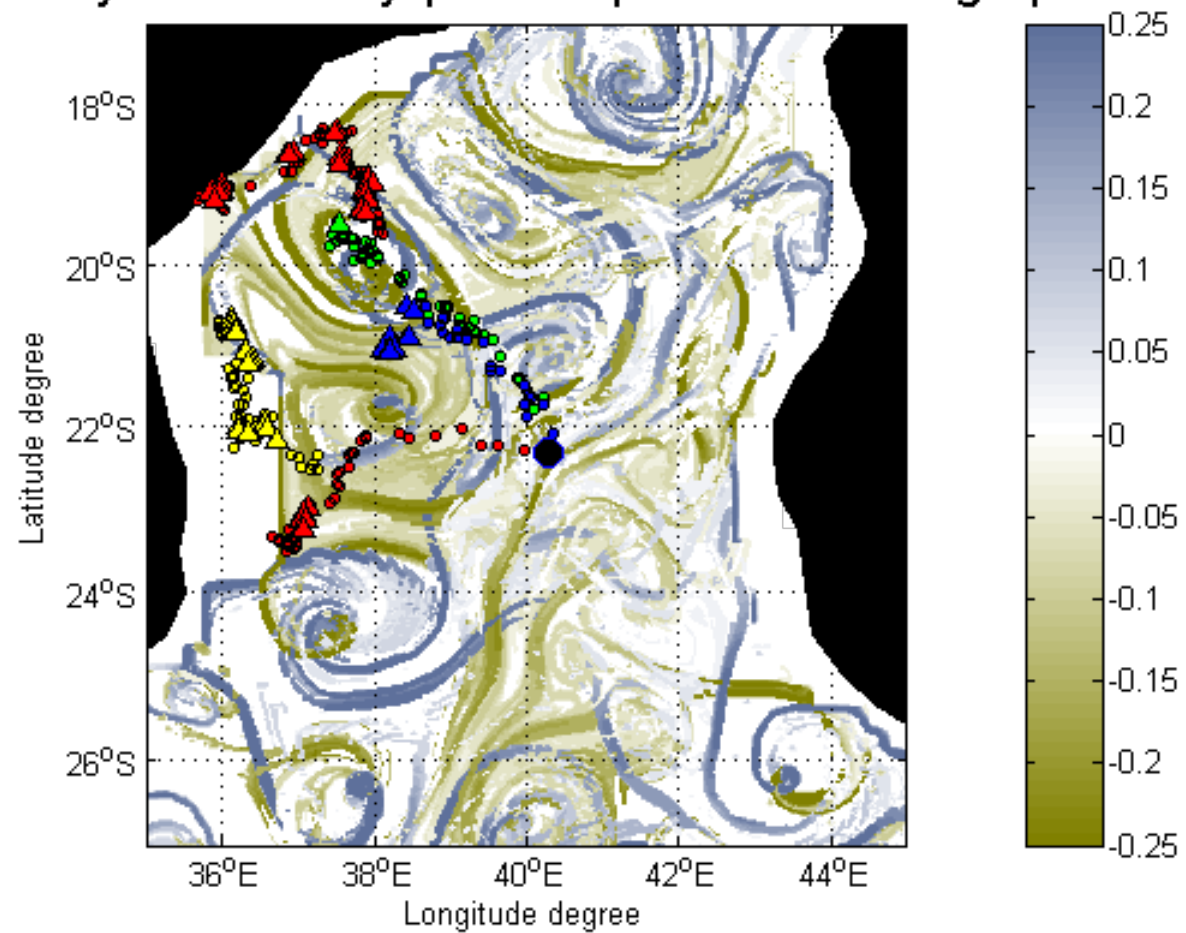
Overlay Finite Size Lyapunov Exponent -1500 long trips



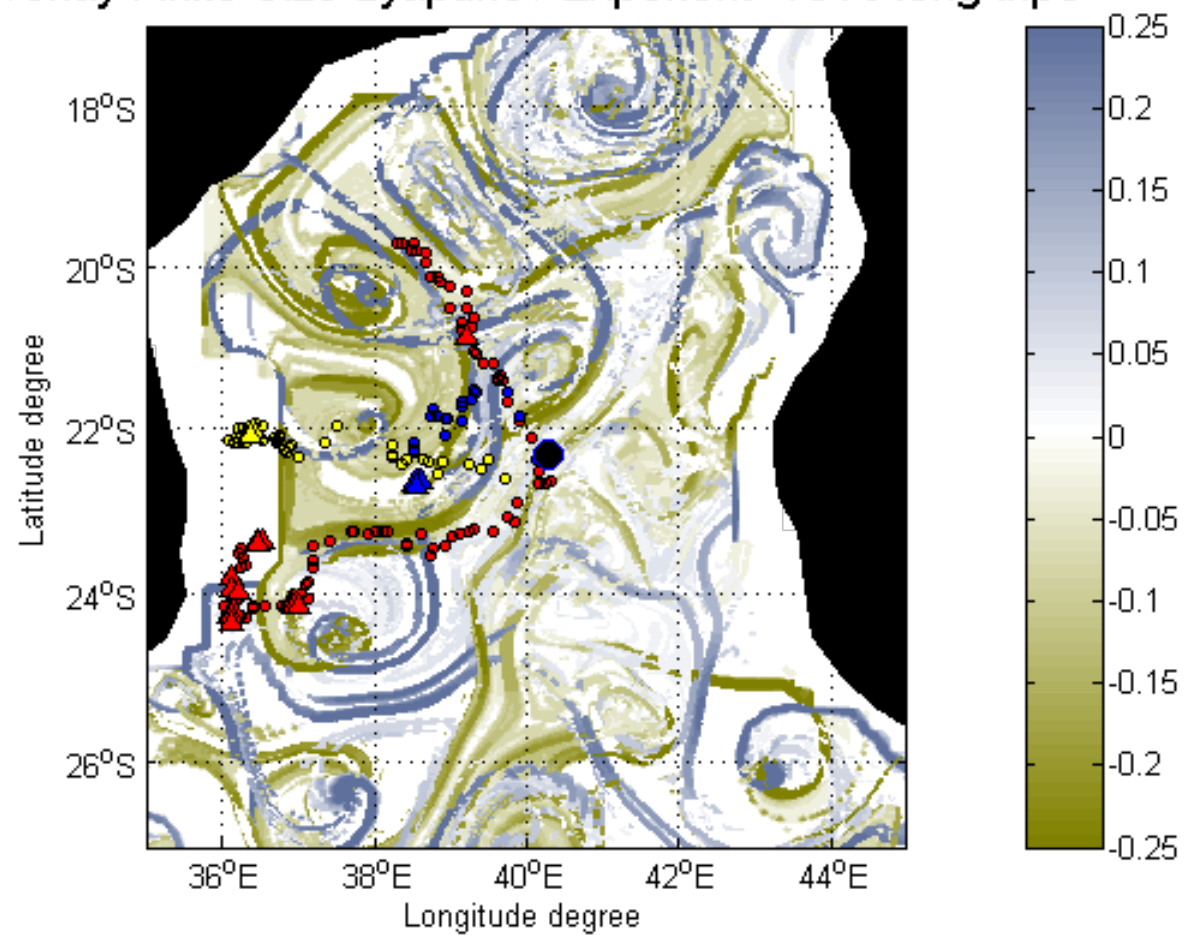
Overlay Finite Size Lyapunov Exponent -1508 long trips



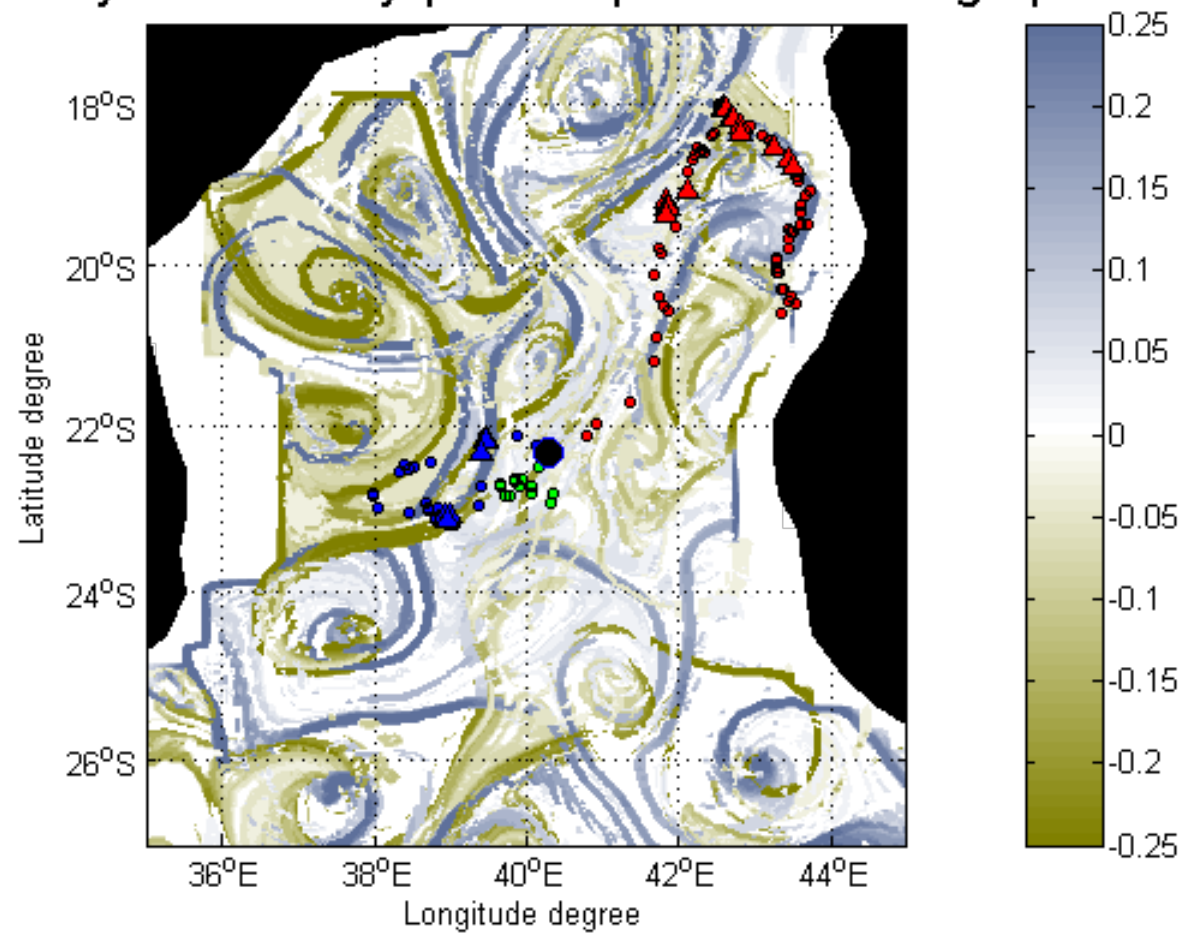
Overlay Finite Size Lyapunov Exponent -1512 long trips



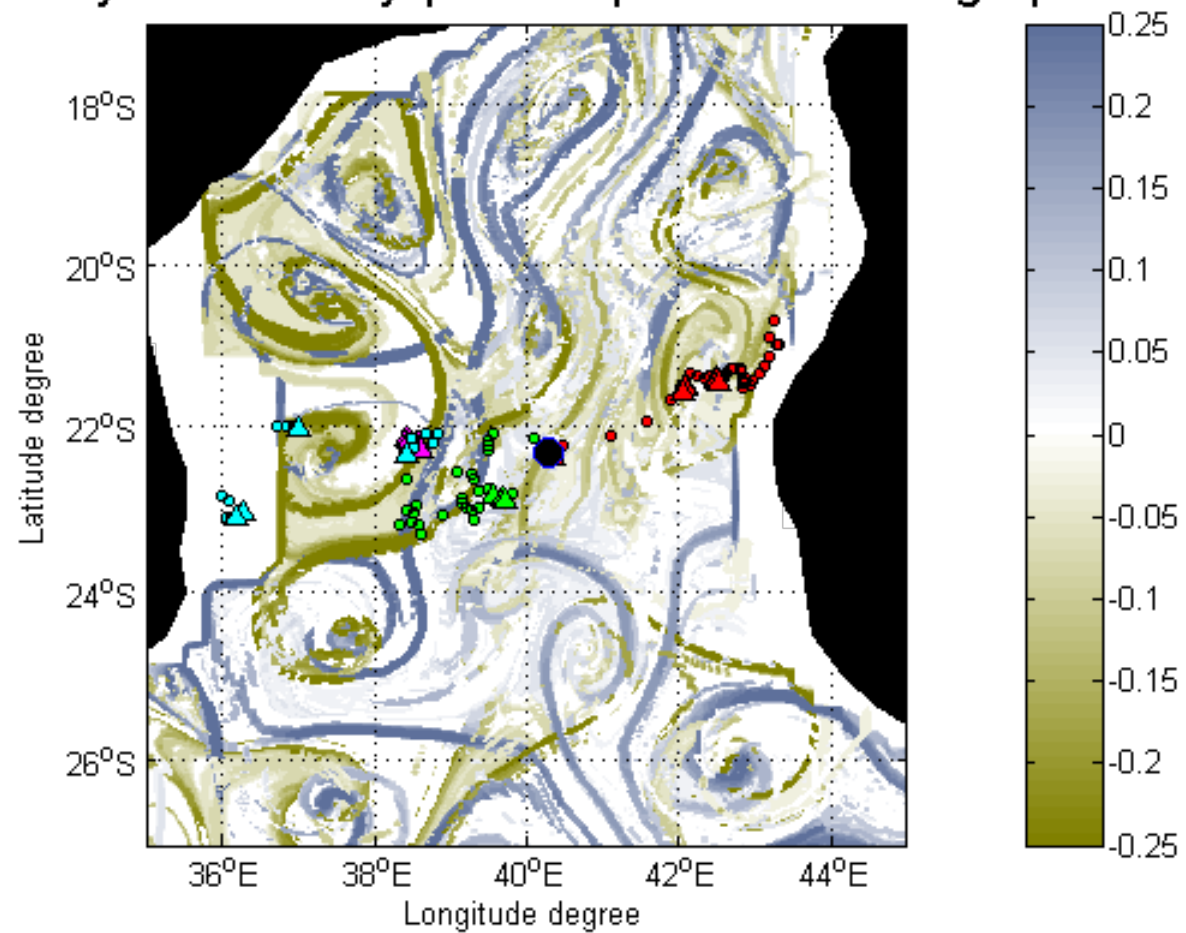
Overlay Finite Size Lyapunov Exponent -1516 long trips



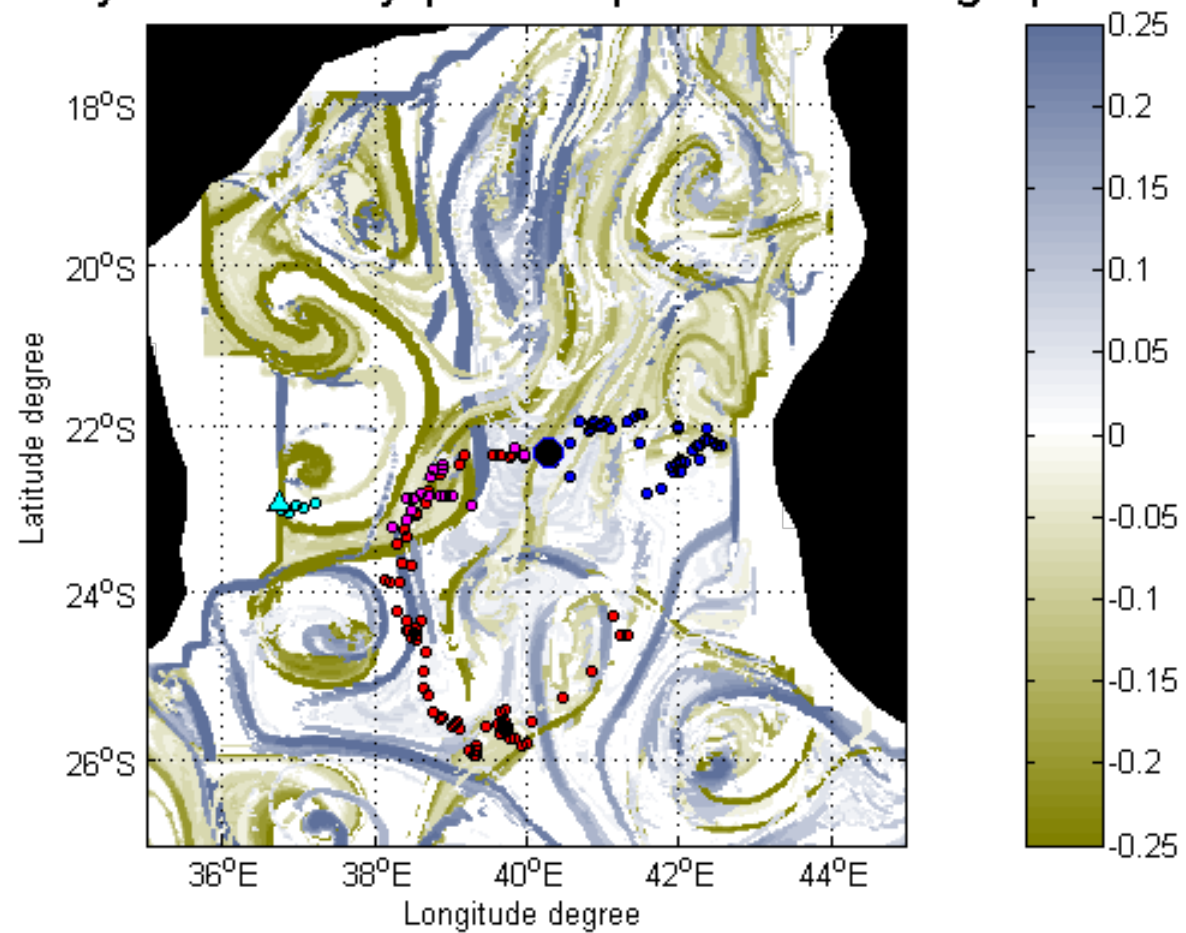
Overlay Finite Size Lyapunov Exponent -1520 long trips



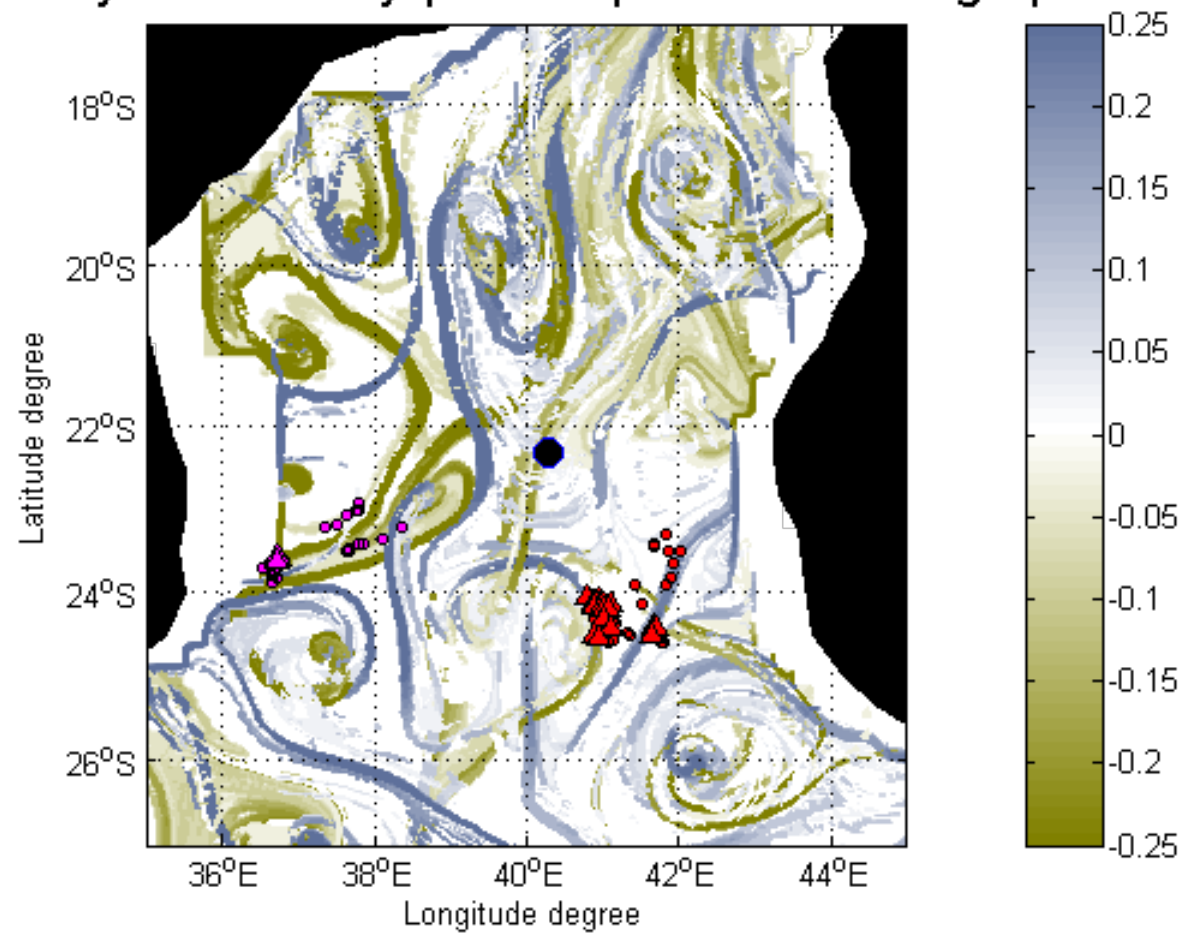
Overlay Finite Size Lyapunov Exponent -1524 long trips



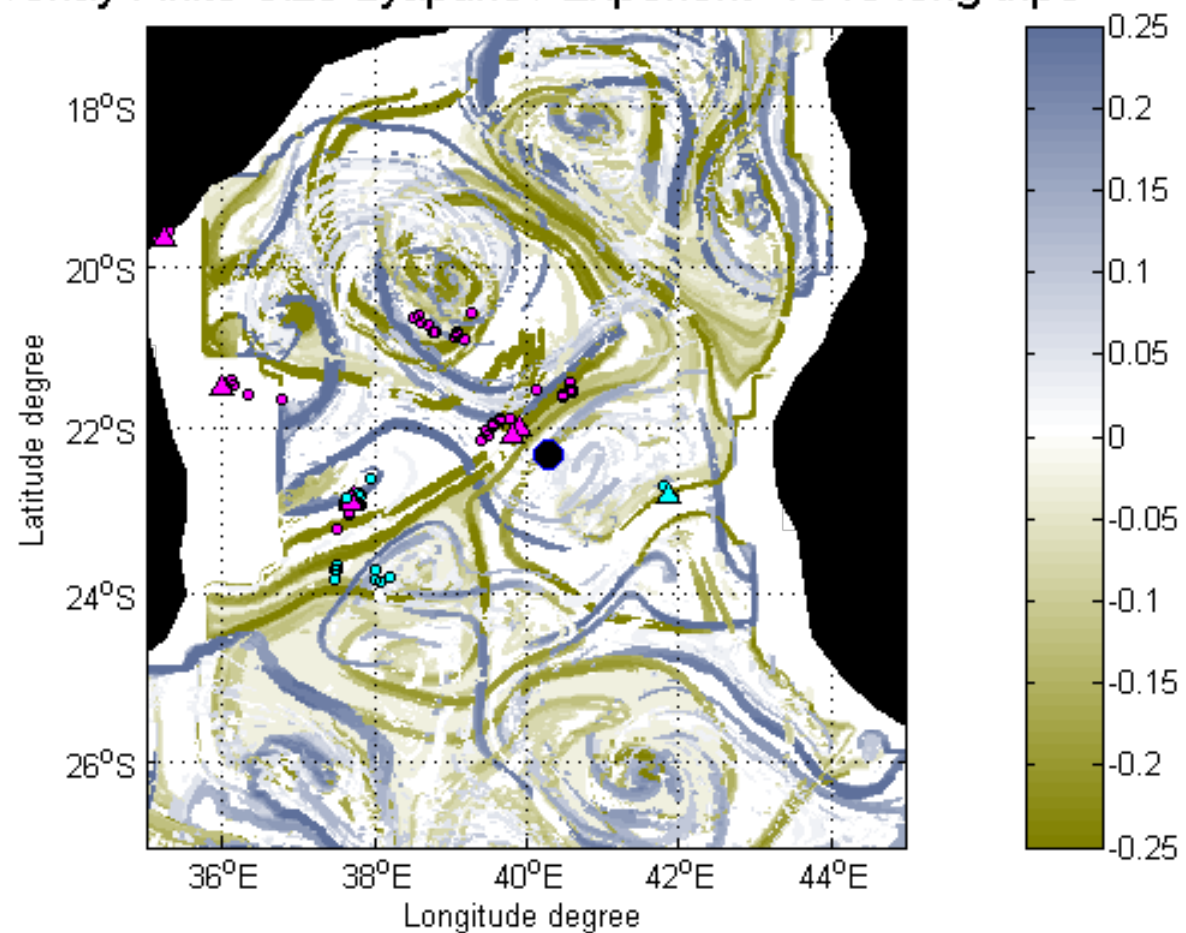
Overlay Finite Size Lyapunov Exponent -1528 long trips



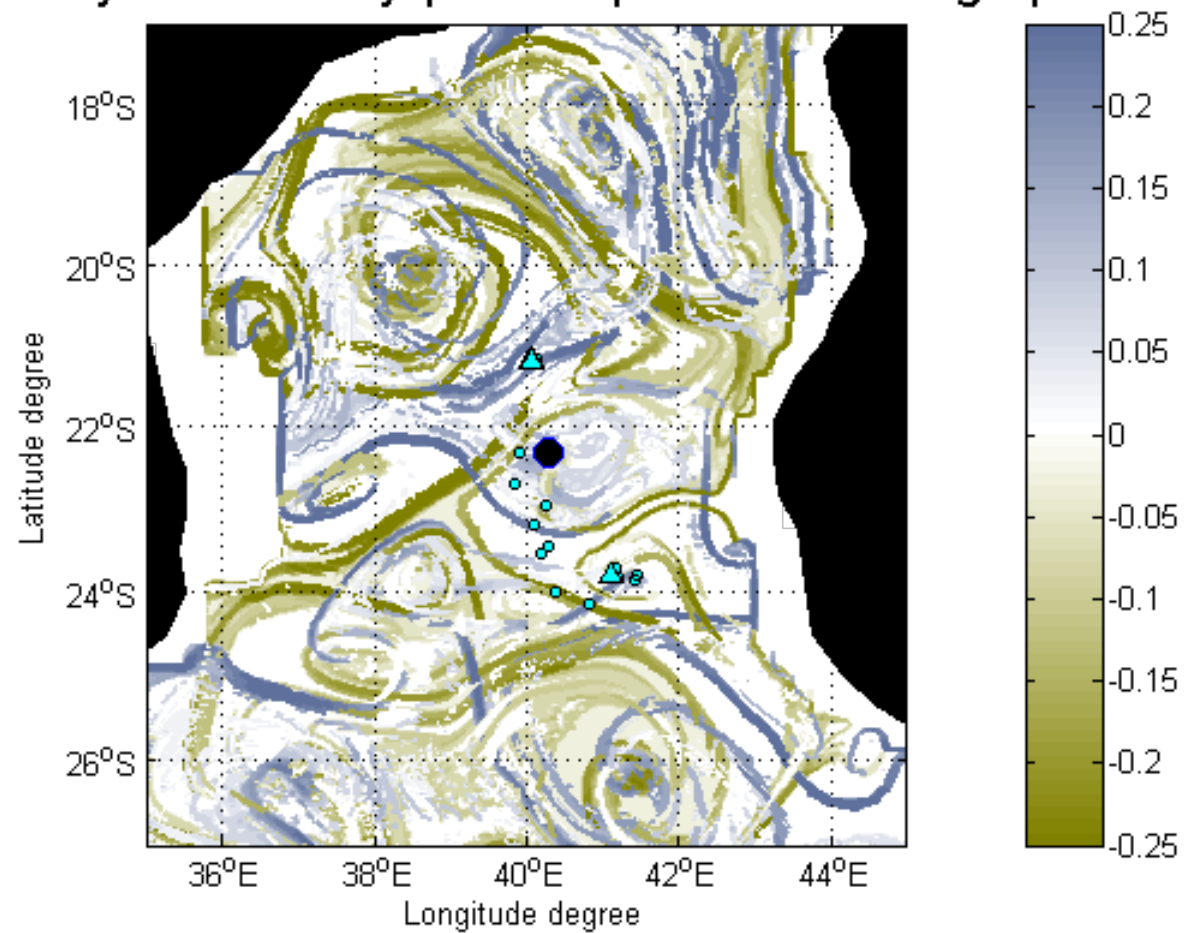
Overlay Finite Size Lyapunov Exponent -1532 long trips



Overlay Finite Size Lyapunov Exponent -1548 long trips



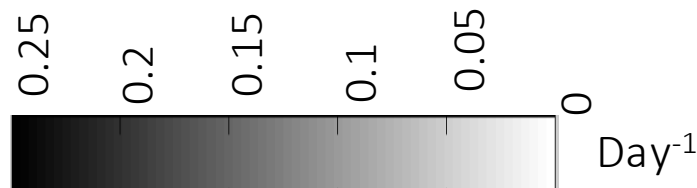
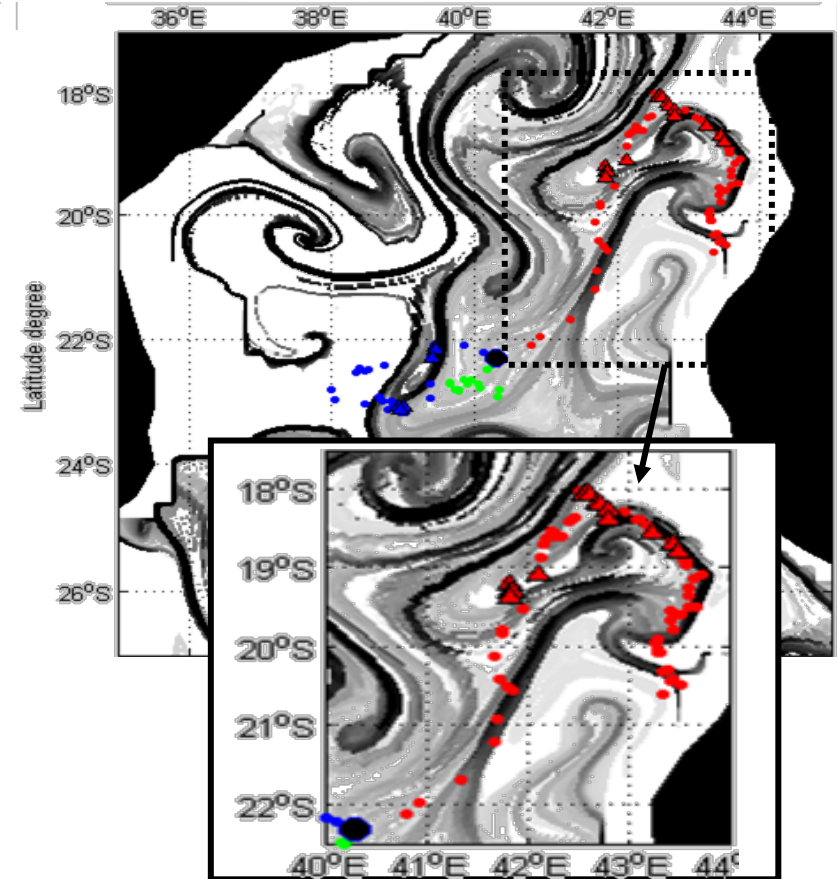
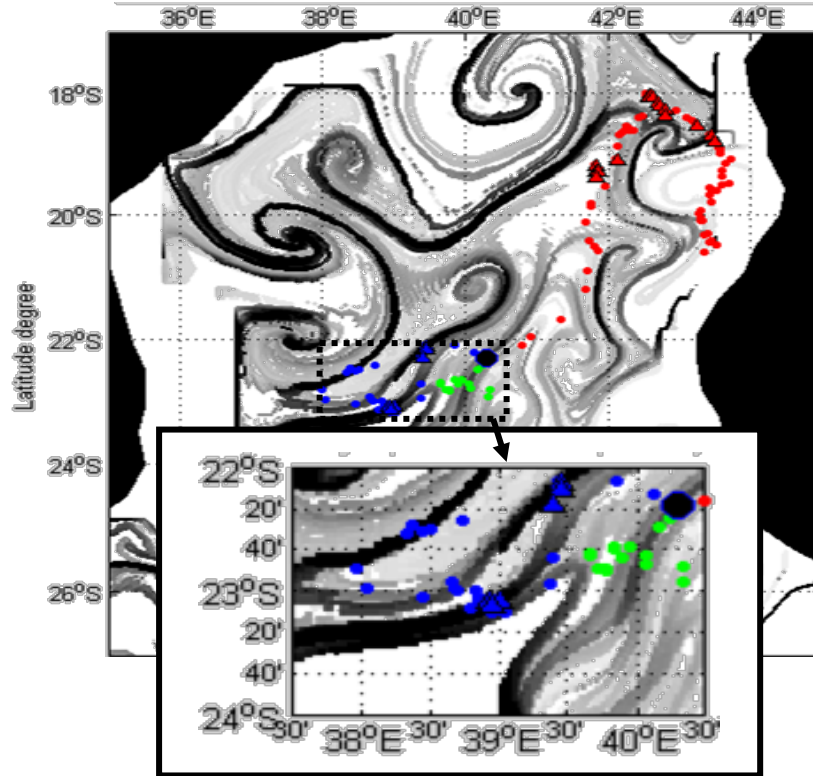
Overlay Finite Size Lyapunov Exponent -1552 long trips



Week of September 24, 2003

Backward FSLE=Attractive LCSs

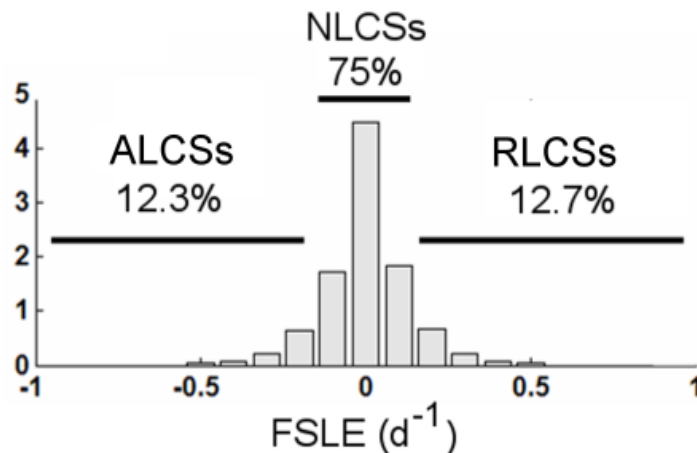
Forward FSLE = Repelling LCSs



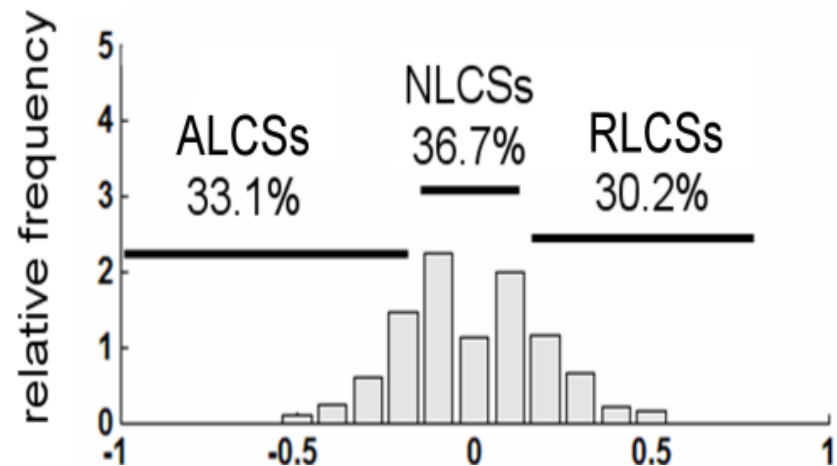
- ▲ foraging patch (flight speed lower than 10 km/h)
- seabird trajectory

Histograms of FSLE values

On the whole area



On the birds positions



ALCS: attracting LCS, i.e. FSLE (backwards) $< -0.1 \text{ day}^{-1}$

RLCS: repelling LCS, i.e. FSLE (forwards) $> 0.1 \text{ day}^{-1}$

NLCS: not LCS (small FSLE)

Despite LCS occupy only 25% of space, 63% of bird's positions are on them